

Table 18-1 Classical Physics

A Bit on Maxwell's Equations

Prerequisites.

- Poincaré's Lemma, which says that on \mathbb{R}^n , every closed form is exact. That is, if $d\omega = 0$, then there exists η with $d\eta = \omega$.
- Integration by parts: $\int \omega \wedge d\eta = -(-1)^{\deg \omega} \int (d\omega) \wedge \eta$ on domains that have no boundary.
- The Hodge star operator \star which satisfies $\omega \wedge \star \eta = \langle \omega, \eta \rangle dx_1 \cdots dx_n$ whenever ω and η are of the same degree.
- The simplest least action principle: the extremes of $q \mapsto \int_a^b (\frac{1}{2} m \dot{q}^2(t) - V(q(t))) dt$ occur when $m\ddot{q} = -V'(q(t))$. That is, when $F = ma$.

Maxwell's equations	
I. $\nabla \cdot E = \frac{\rho}{\epsilon_0}$	(Flux of E through a closed surface) = (Charge inside)/ ϵ_0
II. $\nabla \times E = -\frac{\partial B}{\partial t}$	(Line integral of E around a loop) = $-\frac{d}{dt}$ (Flux of B through the loop)
III. $\nabla \cdot B = 0$	(Flux of B through a closed surface) = 0
IV. $c^2 \nabla \times B = \frac{J}{\epsilon_0} + \frac{\partial E}{\partial t}$	c^2 (Integral of B around a loop) = (Current through the loop)/ ϵ_0 + $\frac{\partial}{\partial t}$ (Flux of E through the loop)
<div> <div>Conservation of charge</div> <div>$\nabla \cdot j = -\frac{\partial \rho}{\partial t}$</div> <div>(Flux of current through a closed surface) = $-\frac{\partial}{\partial t}$ (Charge inside)</div> </div>	
Force law	
$F = q(E + v \times B)$	
Law of motion	
$\frac{d}{dt}(p) = F$, where $p = \frac{mv}{\sqrt{1-v^2/c^2}}$ (Newton's law, with Einstein's modification)	
Gravitation	
$F = -G \frac{m_1 m_2}{r^2} e_r$	

The Feynman Lectures on Physics vol. II, page 18-2

The Action Principle. The *Vector Field* is a compactly supported 1-form A on \mathbb{R}^4 which extremizes the *action*

$$S_J(A) := \int_{\mathbb{R}^4} \frac{1}{2} \|dA\|^2 dt dx dy dz + J \wedge A$$

where the 3-form J is the *charge-current*.

The Euler-Lagrange Equations in this case are $d \star dA = J$, meaning that there's no hope for a solution unless $dJ = 0$, and that we might as well (think Poincaré's Lemma!) change variables to $F := dA$. We thus get

$$dJ = 0 \quad dF = 0 \quad d \star F = J$$

These are the Maxwell equations! Indeed, writing $F = (E_x dx dt + E_y dy dt + E_z dz dt) + (B_x dy dz + B_y dz dx + B_z dx dy)$ and $J = \rho dx dy dz - j_x dy dz dt - j_y dz dx dt - j_z dx dy dt$, we find:

$dJ = 0 \implies$	$\frac{\partial \rho}{\partial t} + \text{div } j = 0$	"conservation of charge"
$dF = 0 \implies$	$\text{div } B = 0$	"no magnetic monopoles"
	$\text{curl } E = -\frac{\partial B}{\partial t}$	that's how generators work!
$d \star F = J \implies$	$\text{div } E = -\rho$	"electrostatics"
	$\text{curl } B = -\frac{\partial E}{\partial t} + j$	that's how electromagnets work!

Exercise. Use the Lorentz metric to fix the sign errors.

Exercise. Use pullbacks along Lorentz transformations to figure out how E and B (and j and ρ) appear to moving observers.

Exercise. With $ds^2 = c^2 dt^2 - dx^2 - dy^2 - dz^2$ use $S = mc \int_{e_1}^{e_2} (ds + eA)$ to derive Feynman's "law of motion" and "force law".

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