

Numerical Simulations of Radial Supercritical Defocusing Waves

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This talk reports on [CSS09] with **G. Simpson**, C. Sulem (U. Toronto) and forthcoming work [CPSS10] also with **P. Pang**.

- 1 Energy Supercritical NLS**
 - Critical Sobolev Regularity
 - Conservation Laws
 - Critical Regimes
- 2 Energy Supercritical Maximal-in-time Theory?**
 - Difficulties in the Energy Supercritical Setting
 - Evolutions with bounded critical norm scatter
- 3 Numerical Simulations of Supercritical Waves**
 - Strauss-Vazquez 1978 Simulations
 - [CSS09] Simulations
 - [CPSS10] Simulations
- 4 Ideas toward Supercritical Maximal-in-time Theory?**

1. Energy Supercritical NLS

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Consider the defocusing monomial NLS initial value problem:

$$\begin{cases} (i\partial_t + \Delta)u = |u|^{p-1}u \\ u(0, x) = u_0(x). \end{cases} \quad (\text{NLS}_p^+(d))$$

Dilation Invariance:

$$u : [0, T] \times \mathbb{R}^d \mapsto \text{solves } \text{NLS}_p^+(d)$$

\Updownarrow

$\boxed{\mathbb{R}^d}$

$$\forall \lambda > 0, u_\lambda : [0, \lambda^2 T] \times \mathbb{R}^d \mapsto \text{solves } \text{NLS}_p^+(d)$$

where

$$u_\lambda(\tau, y) = \left(\frac{1}{\lambda}\right)^{\frac{2}{p-1}} u\left(\frac{\tau}{\lambda^2}, \frac{y}{\lambda}\right).$$

Critical Sobolev Regularity

A simple calculation shows that

$$\|D^s u_\lambda(\tau, \cdot)\|_{L^2} = \left(\frac{1}{\lambda}\right)^{\frac{2}{p-1} + s - \frac{d}{2}} \|D^s u(\tau)\|_{L^2}.$$

We encounter a **dilation invariant Sobolev space norm** when

$$s = s_c = \frac{d}{2} - \frac{2}{p-1}.$$

The space $\dot{H}^{s_c}(^d)$ plays a basic role in theory for $NLS_p(^d)$.

Conservation Laws for $NLS_p^+(\mathbb{R}^d)$

Time invariant quantities:

$$\text{Mass} = \int |u(t, x)|^2 dx.$$

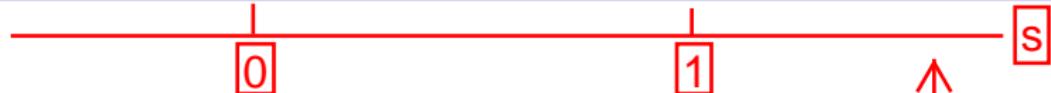
$$\text{Momentum} = 2\Im \int \bar{u}(t) \partial_x u(t) dx.$$

$$\text{Energy} = H[u(t)] = \int_T \frac{1}{2} |\partial_x u(t)|^2 dx + \frac{2}{p+1} |u(t)|^{p+1} dx.$$

Sobolev Controls:

- Mass controls the L^2 (a.k.a. H^0) norm.
- Momentum is at the $H^{1/2}$ regularity level.
- Energy controls H^1 norm.

Criticality Regimes $NLS_p^+(\mathbb{R}^d)$



Recall the critical Sobolev regularity index for $NLS_p^+(\mathbb{R}^d)$:

$$s = s_c = \frac{d}{2} - \frac{2}{p-1}.$$

Conservation laws identify 5 critical regimes:

- $s_c < 0$, Mass Subcritical
- $s_c = 0$, Mass Critical
- $0 < s_c < 1$, Mass Supercritical/Energy Subcritical
- $s_c = 1$, Energy Critical
- $s_c > 1$, Energy Supercritical.

Global Wellposedness for Energy Subcritical Regime

- H^1 -Local Wellposedness Lifetime depends on H^1 norm:

$$T_{lwp} \sim \|u(0)\|_{H^1}^{-\gamma}.$$

- Energy Conservation implies H^1 *a priori* control:

$$\|u(t)\|_{H^1}^2 \leq \text{Energy}[u(0)], \forall t.$$

- Iterate the local theory with **uniform** time steps \implies GWP.

Remarks:

- Maximal-in-time behavior of solutions? Not clear...
- Low regularity relaxations (high/low frequency, I -method)
- Polynomial in time bounds on $\|u(t)\|_H^s$, $s \gg 1$.

GWP for Energy Critical Regime

LWP iteration FAILS:

- Energy implies H^1 *a priori* control: $\|u(t)\|_{H^1}^2 \leq \text{Energy}[u(0)]$.
- H^1 -LWP lifetime not controlled by H^1 norm.
- H^1 -LWP iteration does NOT imply GWP.

Energy Critical Scattering via Critical Elements Method:

- Induction on Energy Approach [Bou99]
 - $NLS_5^+(\mathbb{R}^3)$, $NLS_3^+(\mathbb{R}^4)$ radial.
 - $NLS_5^+(\mathbb{R}^3)$ GWP and Scatters [CKS⁺08].
 - 4 and higher dimensions [RV05], [Vis07].
- Extract and kill critical elements. [KM08a]
 - Progress on focusing problems beneath soliton threshold.
 - Mass Critical Case [KVZ08], [Dod10]
 - Other equations, including Navier-Stokes [KK09]

2. Energy Supercritical Maximal-in-time Theory?

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- This theory is **wide open**.
- We have no *a priori* control on Sobolev norms with $s > 1$.
- Same issue obstructs Navier-Stokes initial value problem?
- Idea? First resolve $NLS_3^+(\mathbb{R}^3)$? Use the momentum?

Critical Norm Bounded \implies Scattering

Bounded H^{s_c} norm **Assumption:** Suppose $H^{s_c} \ni u(0) \mapsto u$ solves NLS or NLW with $s_c > 1$ and $\exists K < 0$ such that

$$\sup_t \|u(t)\|_{H^{s_c}} < K.$$

- $NLW_p^{(d)}$:
Radial + Bounded H^{s_c} norm \implies scattering [KM08b].
- $NLS_p^+^{(d)}$: Bounded H^{s_c} norm \implies scattering [KV10].

Question: What is the behavior of $\|u(t)\|_{H^{s_c}}$?

3. Numerical Simulations of Supercritical Waves

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- Supercritical NLW behaves nicely [SV78].
- $NLS_5^+(^5)$ has bounded H^2 norm [CSS09].
- Energy supercritical NLW behaves similarly [CPSS10].

Strauss-Vazquez 1978 simulations

JOURNAL OF COMPUTATIONAL PHYSICS 28, 271-278 (1978)

Note

Numerical Solution of a Nonlinear Klein-Gordon Equation

We compute the solutions of the equation $u_{tt} - \Delta u + m^2 u + g u^p = 0$ for p odd and $m, g > 0$. Our computations show that (i) the solutions remain bounded as $t \rightarrow \infty$, (ii) the amplitude decreases as p increases, and (iii) the number of oscillations increases as p increases. Because of (i), theoretical results imply that the amplitude goes to zero like $O(t^{-3/2})$ as $t \rightarrow \infty$.

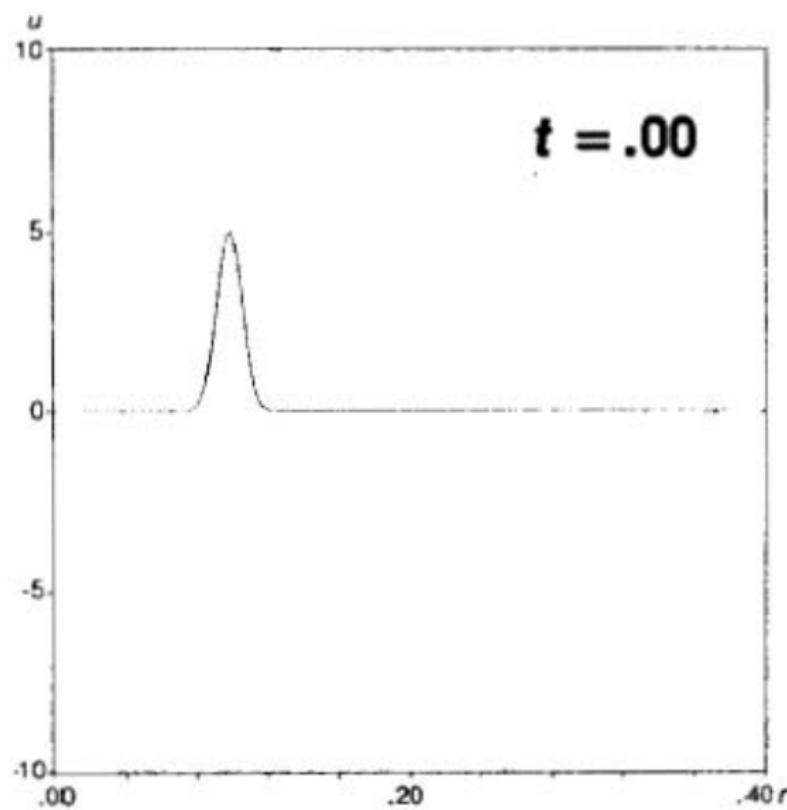
1. ANALYSIS

The nonlinear Klein-Gordon equation (NLKG)

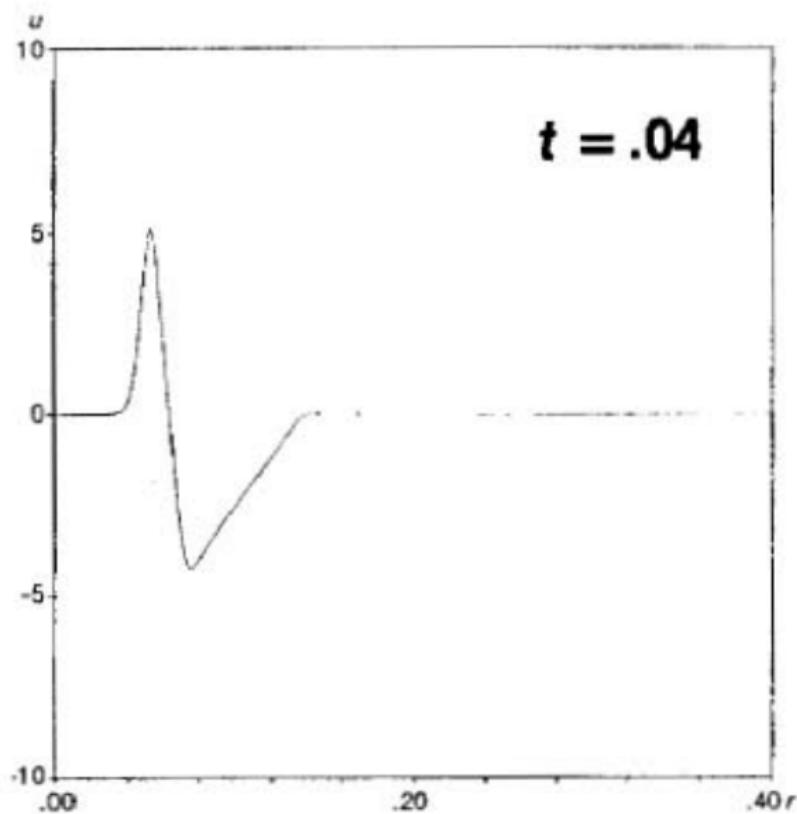
$$u_{tt} - \Delta u + m^2 u + G'(u) = 0, \quad (1)$$

$$u(x, 0) = \phi(x), \quad u_t(x, 0) = \Psi(x) \quad (2)$$

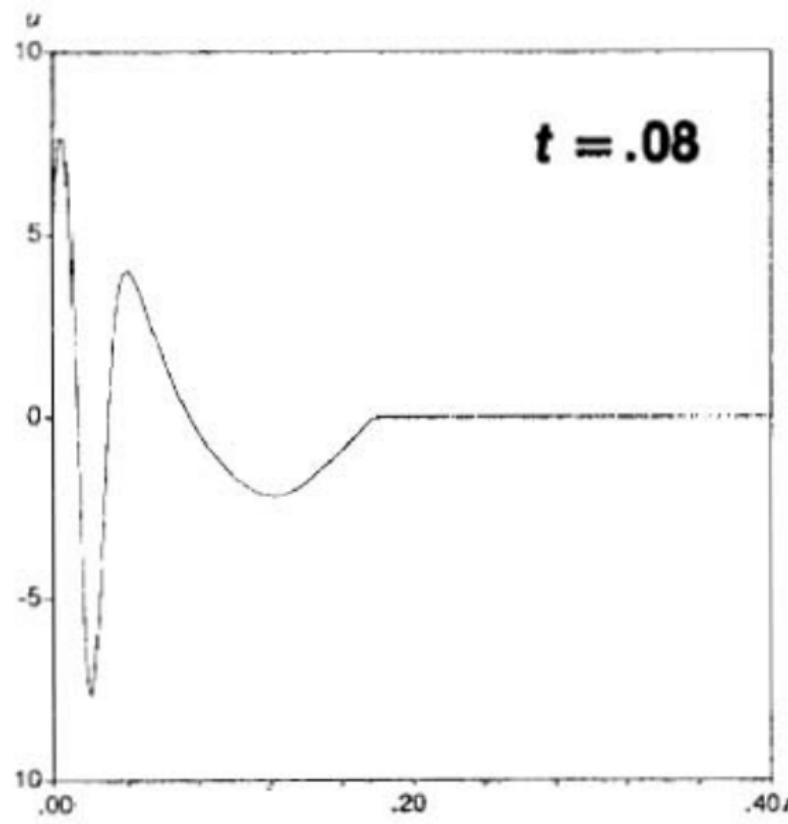
Strauss-Vazquez 1978 simulations



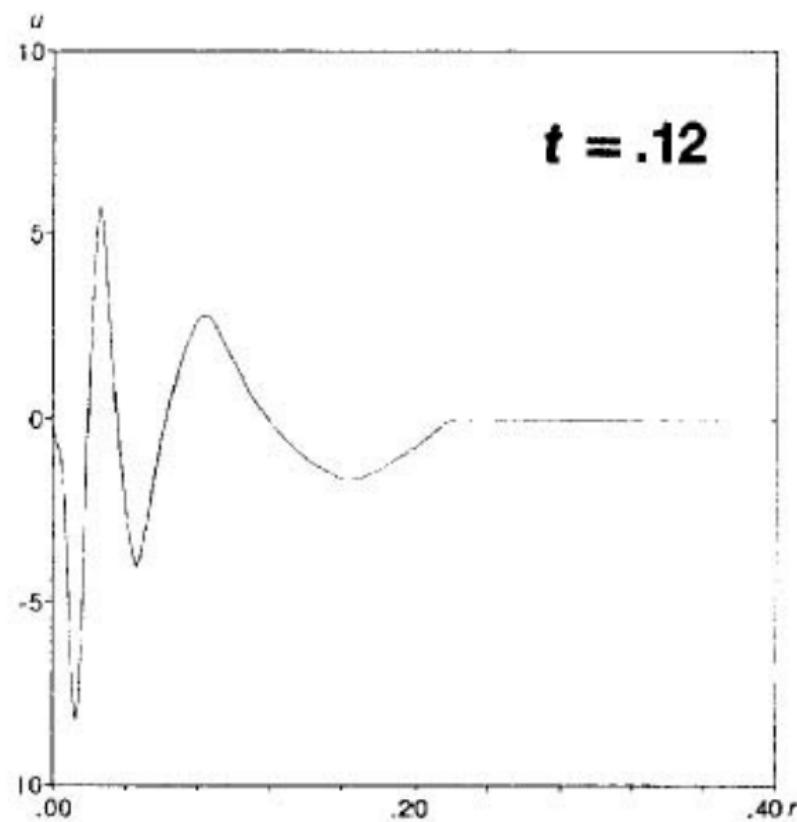
Strauss-Vazquez 1978 simulations



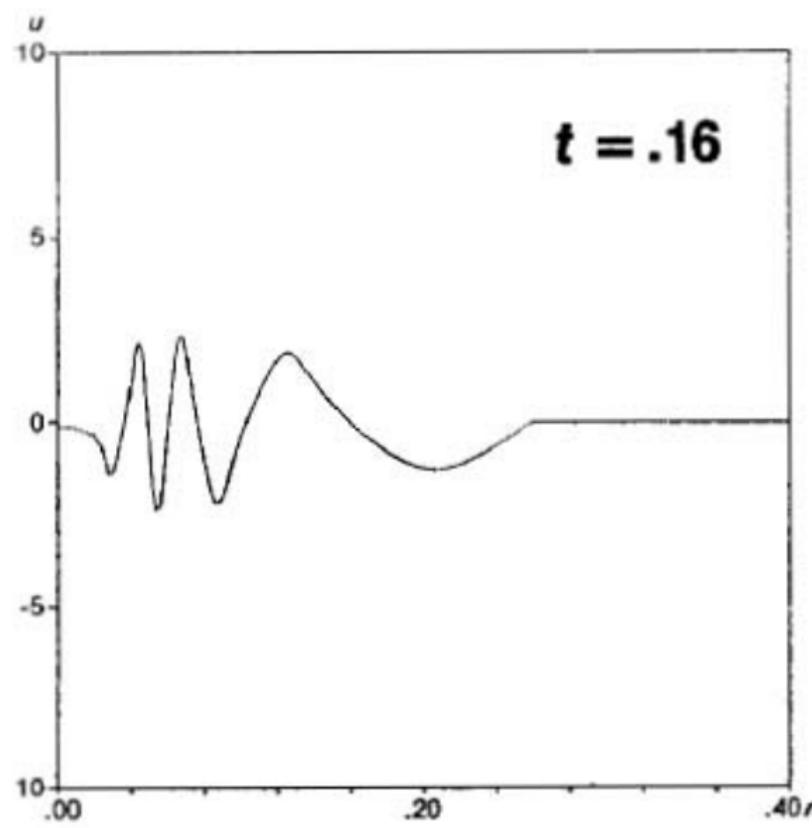
Strauss-Vazquez 1978 simulations



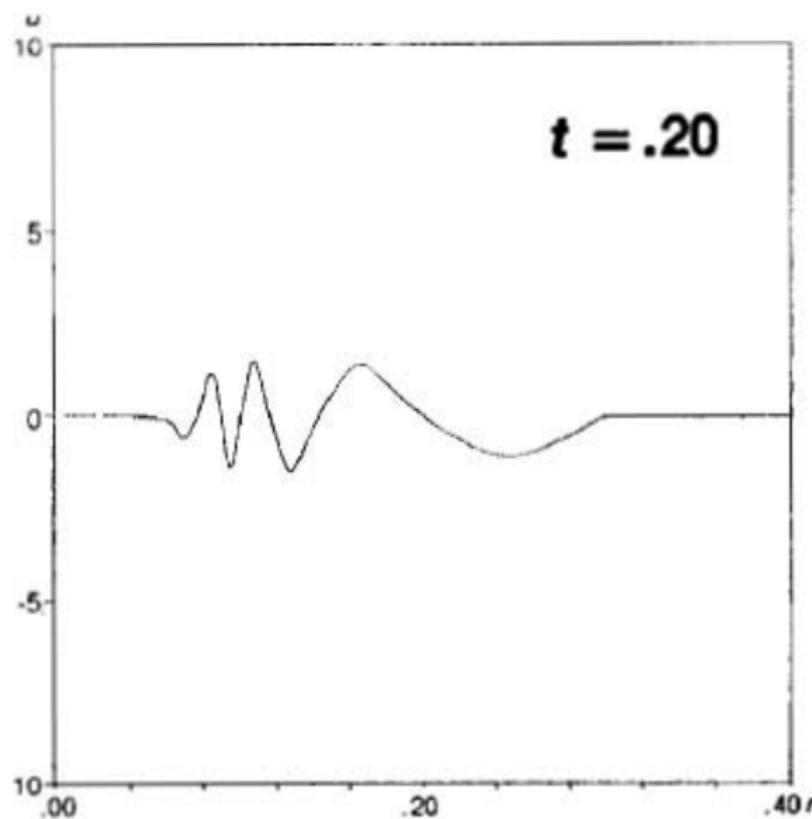
Strauss-Vazquez 1978 simulations



Strauss-Vazquez 1978 simulations



Strauss-Vazquez 1978 Simulations



C-Simpson-Sulem simulations

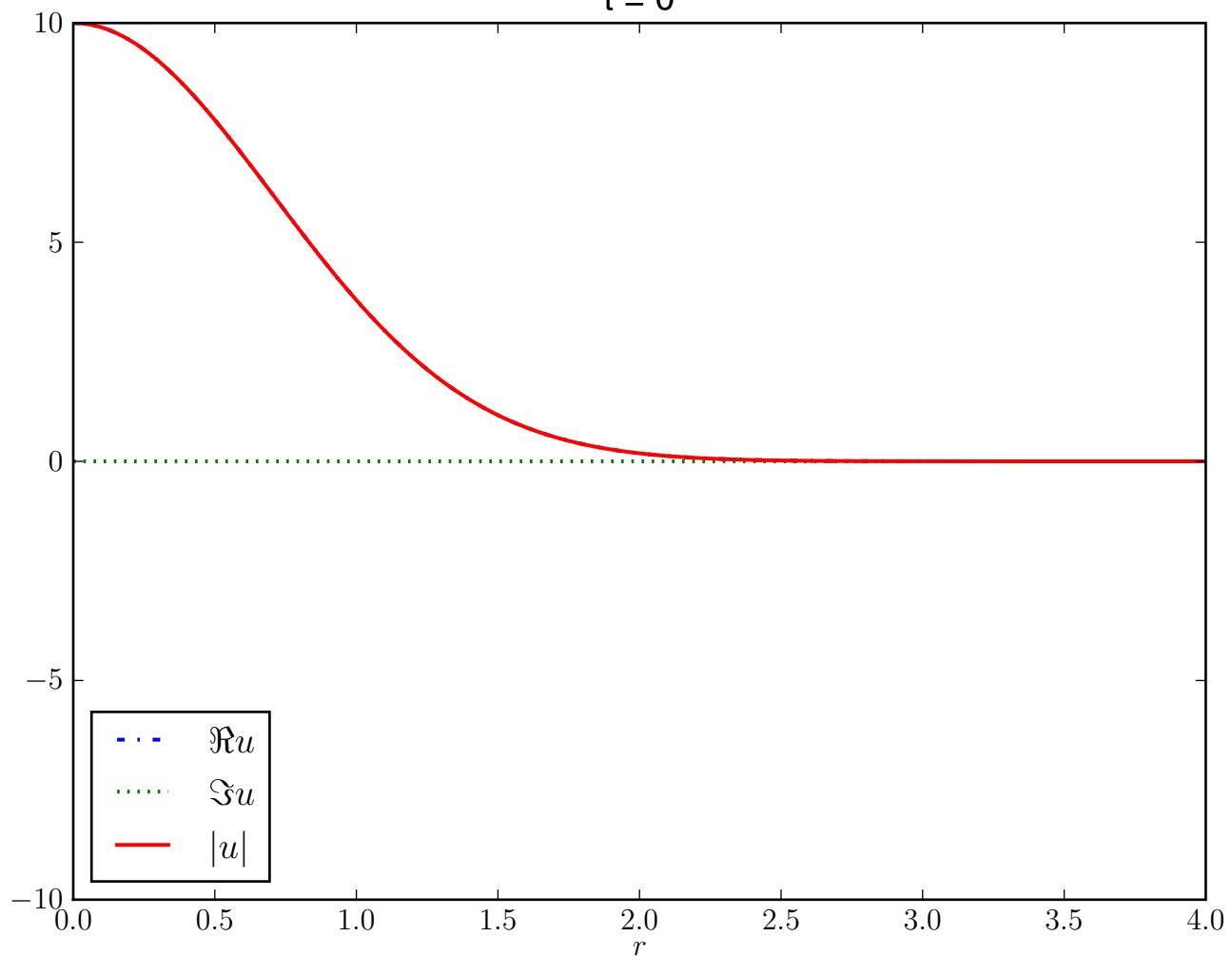
C-Simpson-Sulem simulations

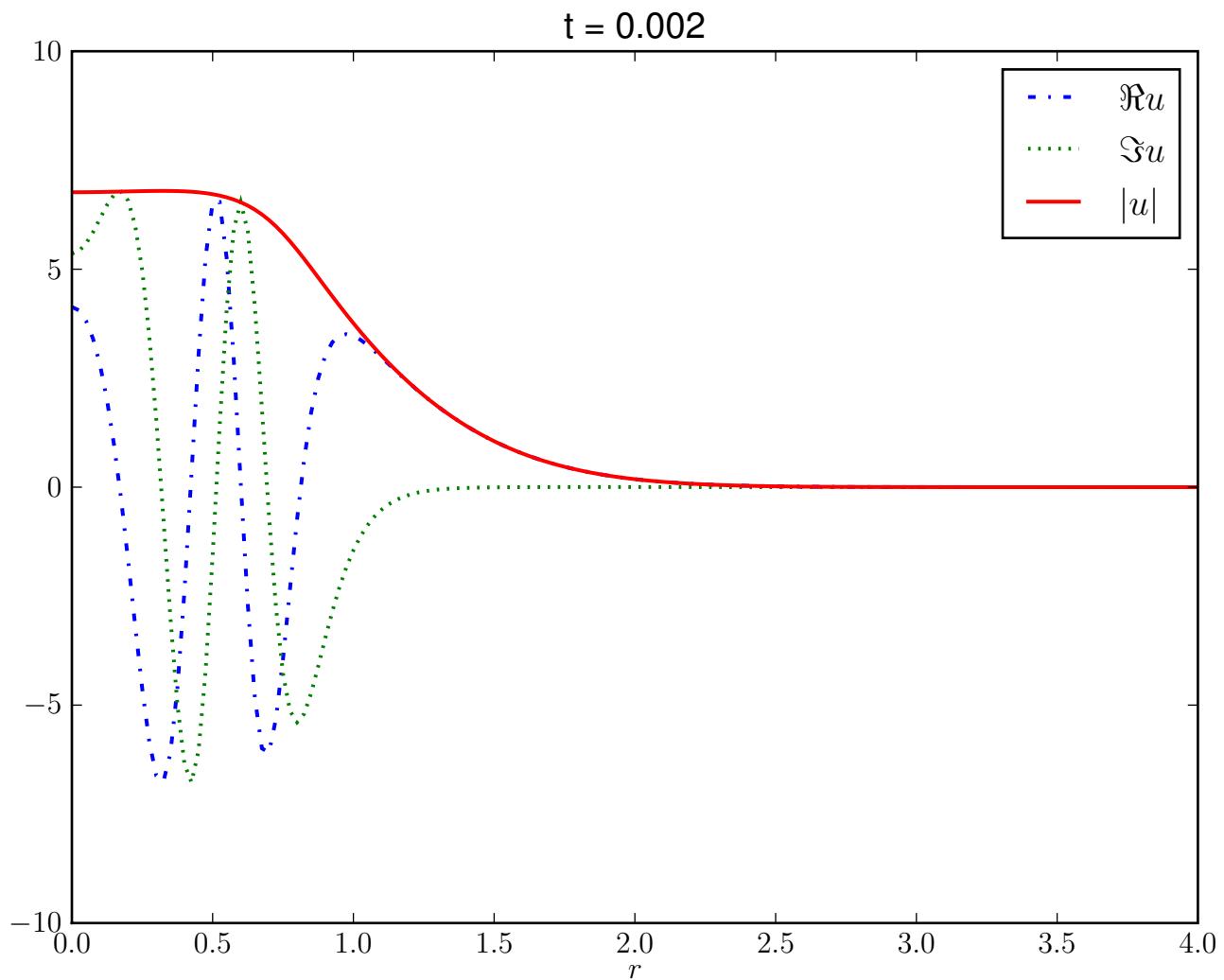
Four simulations of radial data:

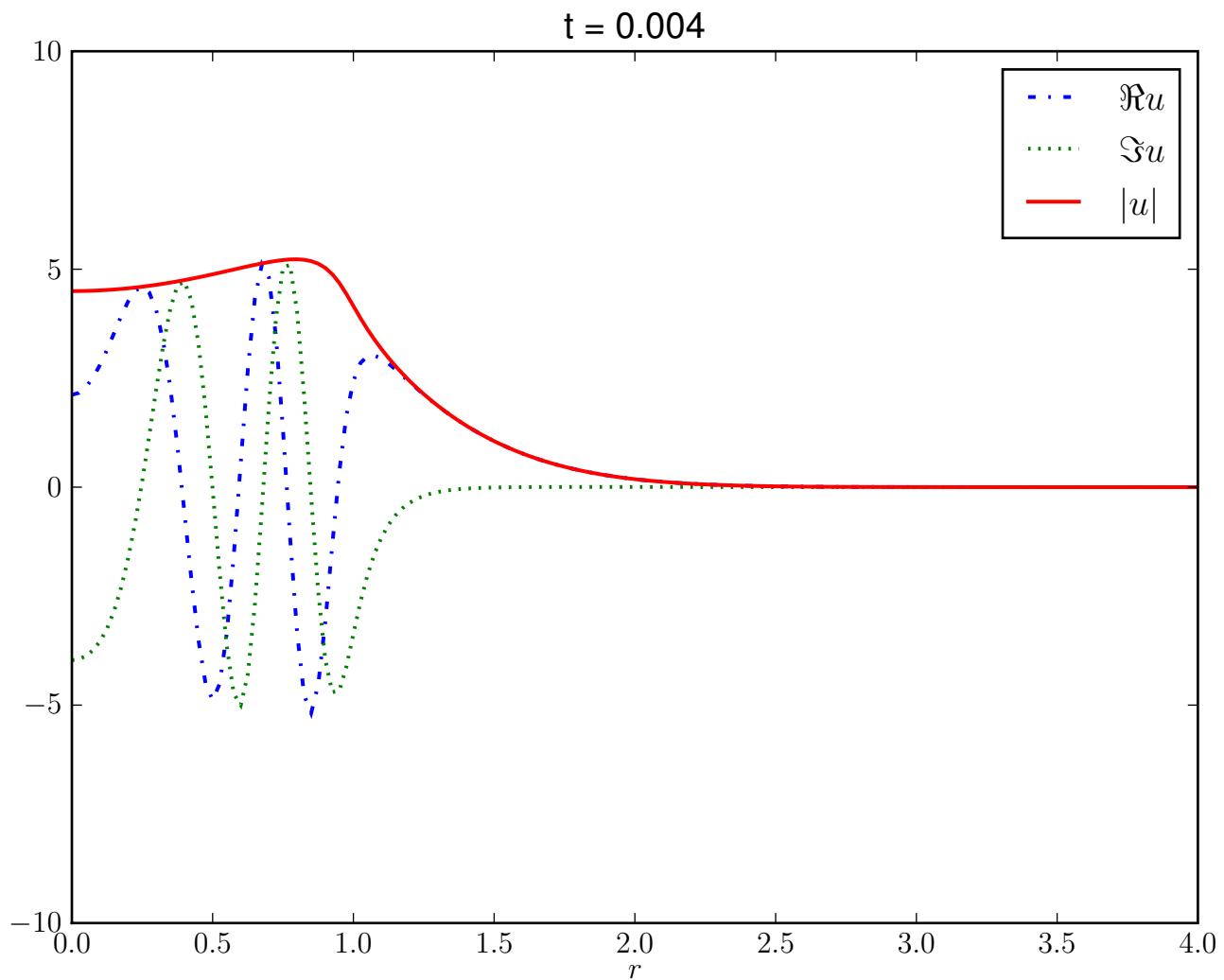
- Centered Gaussian
- Phased Centered Gaussian
- Phased Centered Gaussian (Linear flow diagnostic)
- Spherical Ring

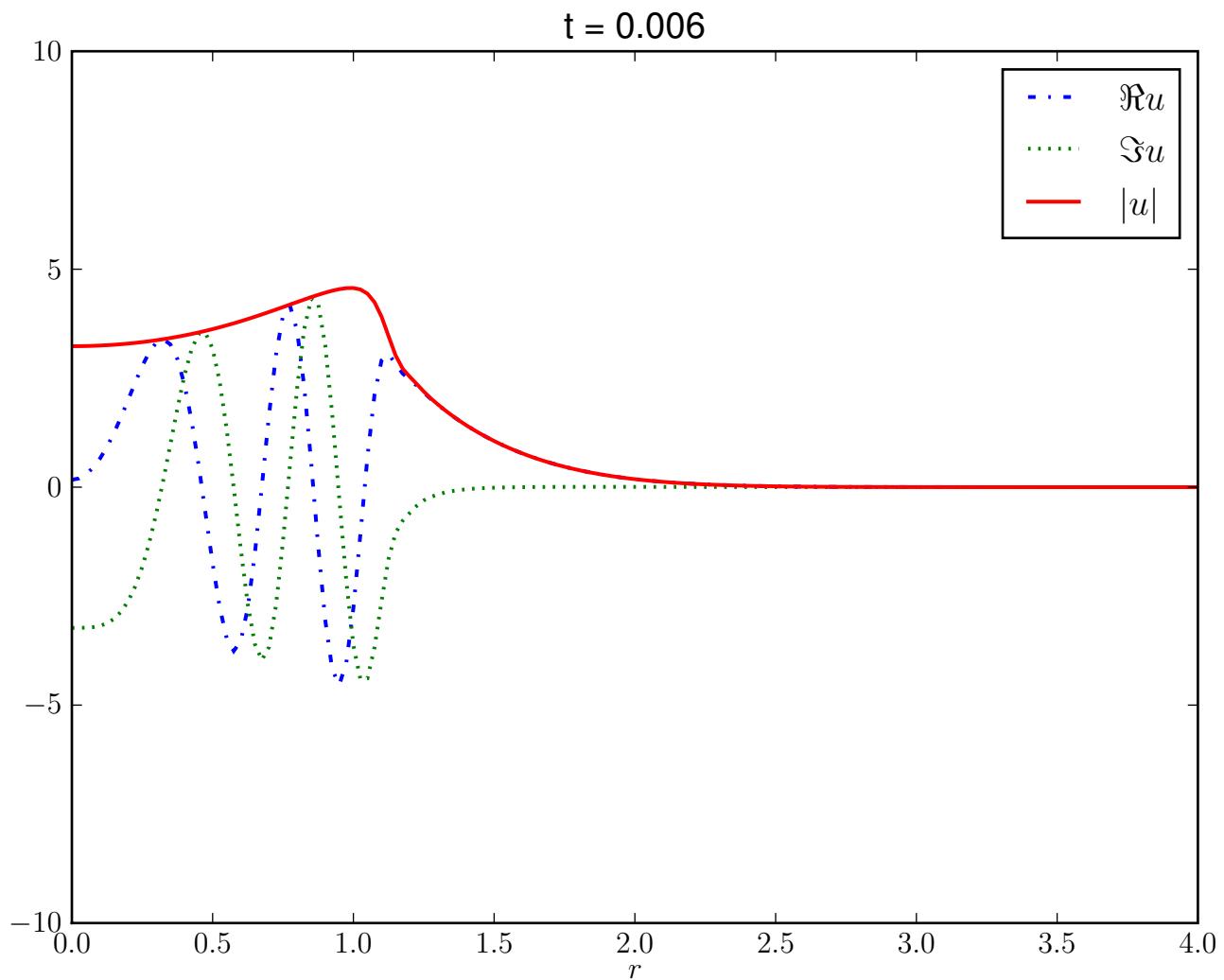
Centered Gaussian Initial Data

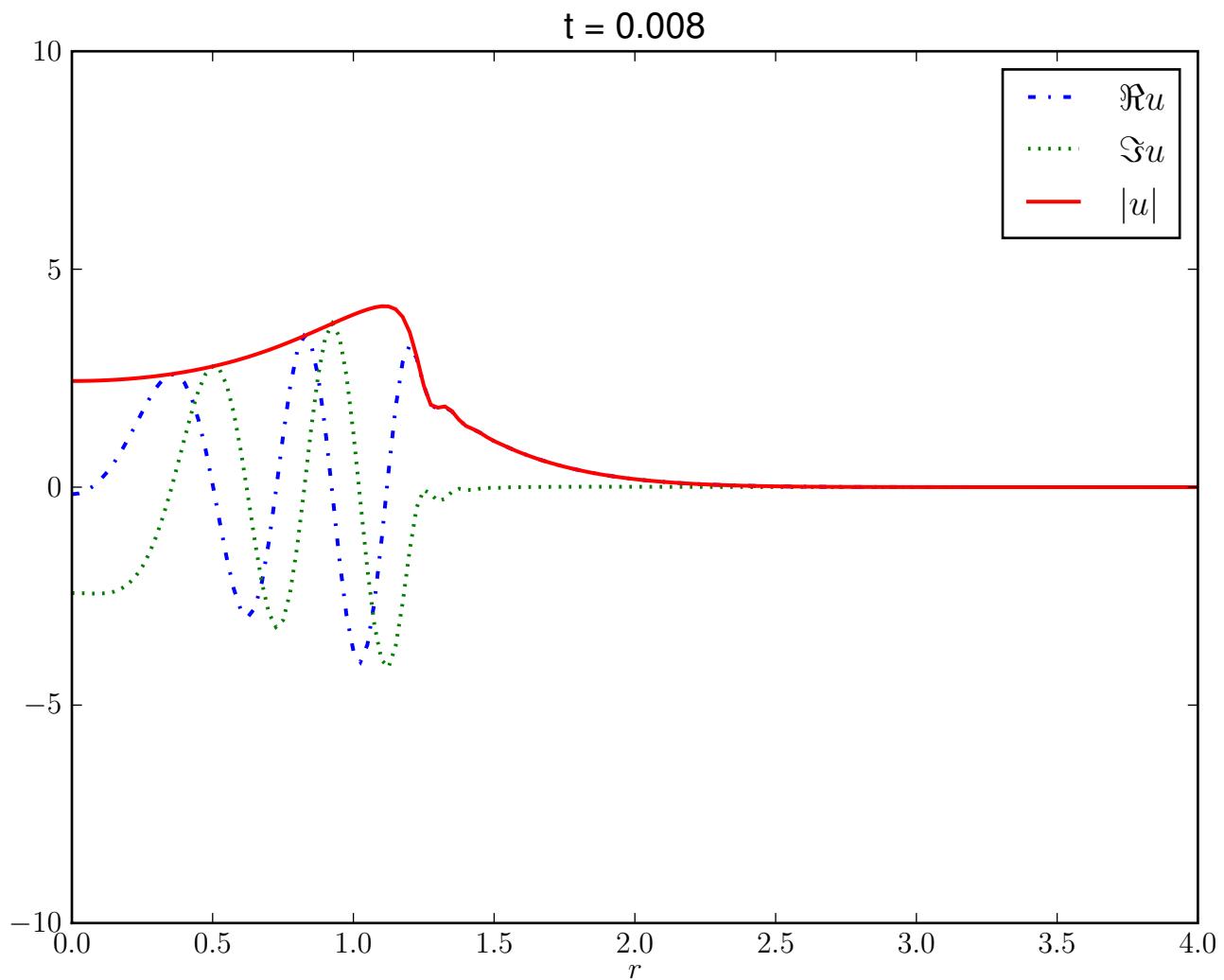
$t = 0$

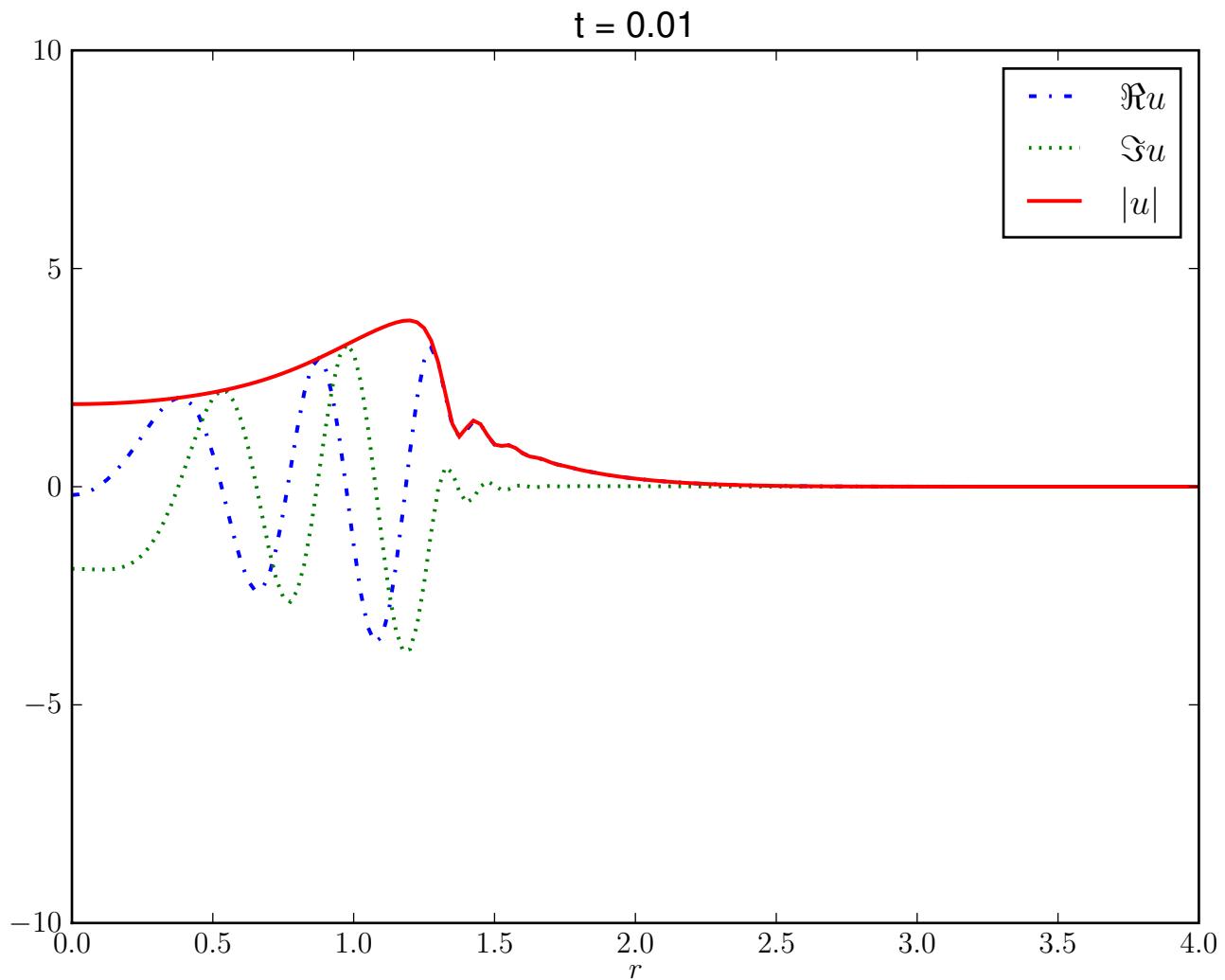


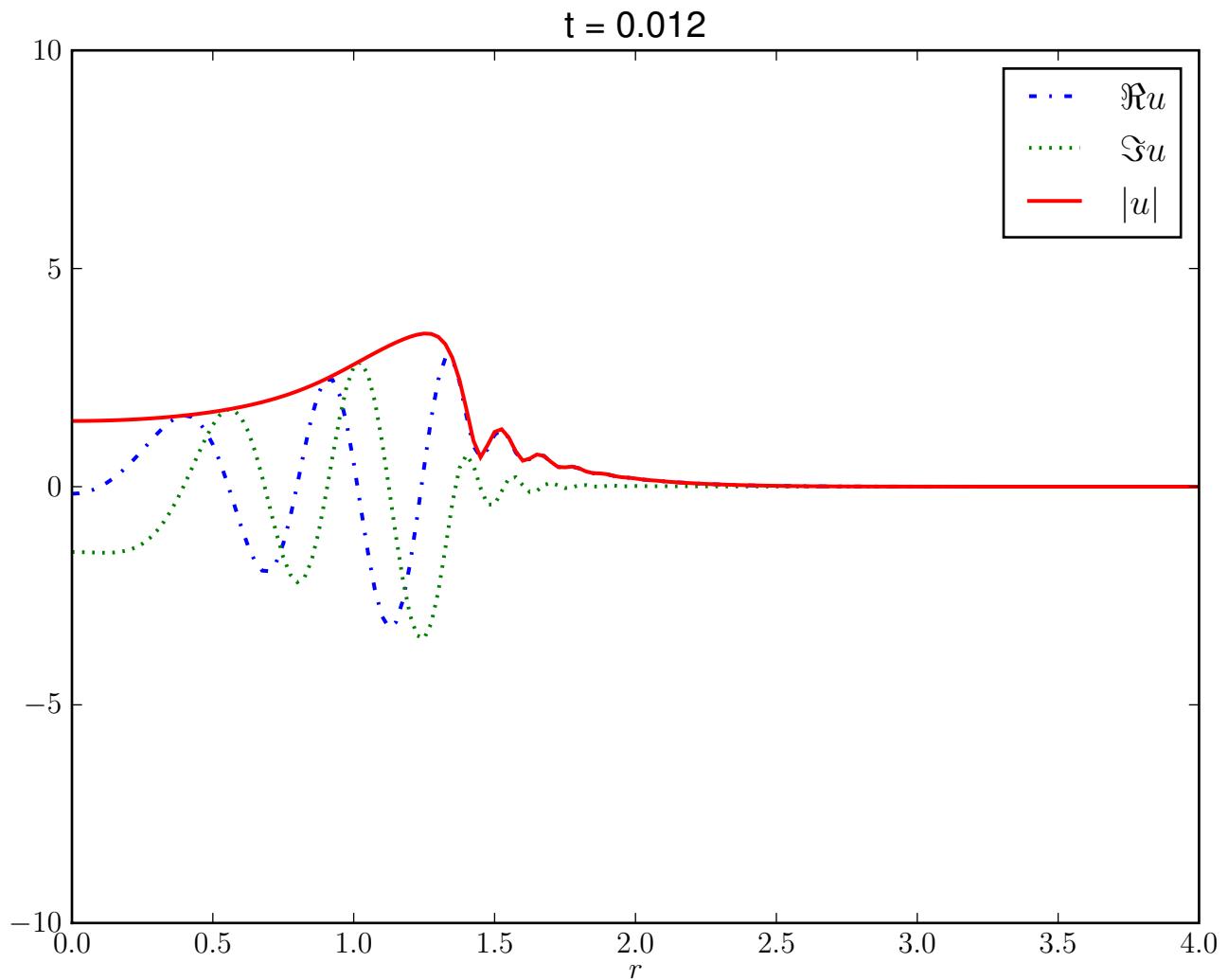


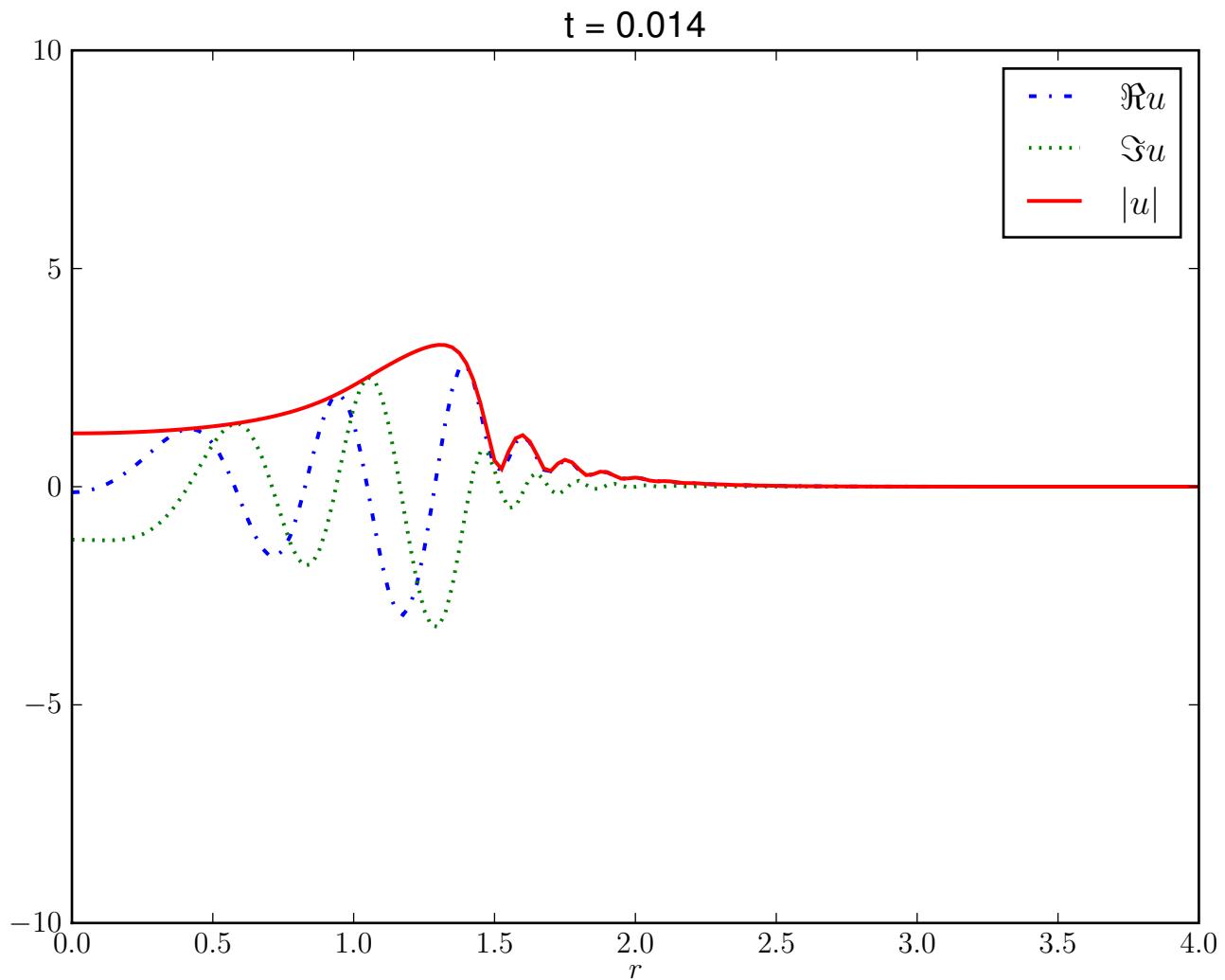


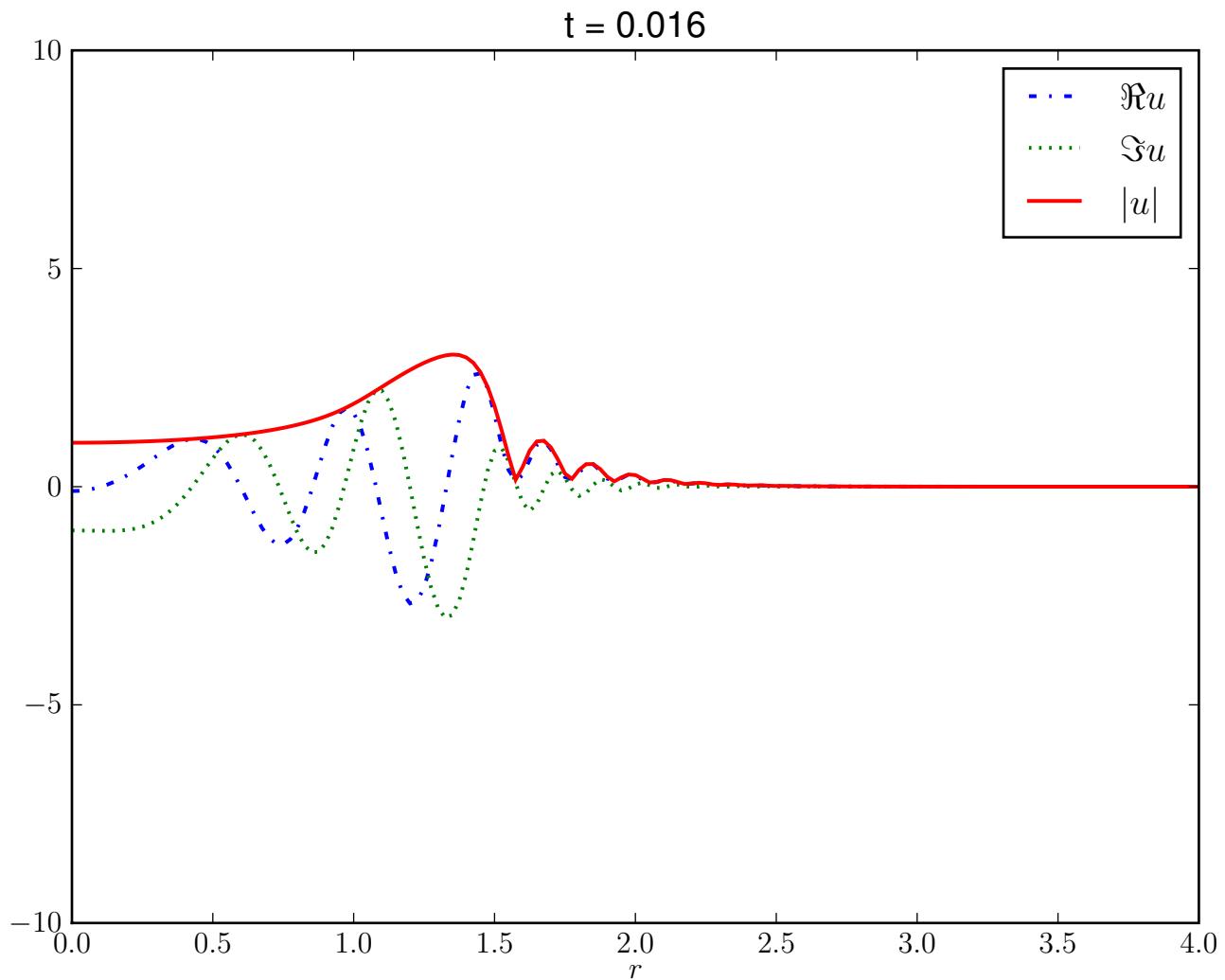


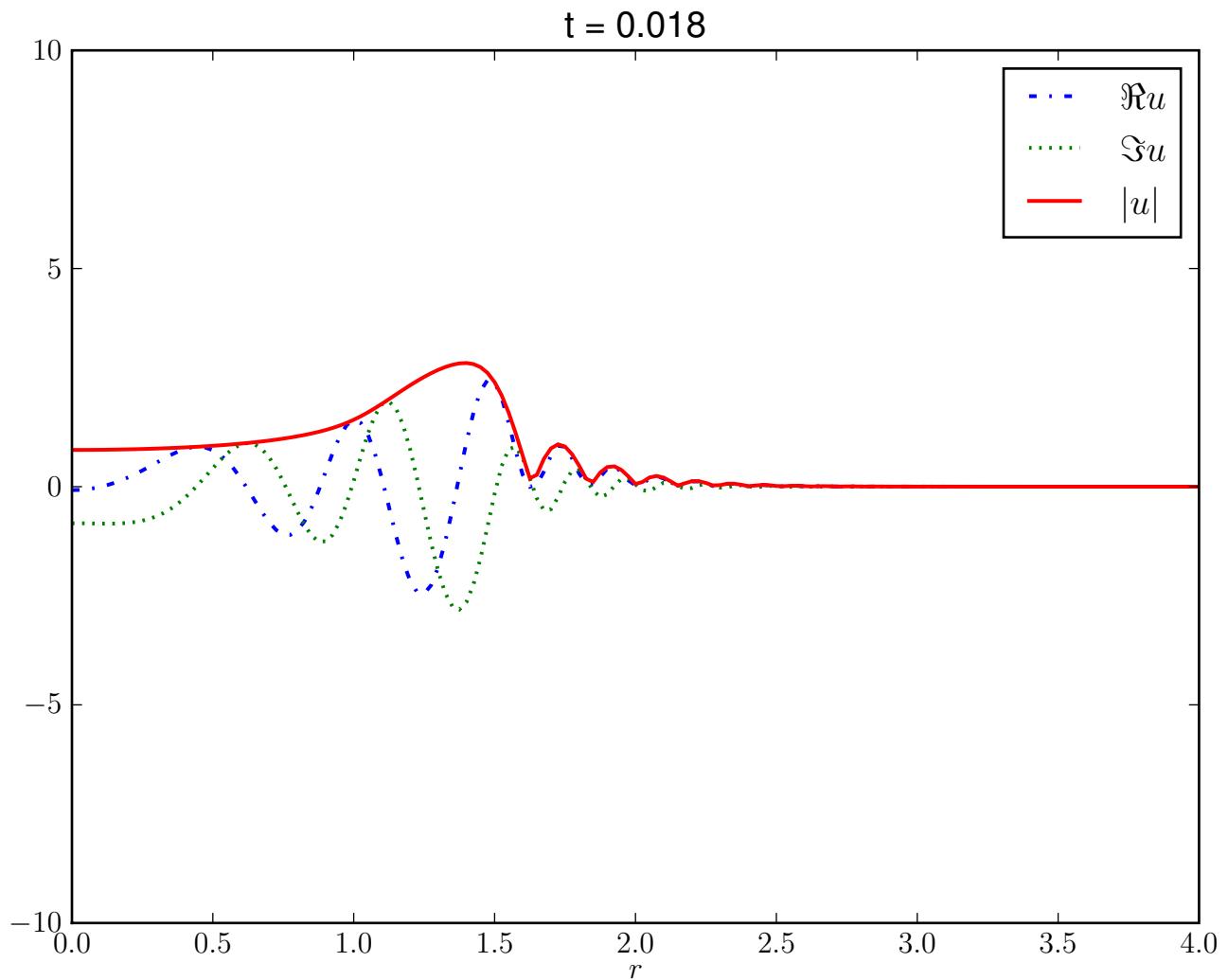


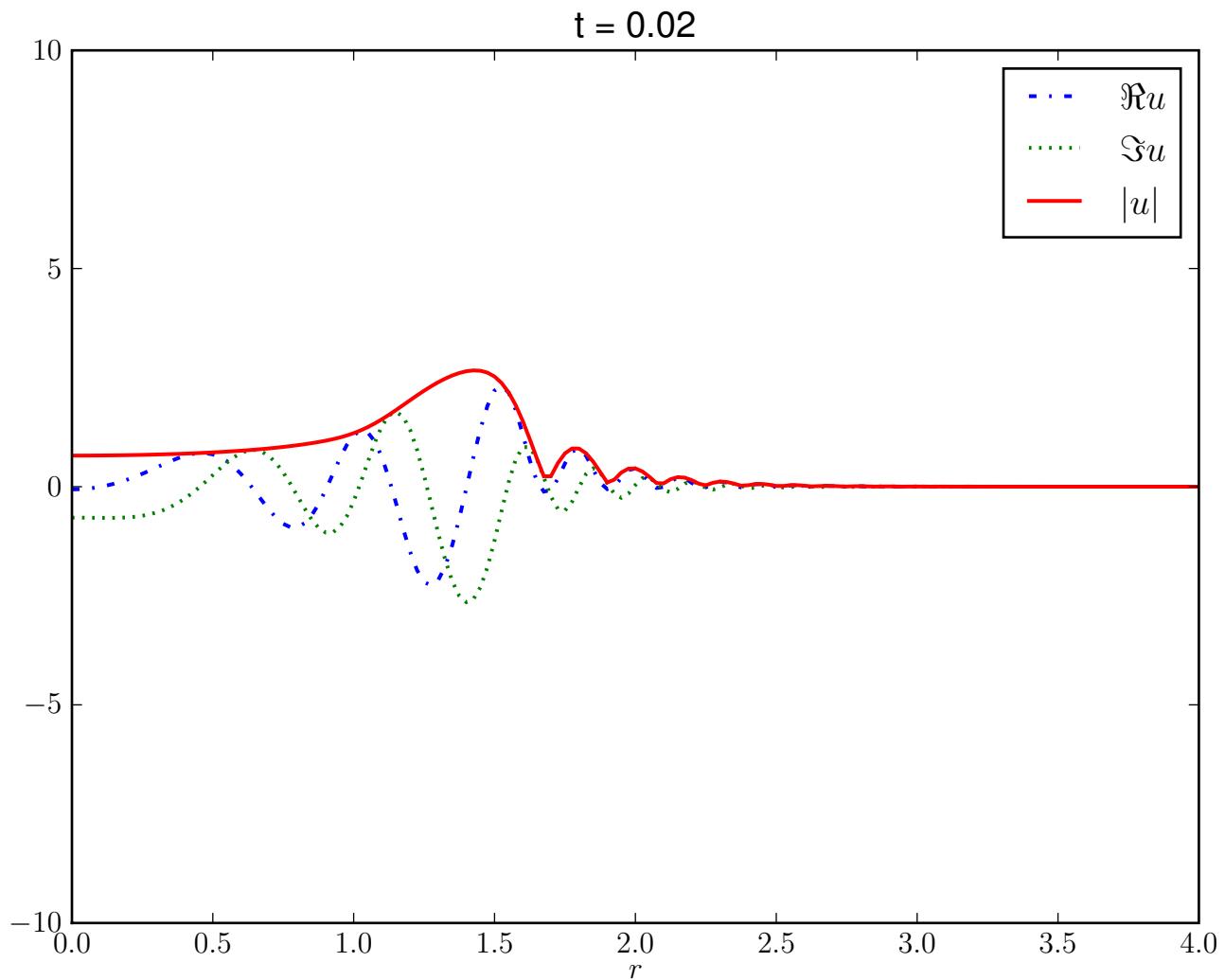


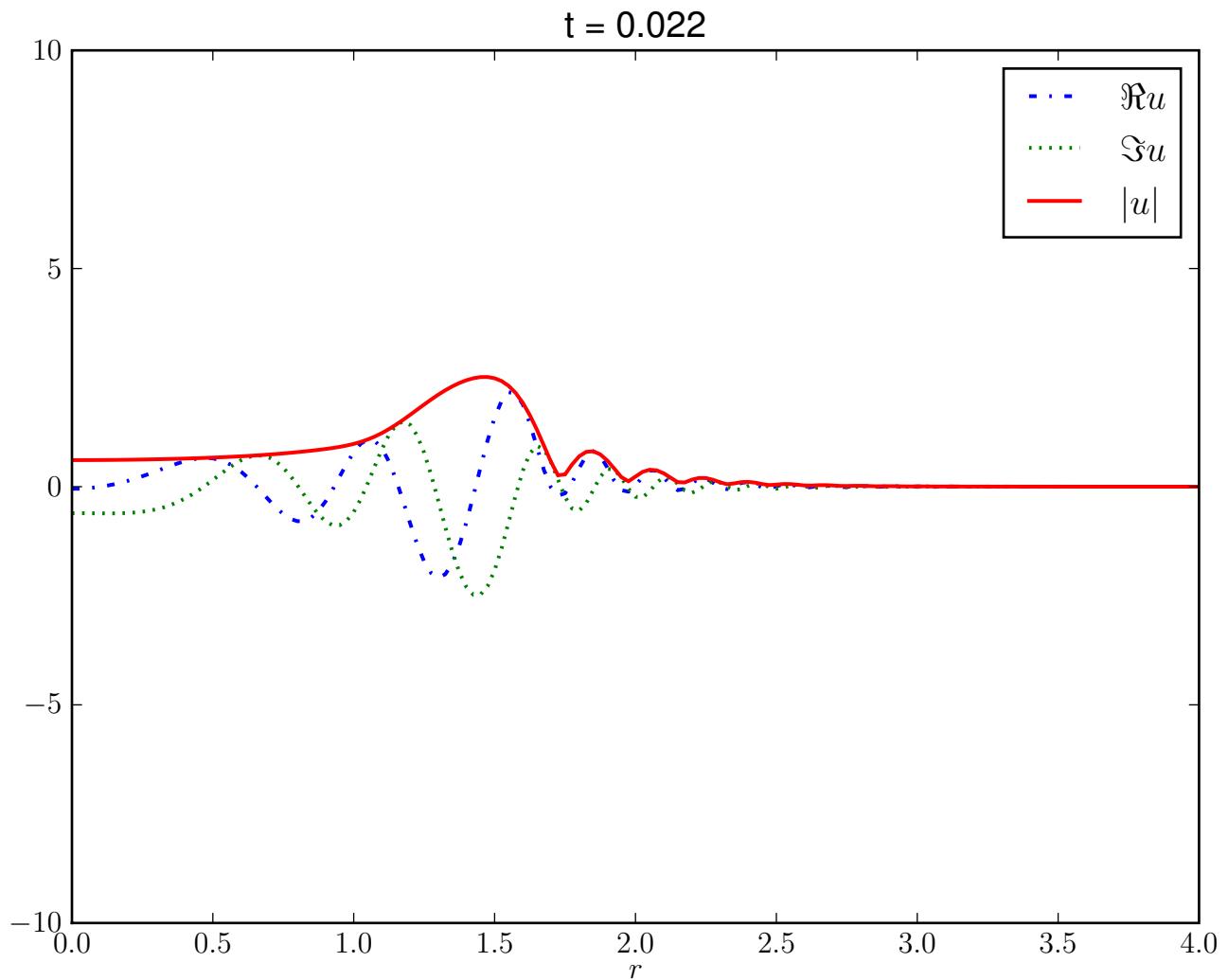


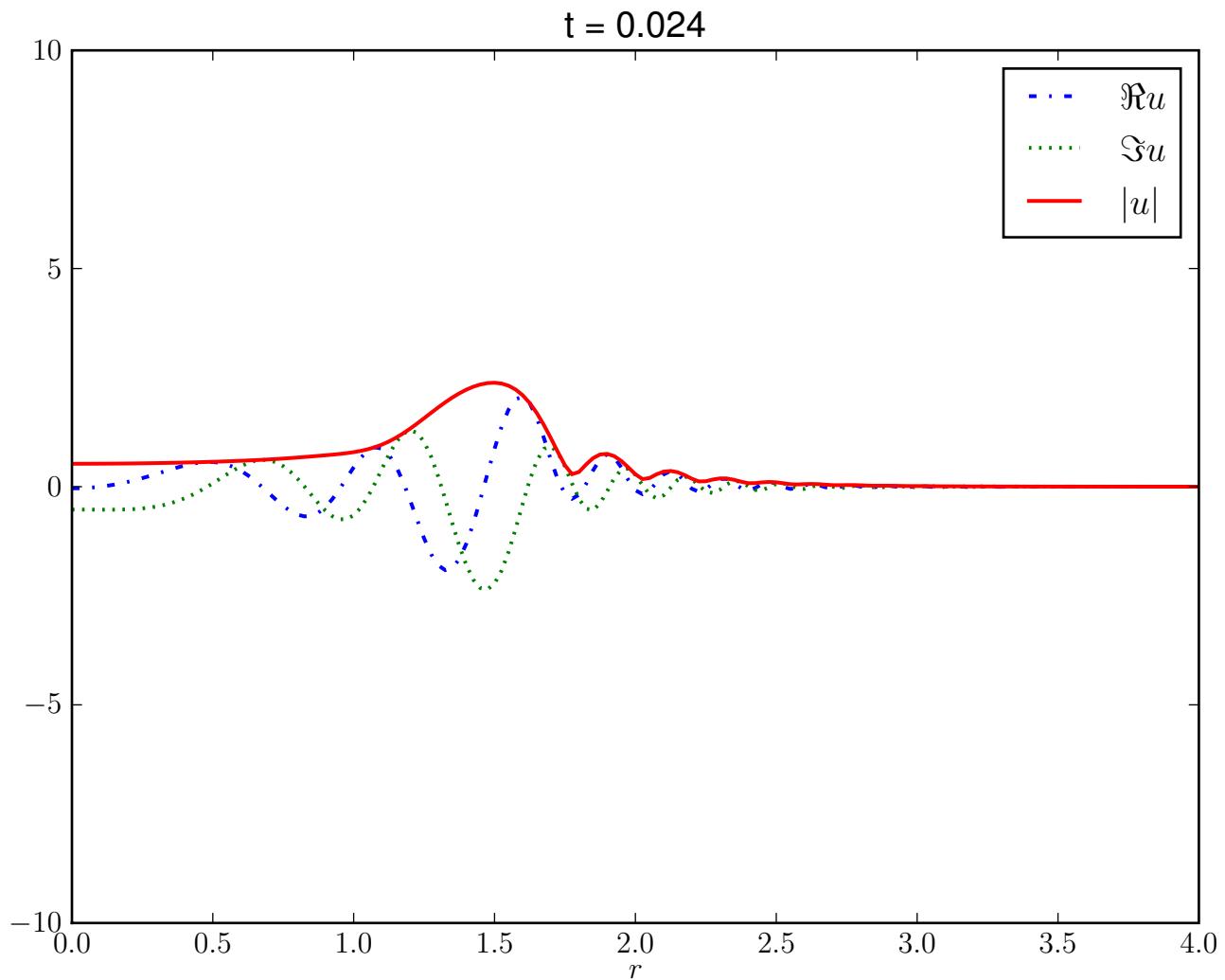


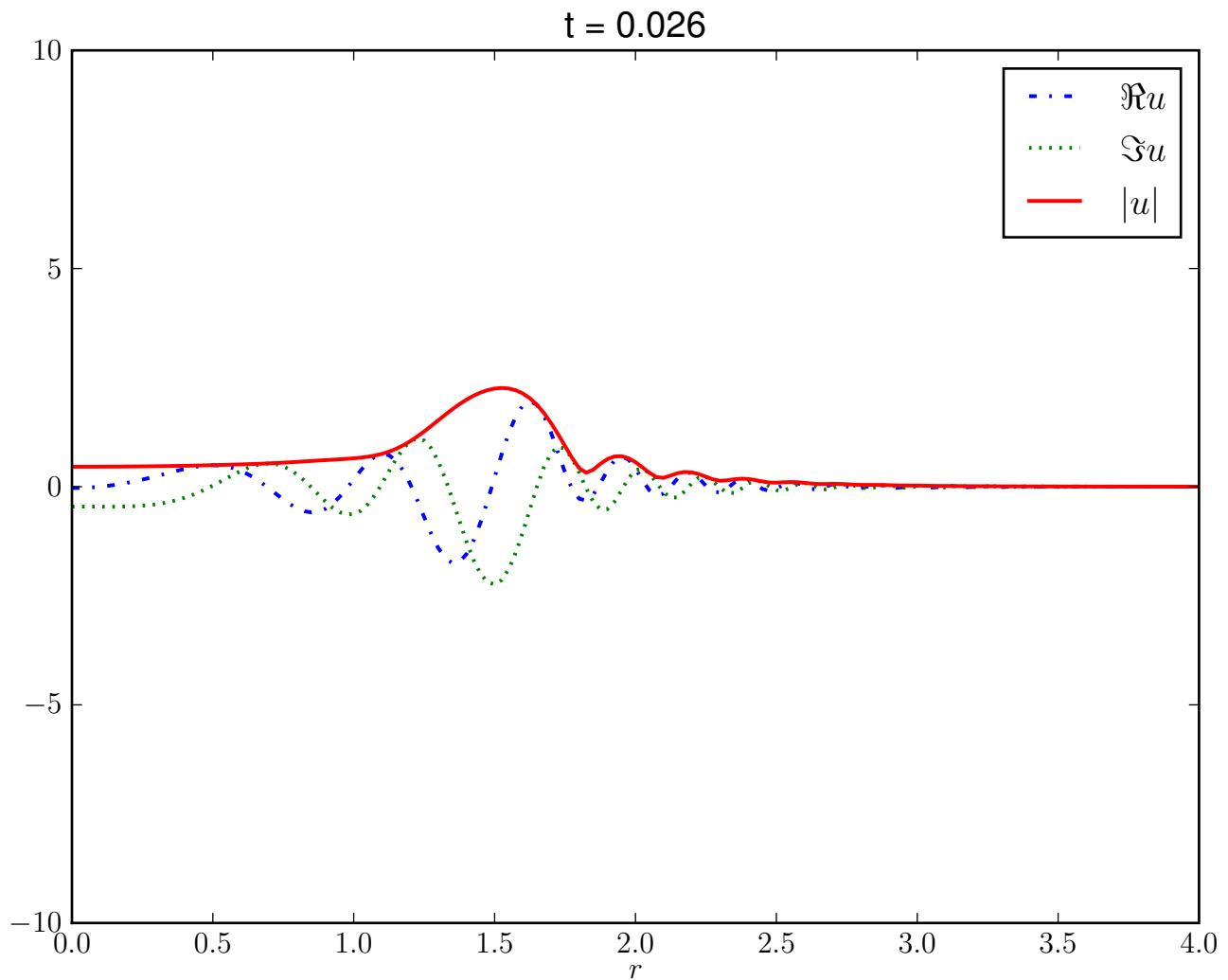


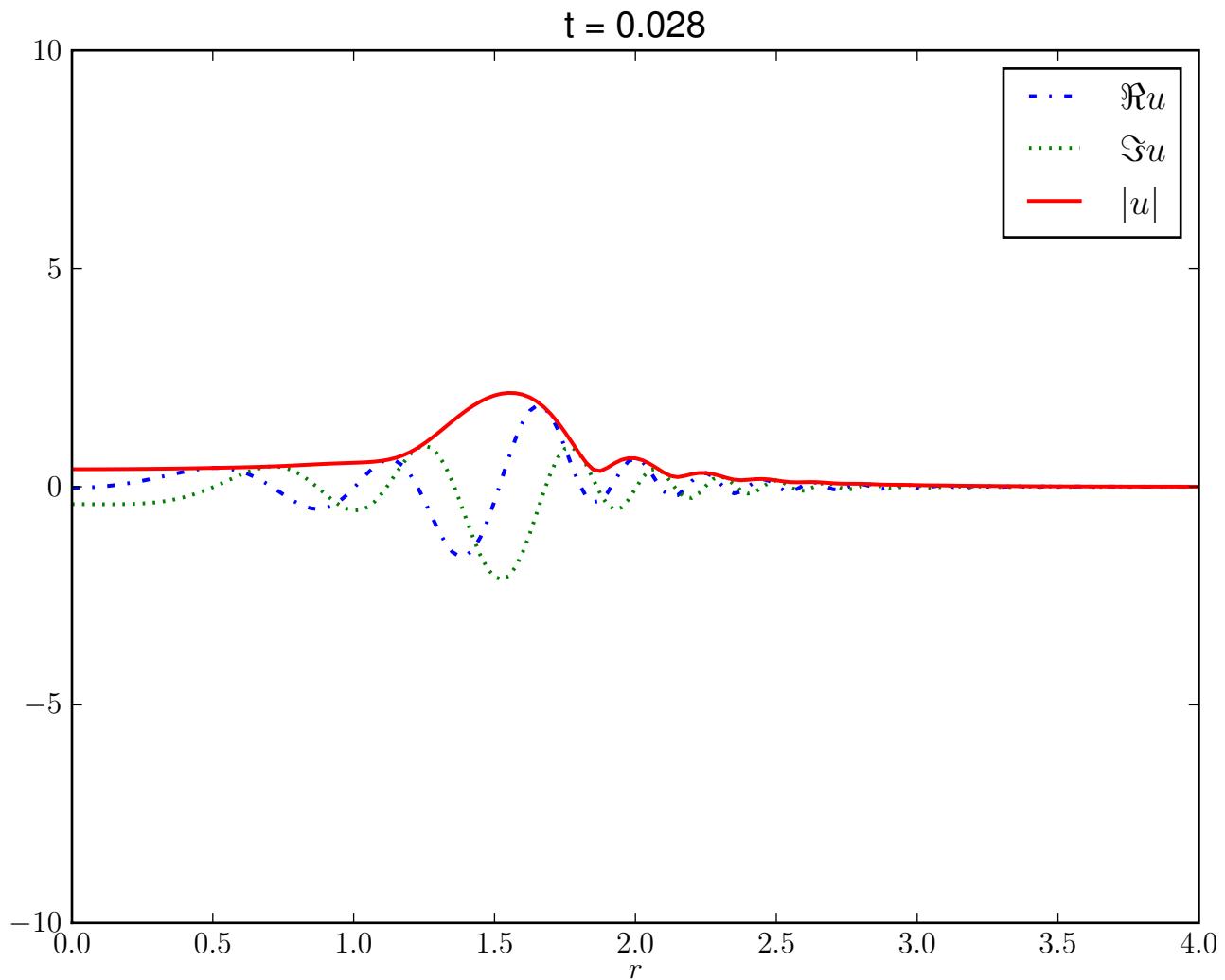


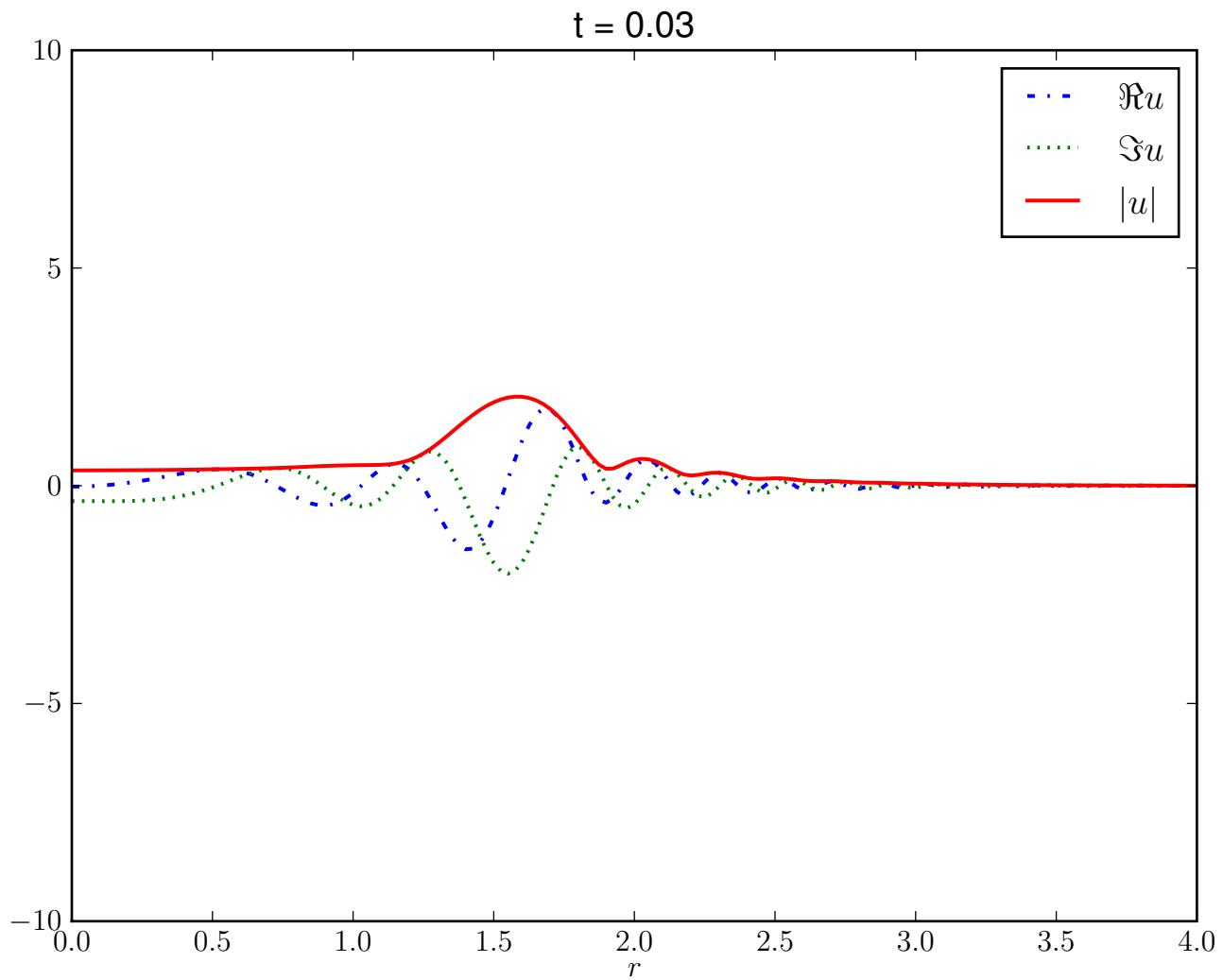


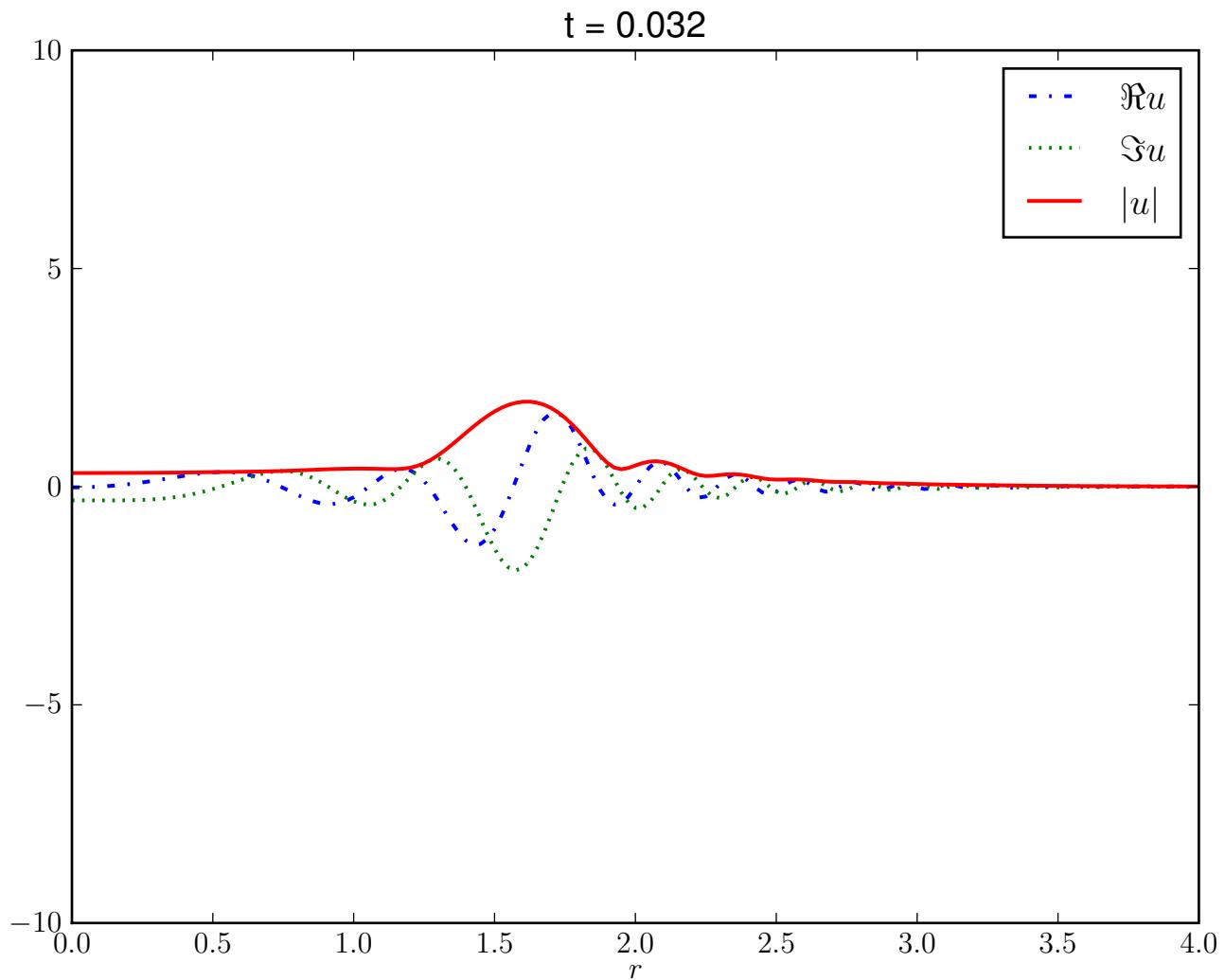


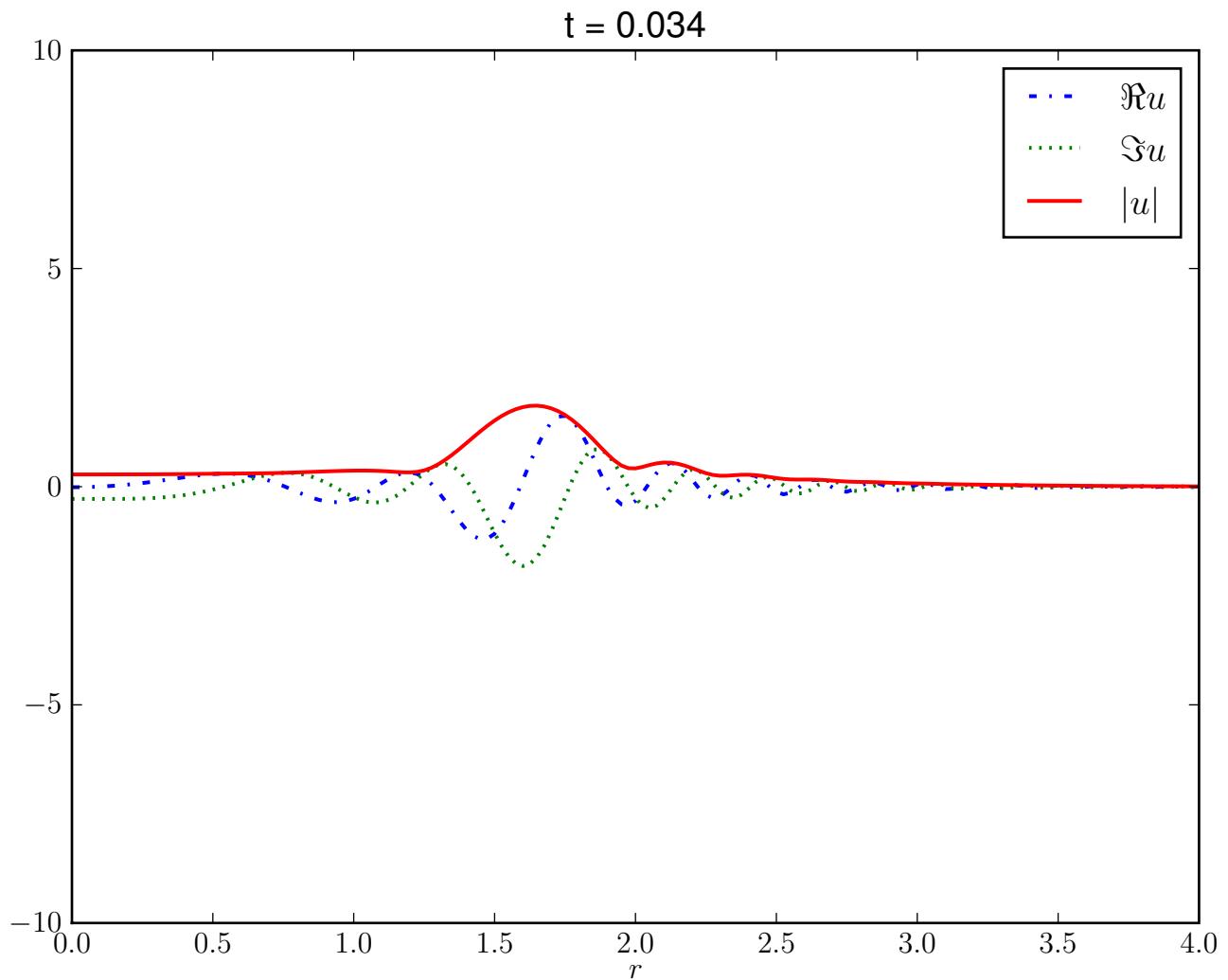


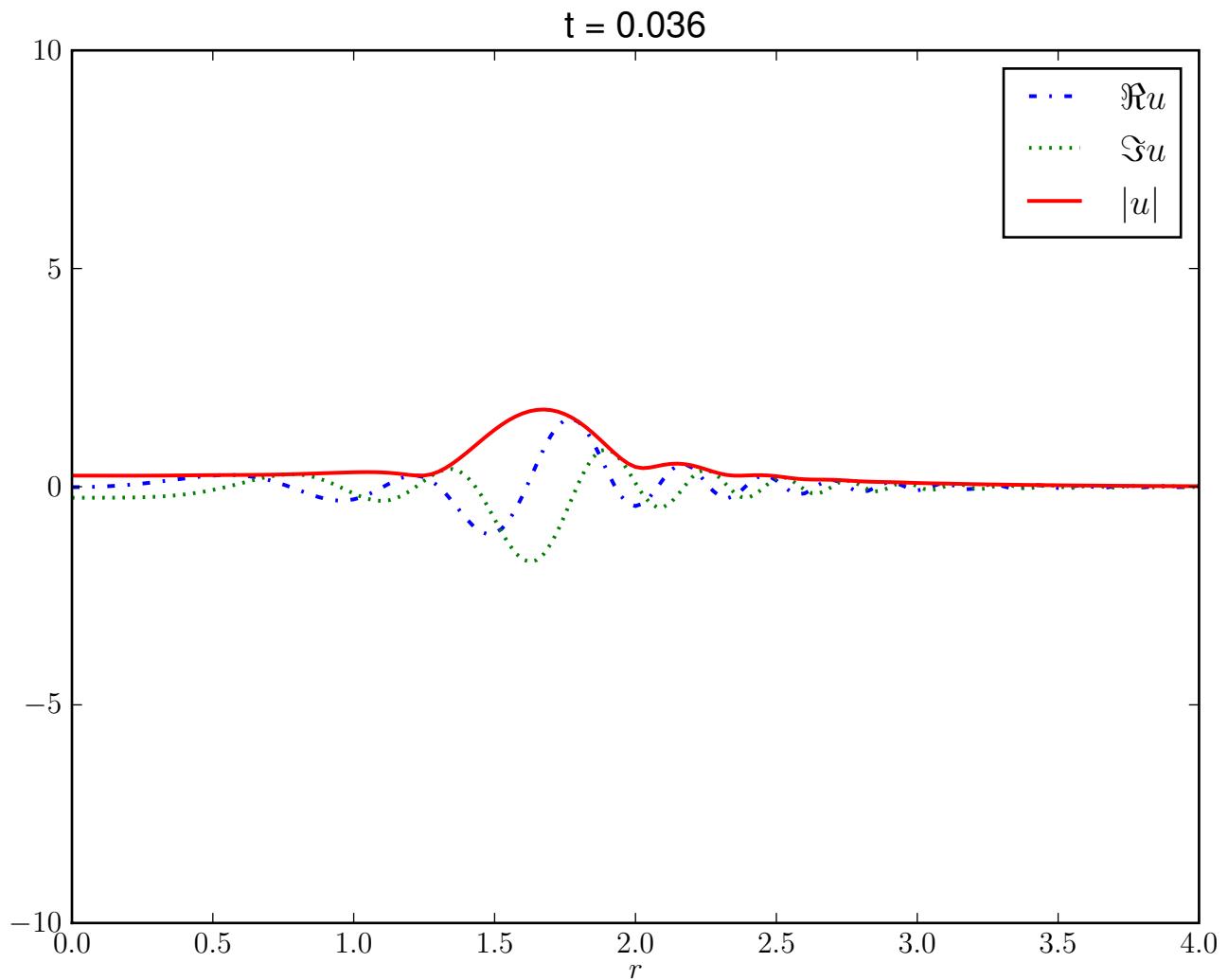


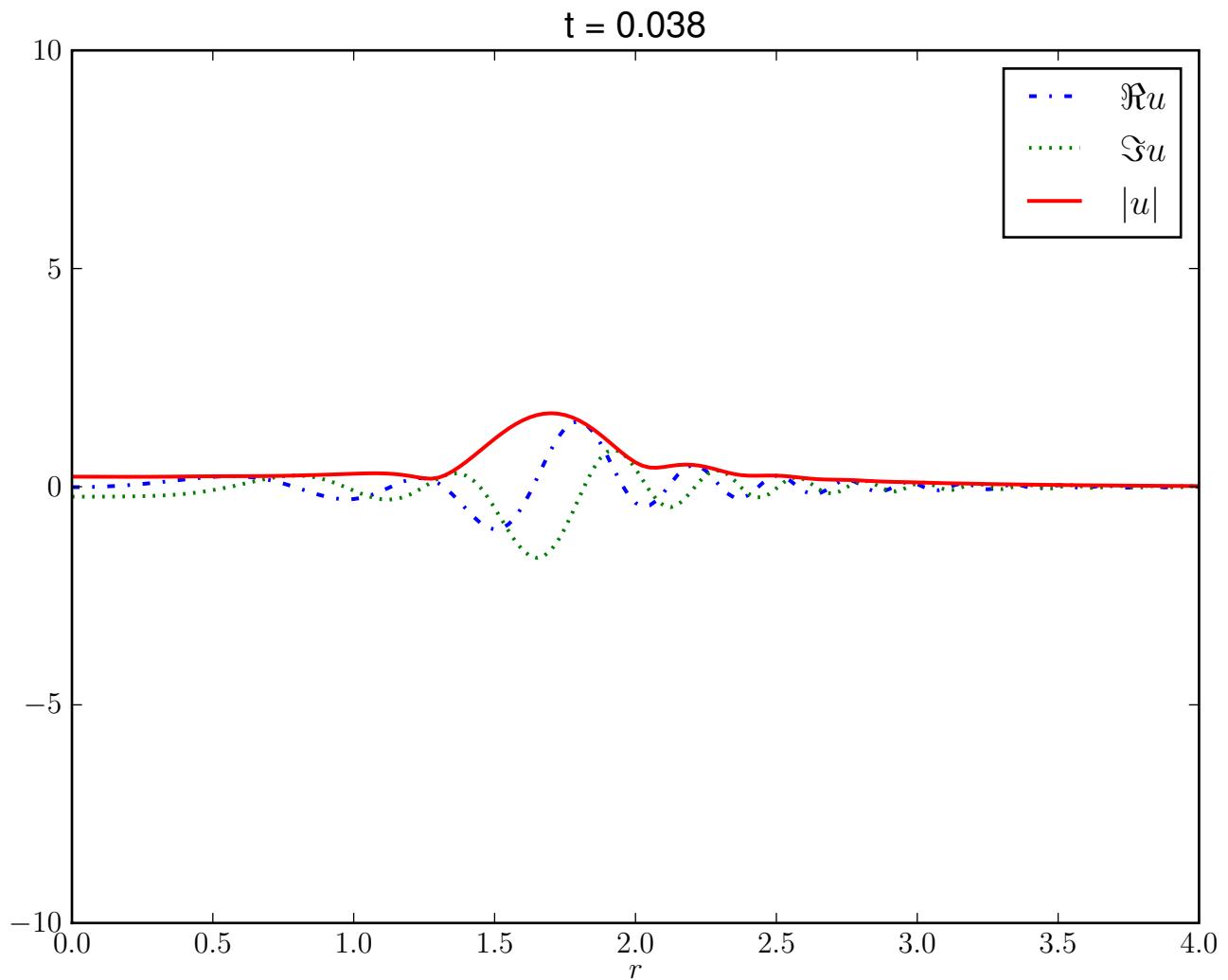


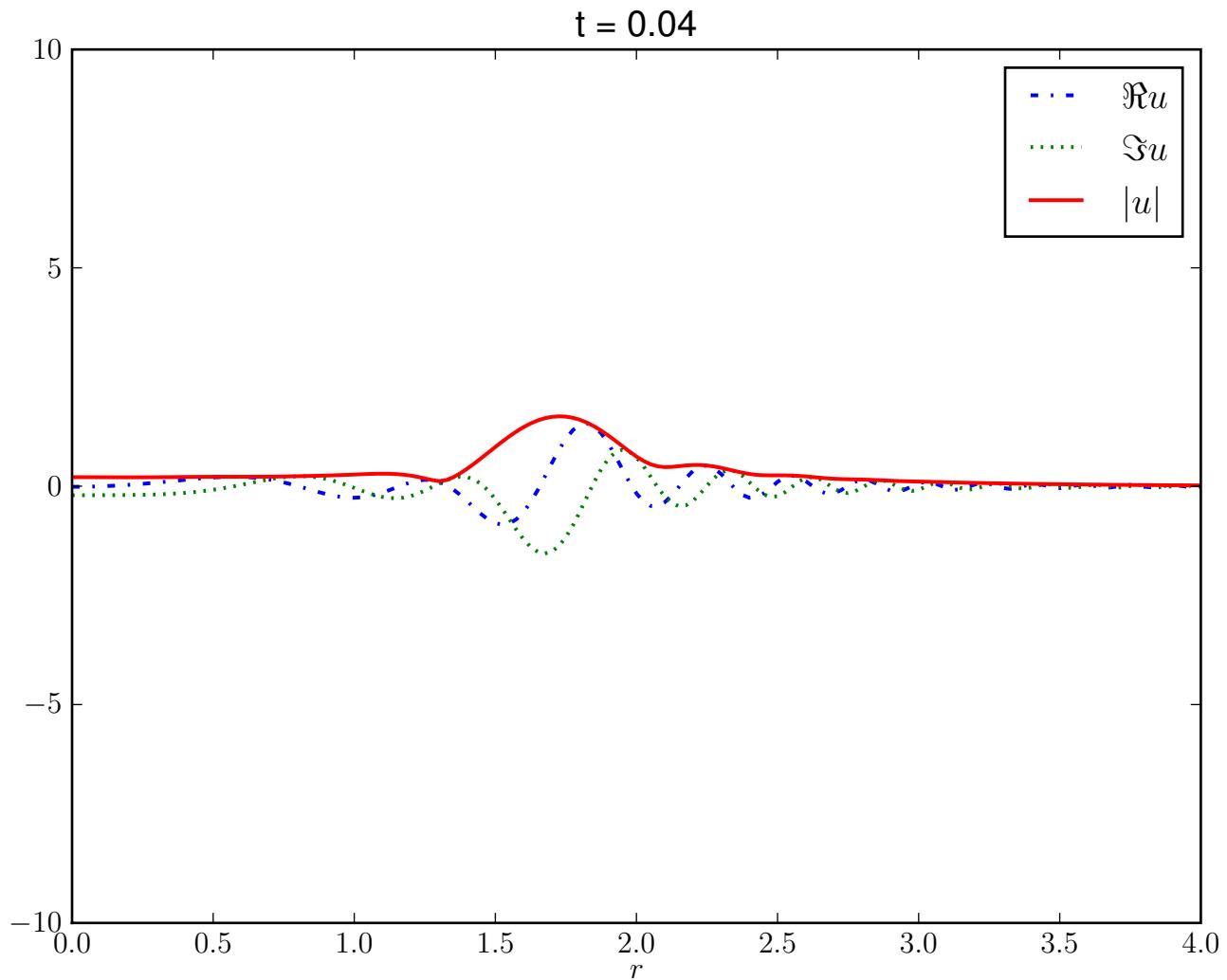




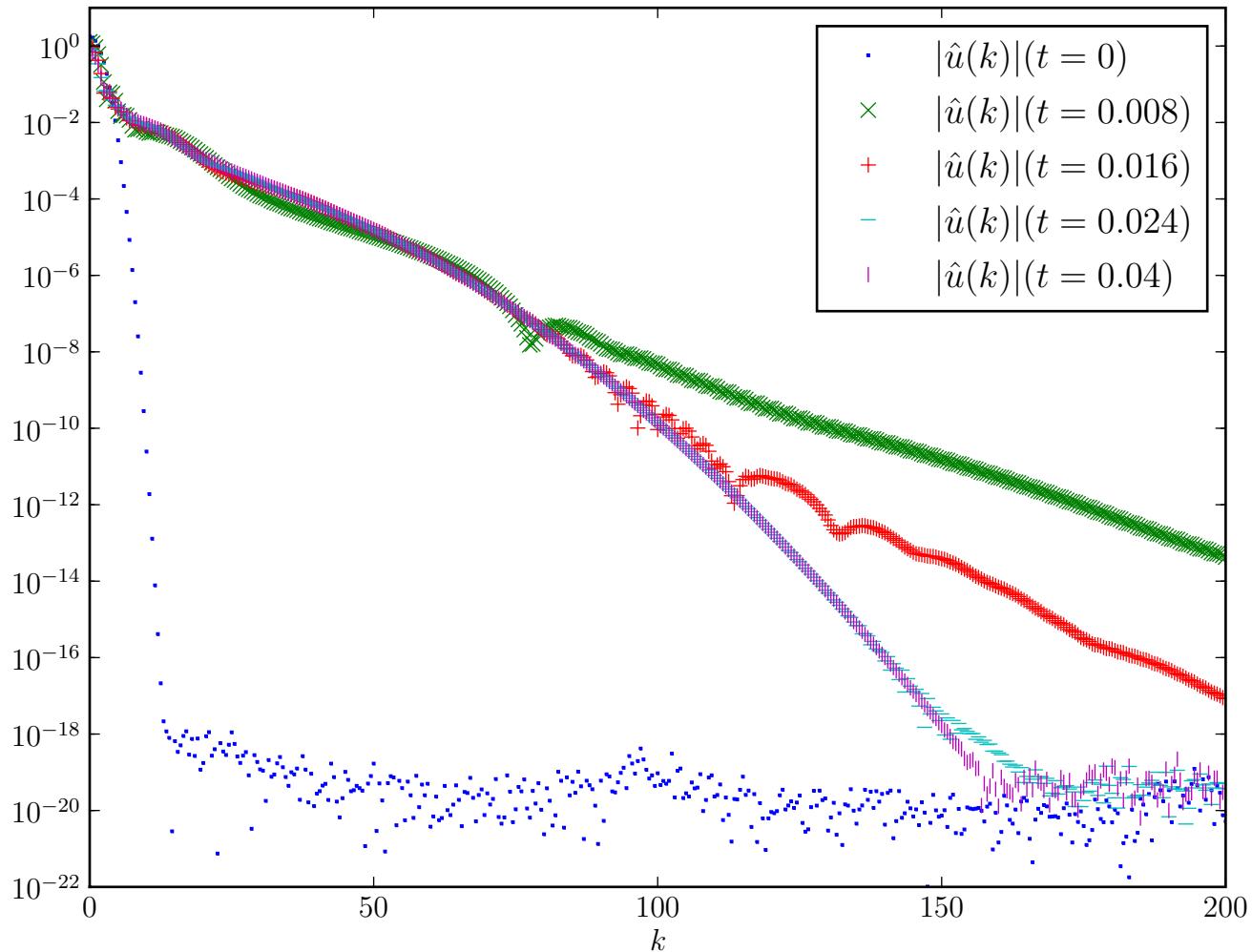




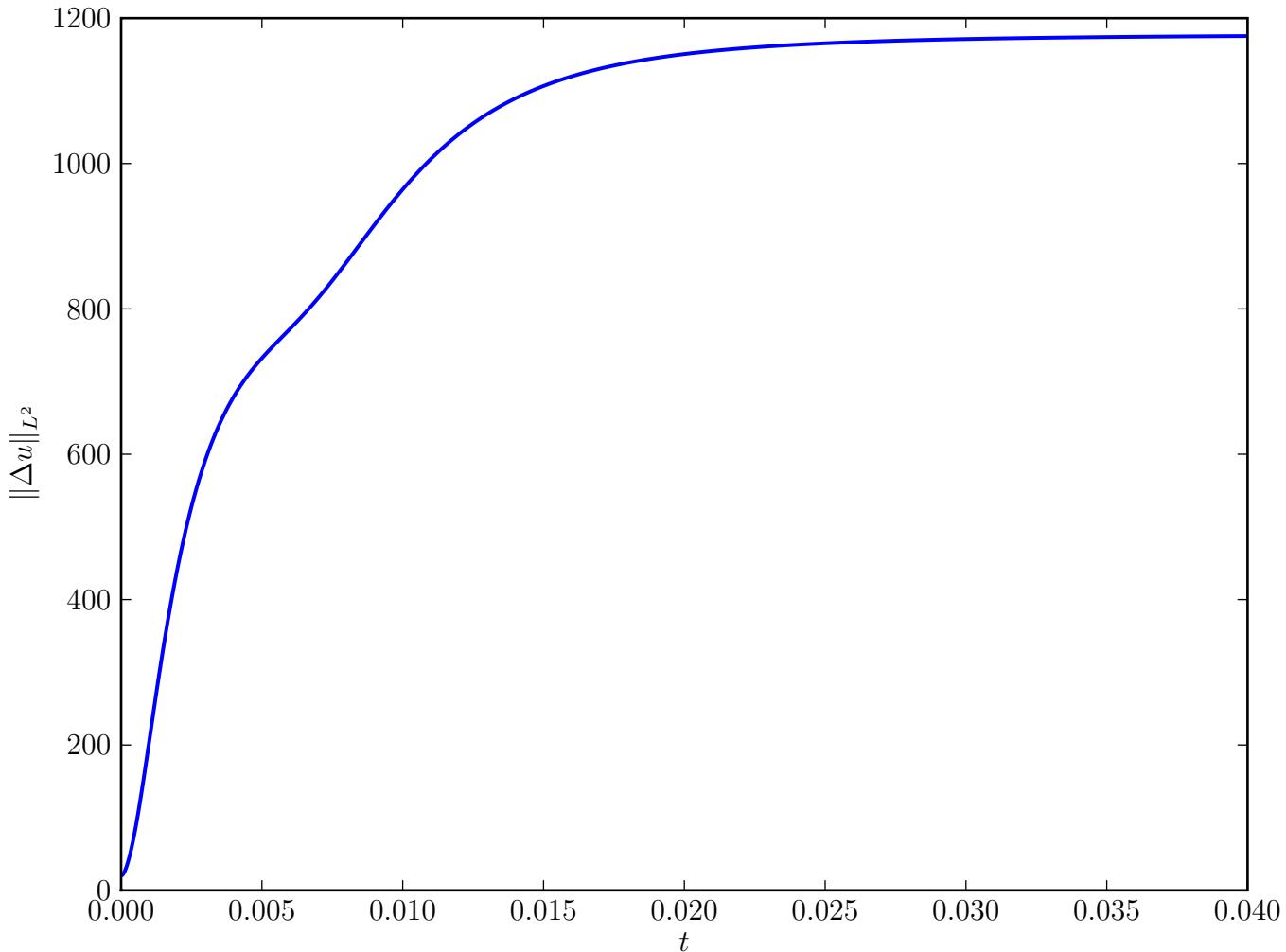




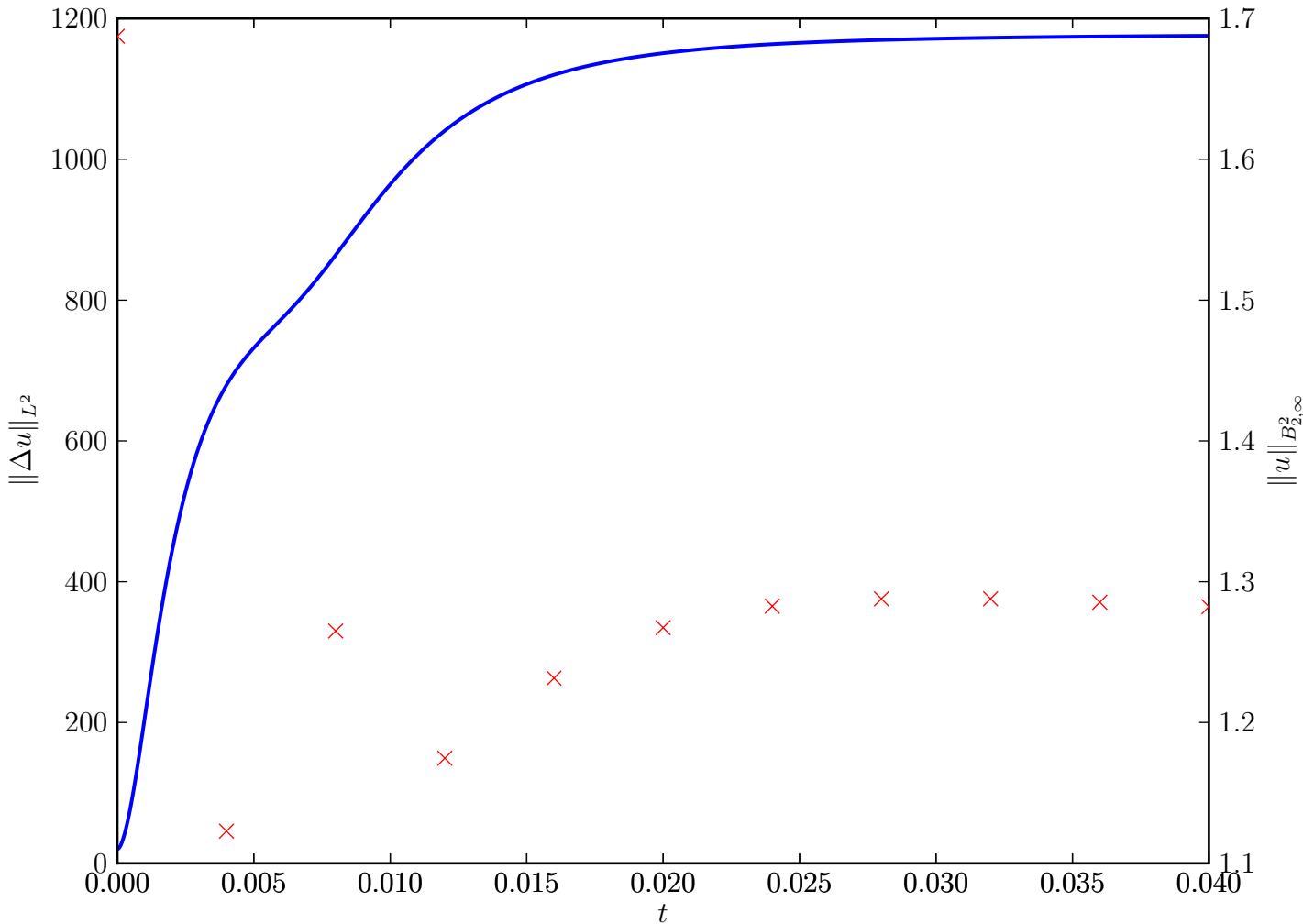
Centered Gaussian Fourier transform snapshots along nonlinear flow



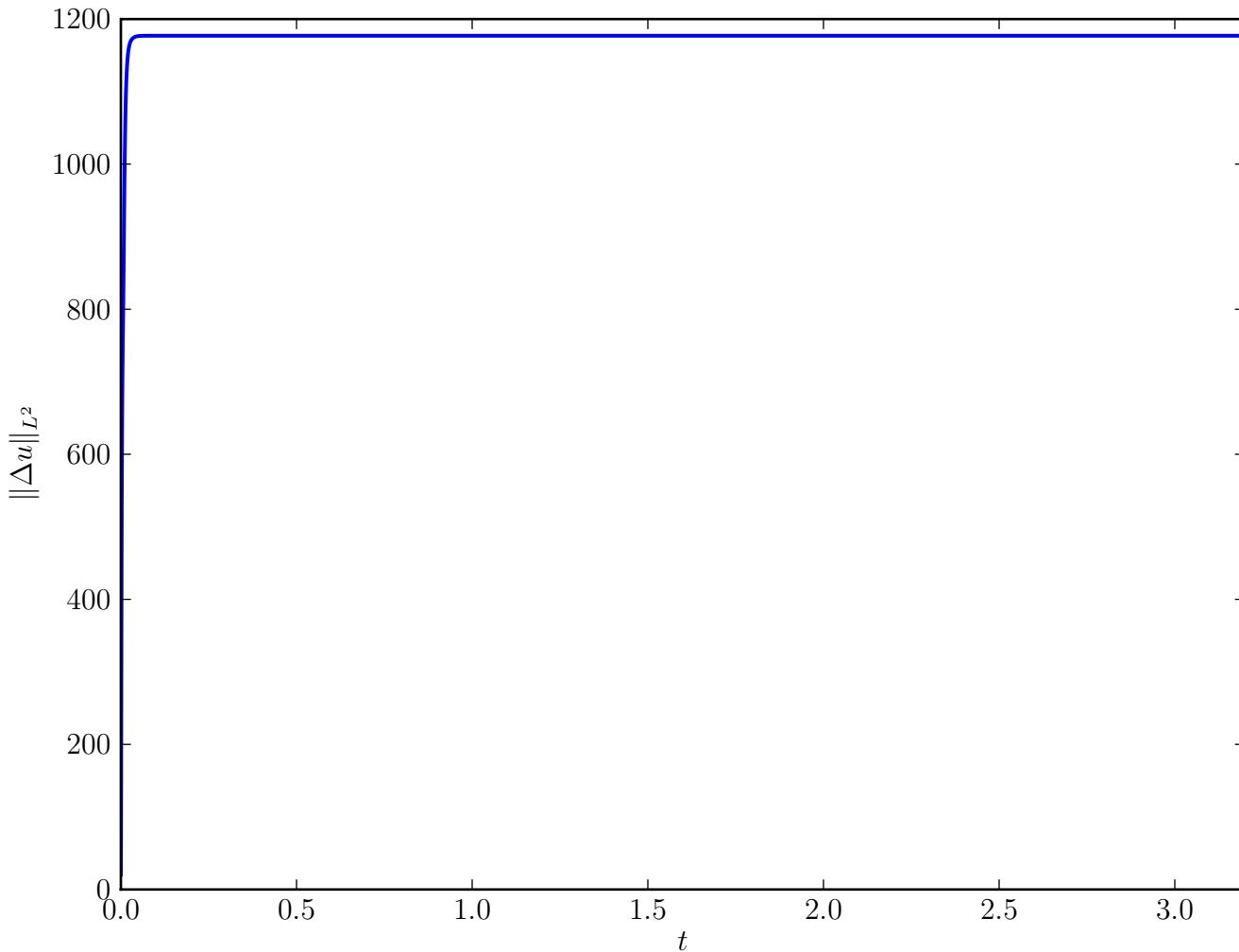
H² norm: Centered Gaussian along nonlinear flow



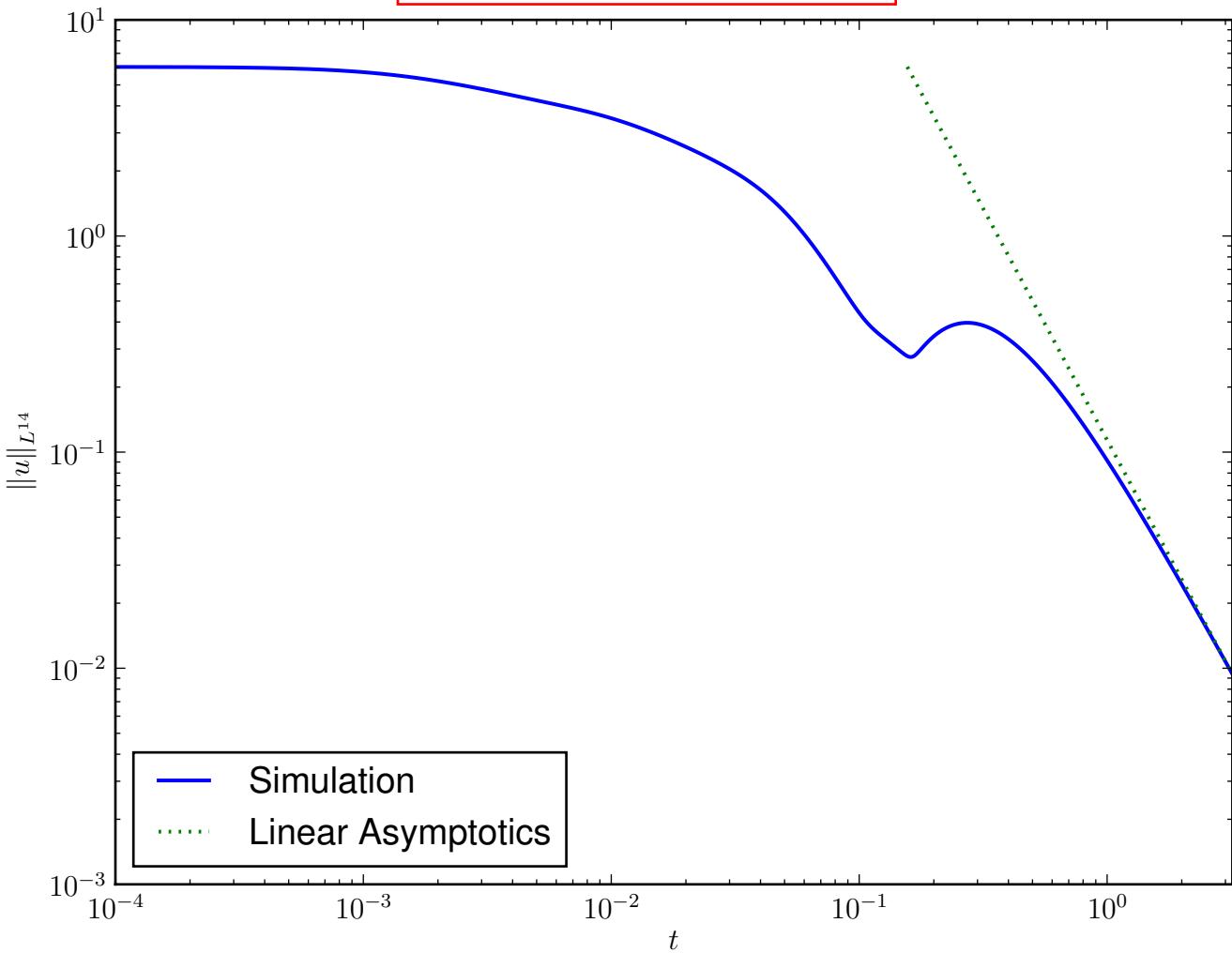
Sobolev vs. Besov: Centered Gaussian along nonlinear flow



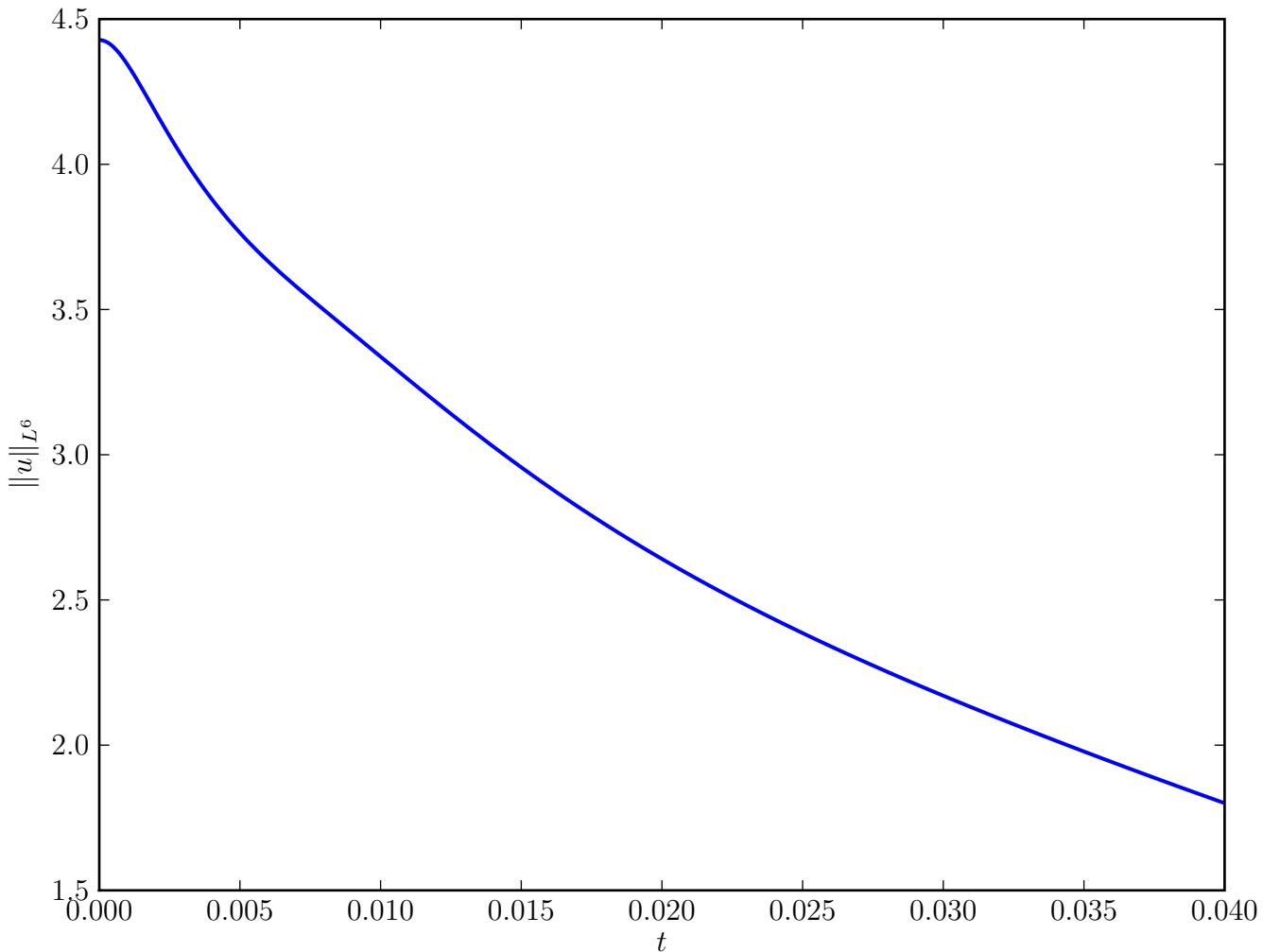
Longer time H² norm: Centered Gaussian along nonlinear flow



Strichartz L¹⁴_x decay asymptotics



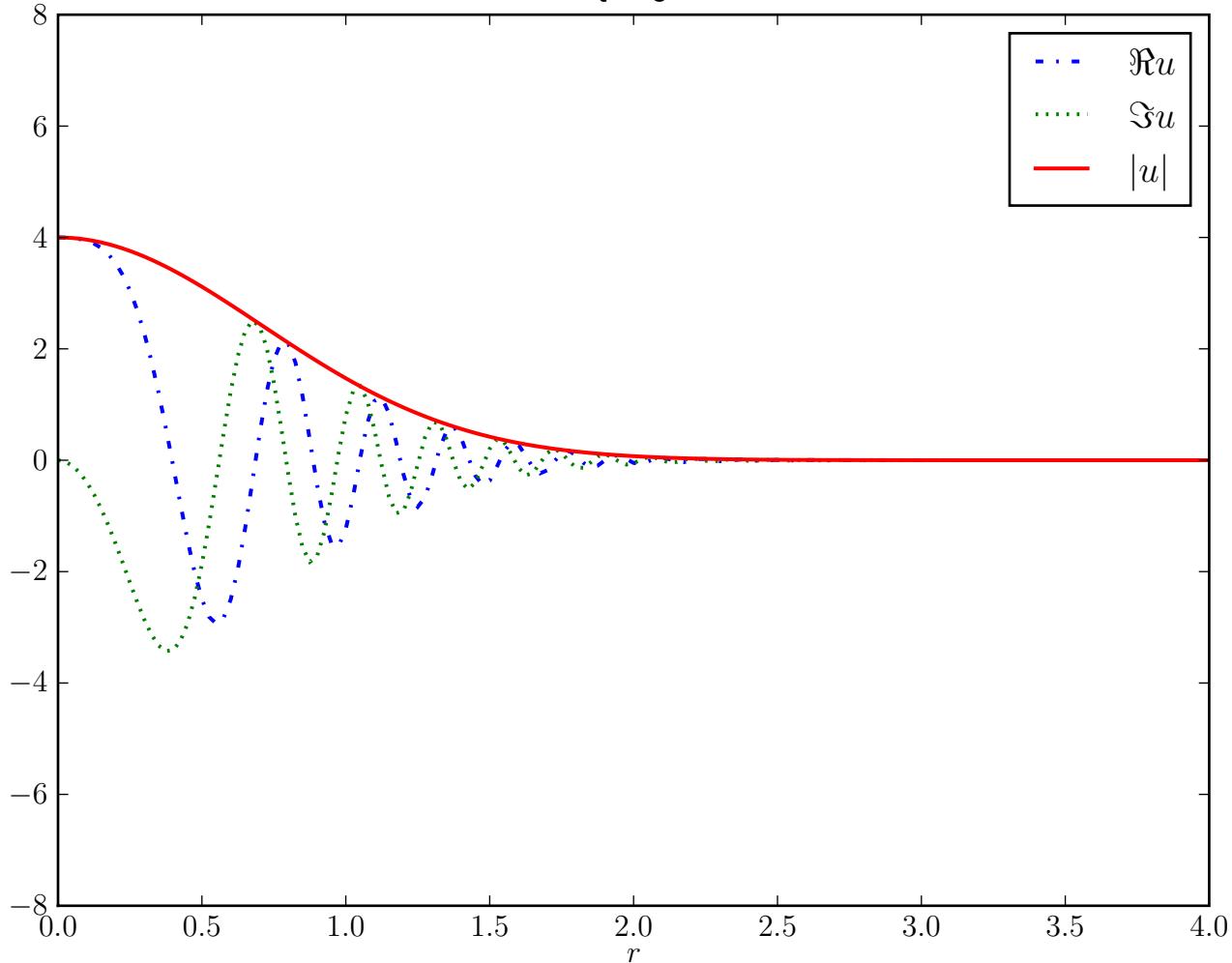
Potential energy norm decay: Centered Gaussian along nonlinear flow

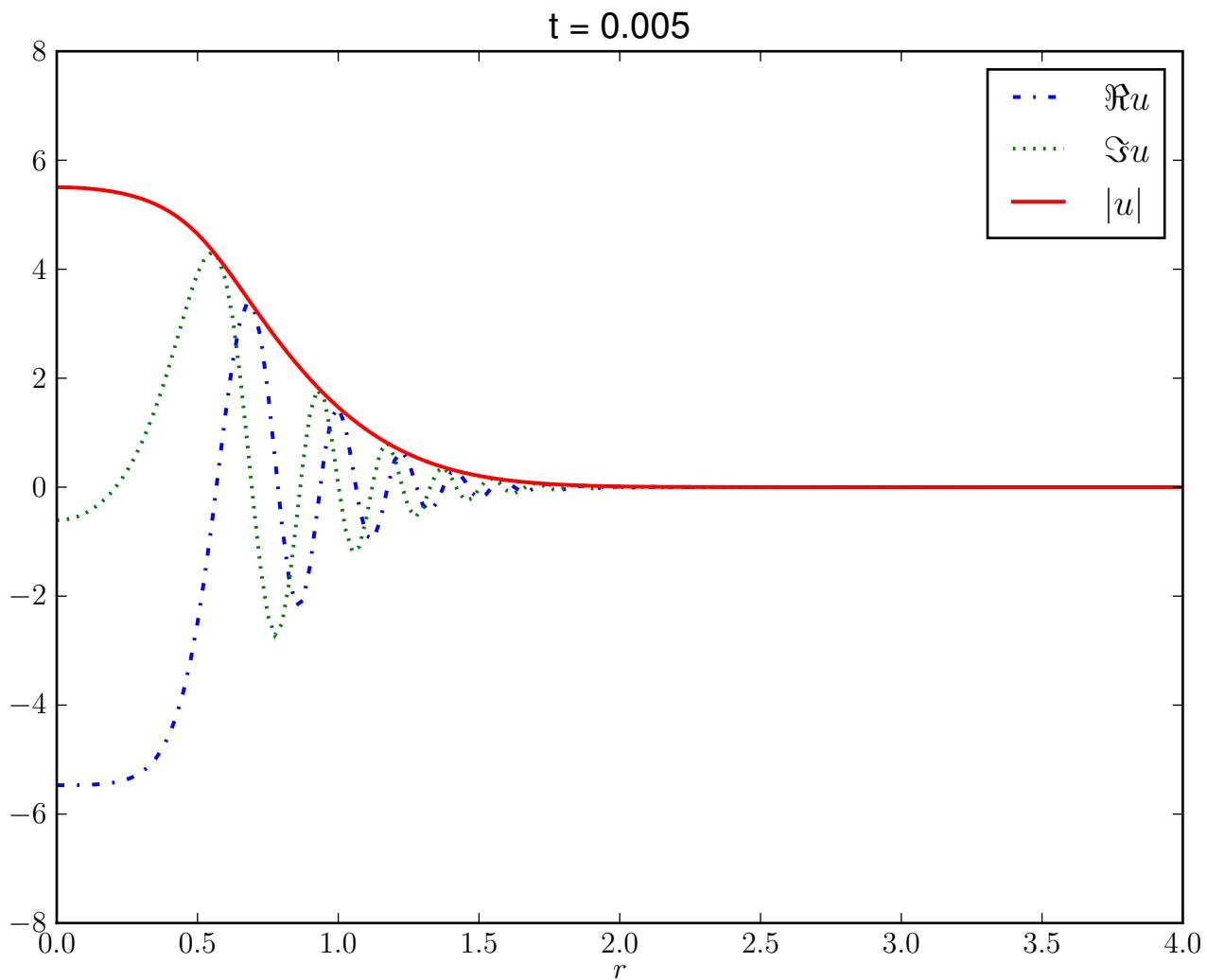


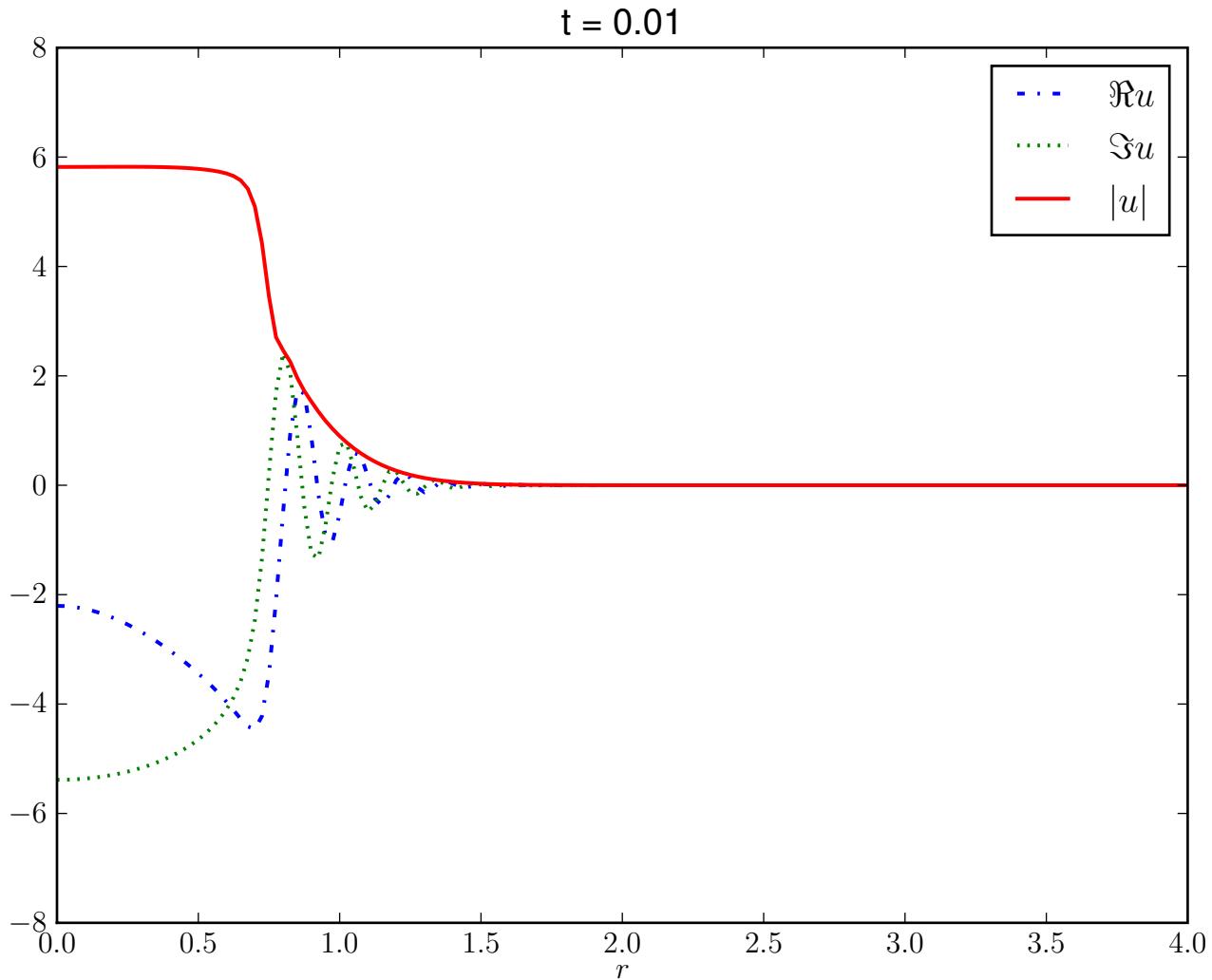
Phased Centered Gaussian Initial Data

Nonlinear Flow

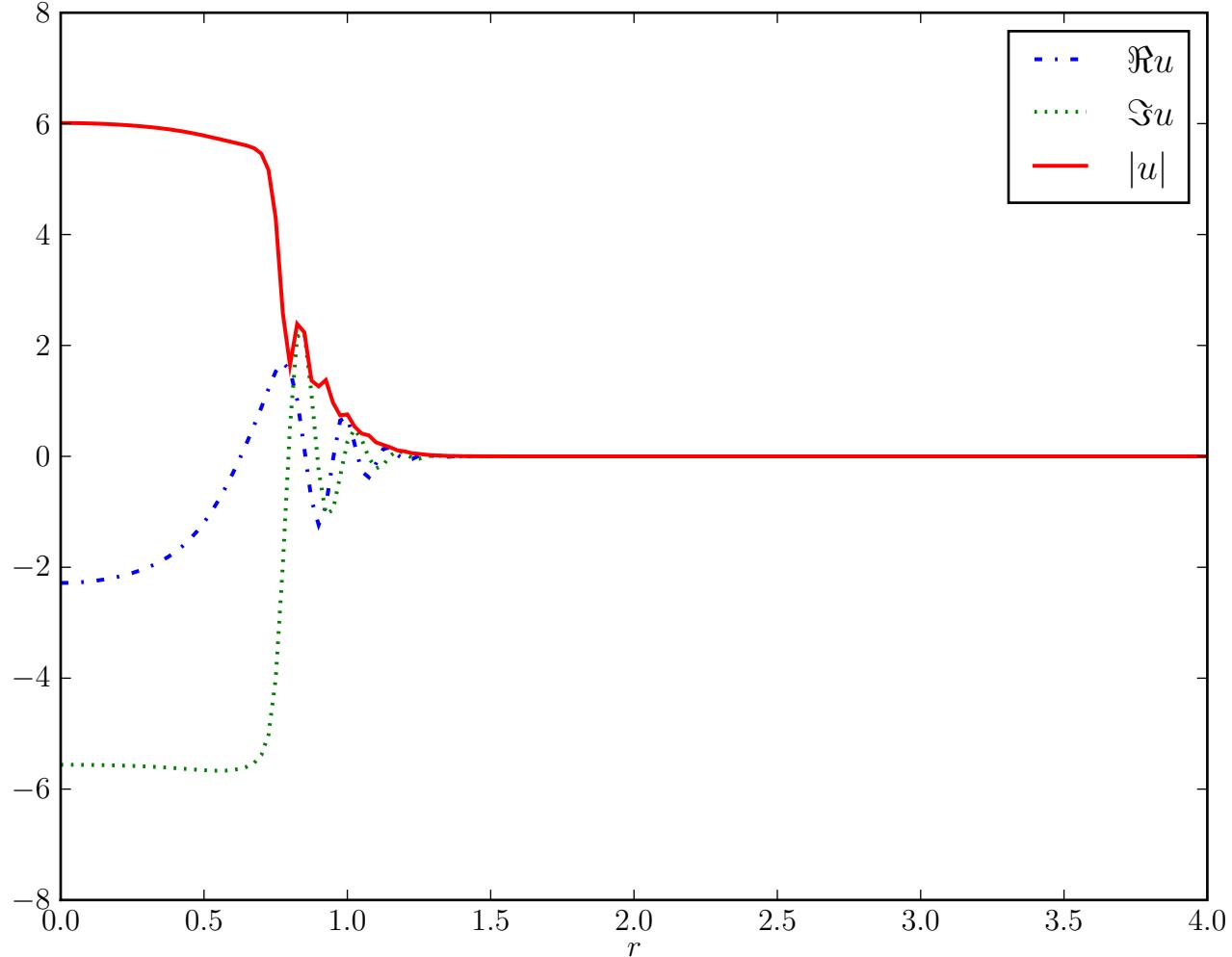
$t = 0$



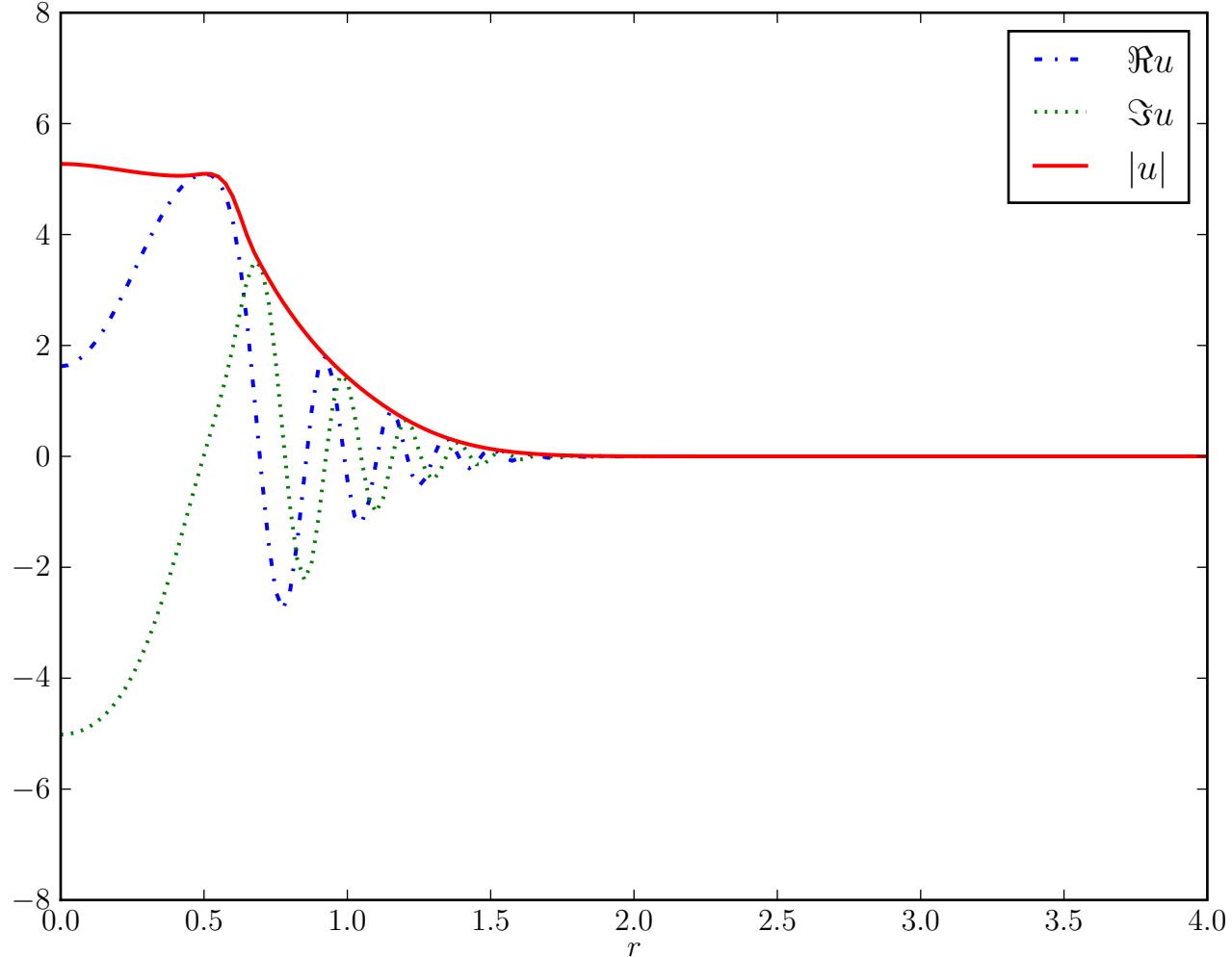




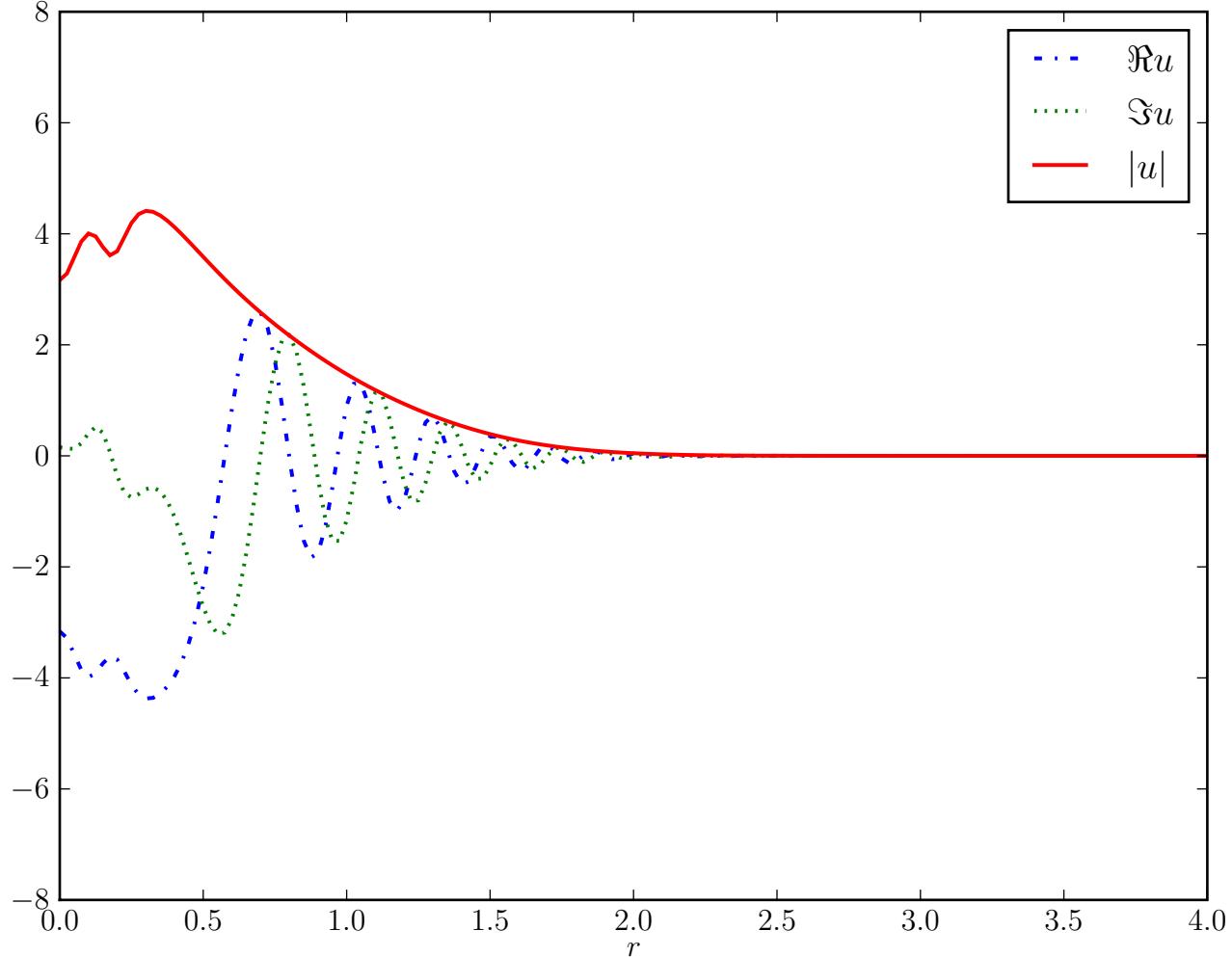
$t = 0.015$



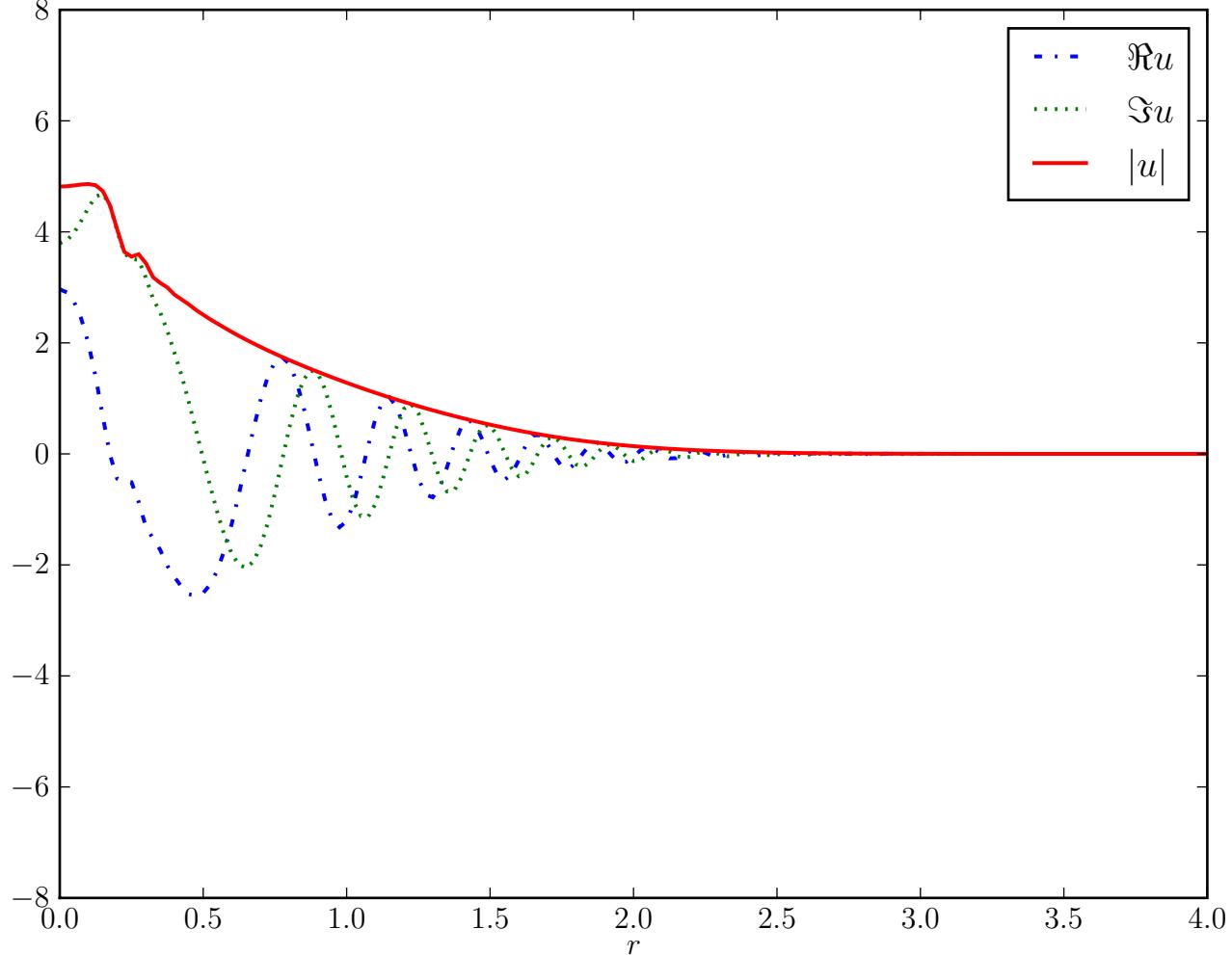
$t = 0.02$



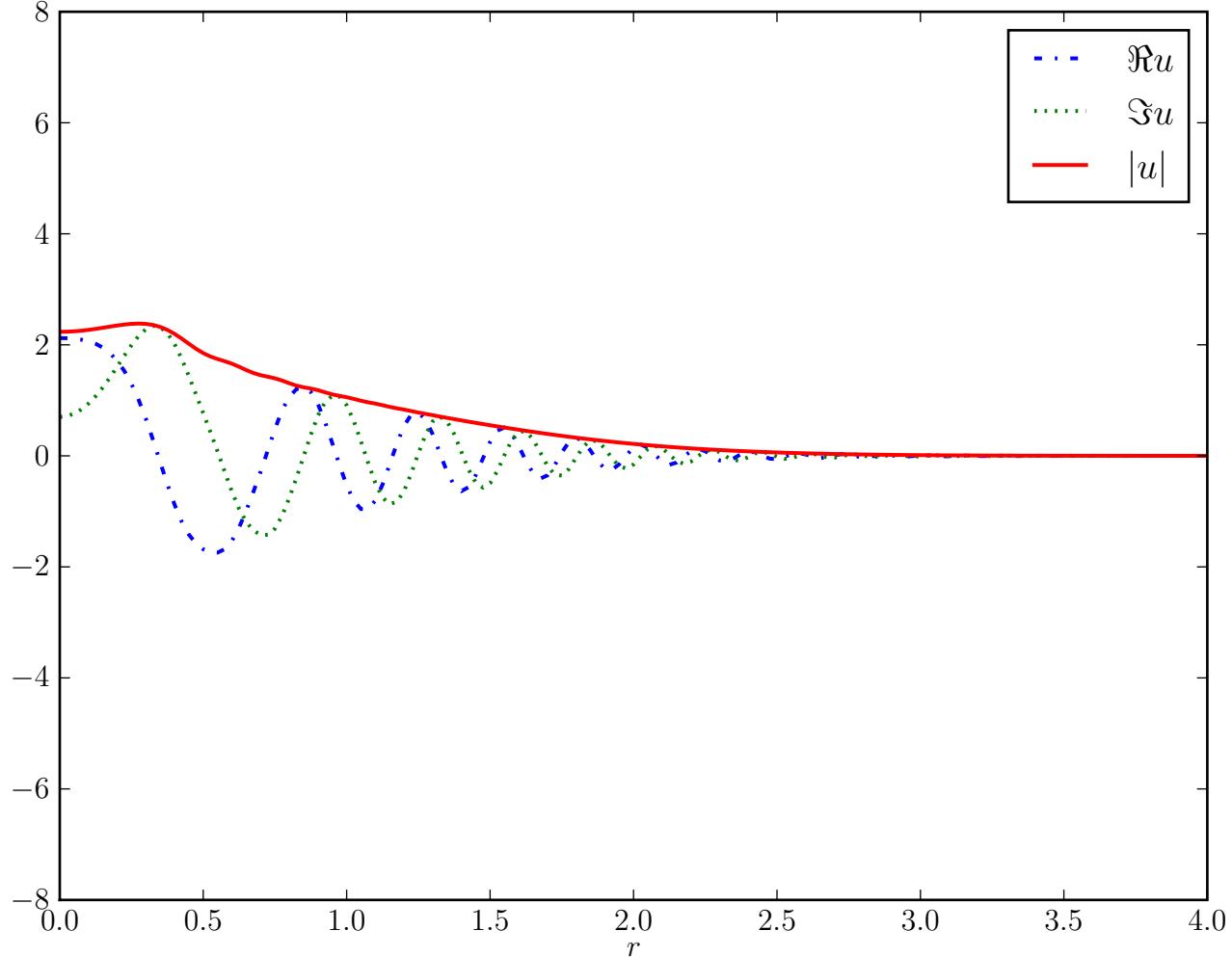
$t = 0.025$



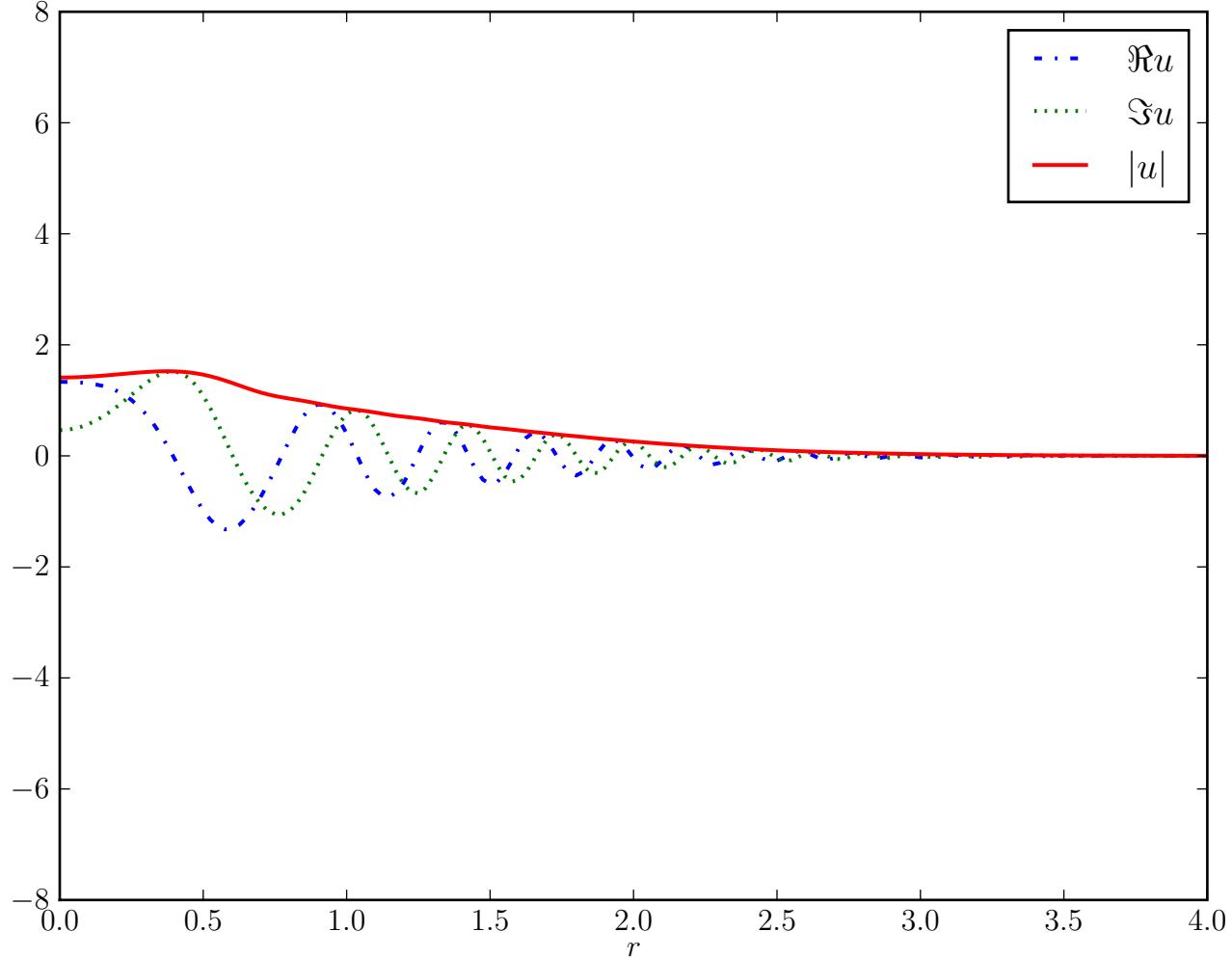
$t = 0.03$



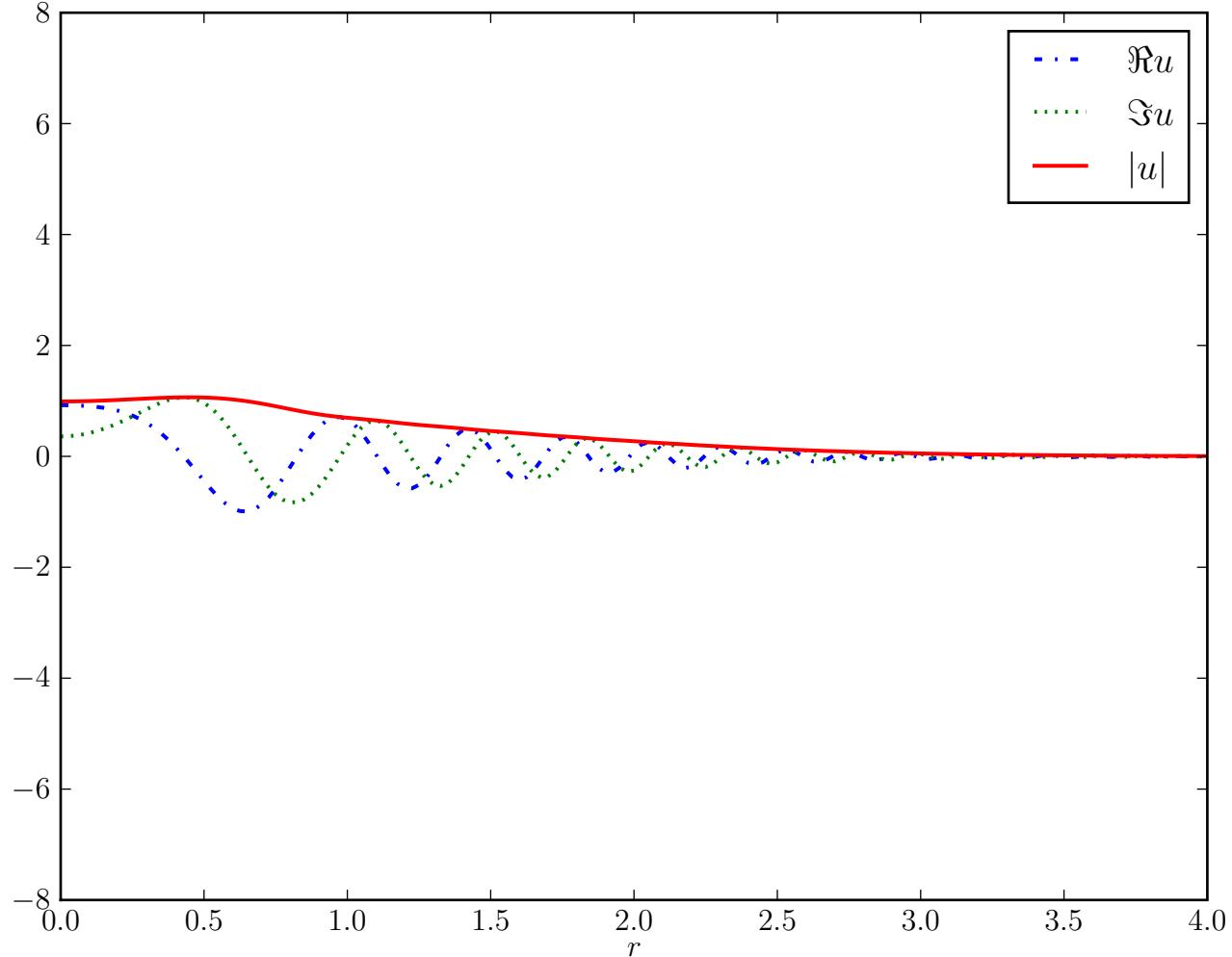
$t = 0.035$



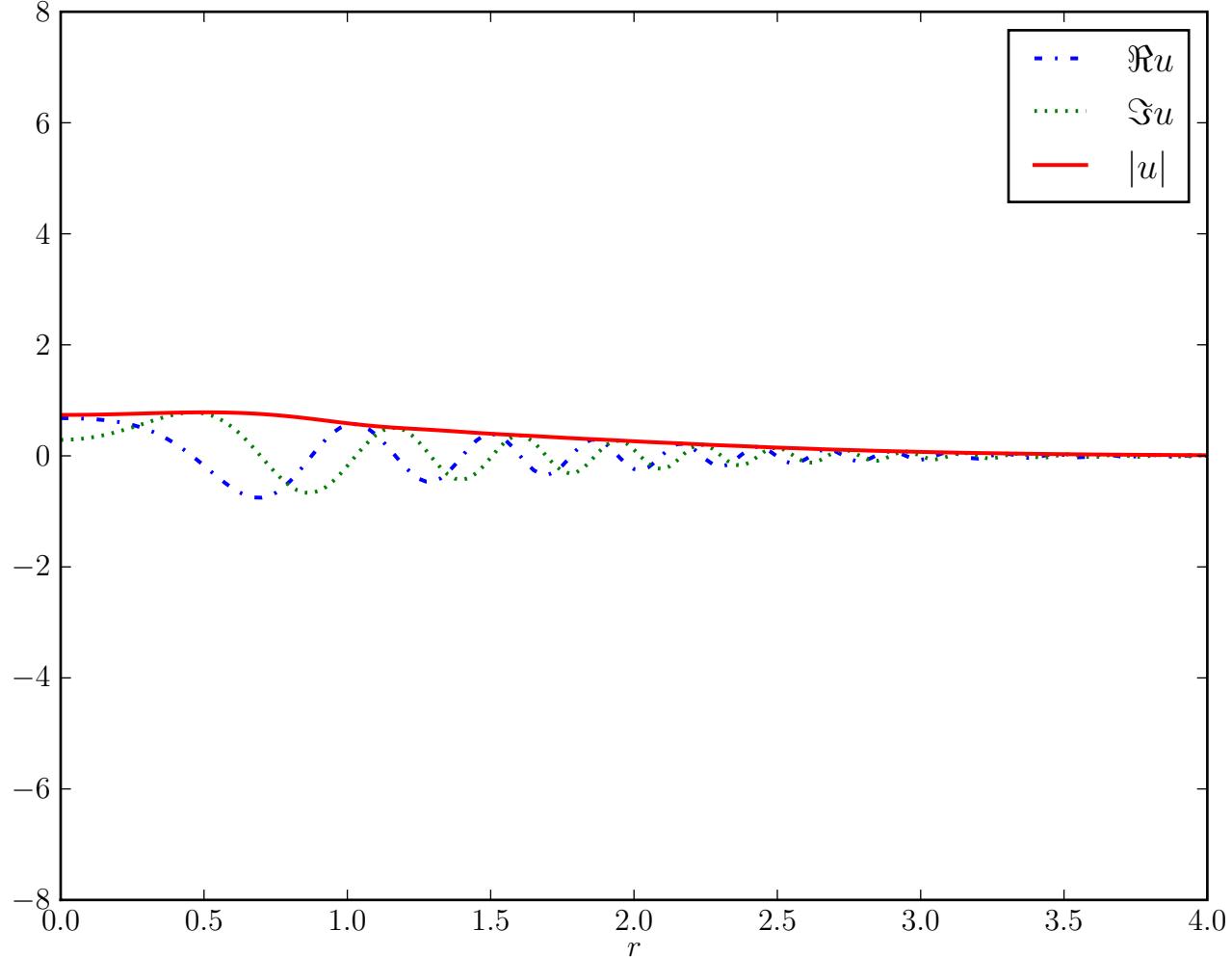
$t = 0.04$



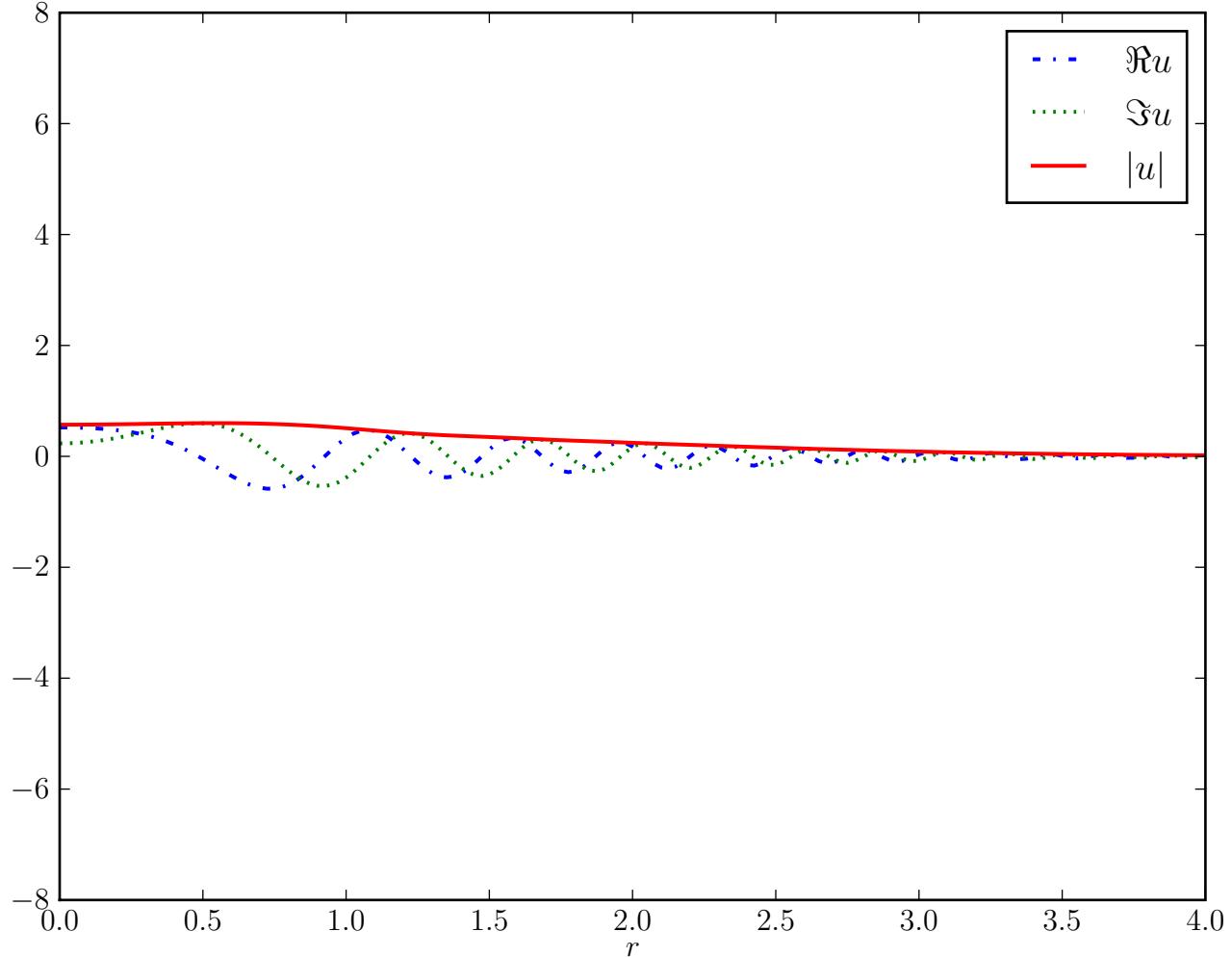
$t = 0.045$



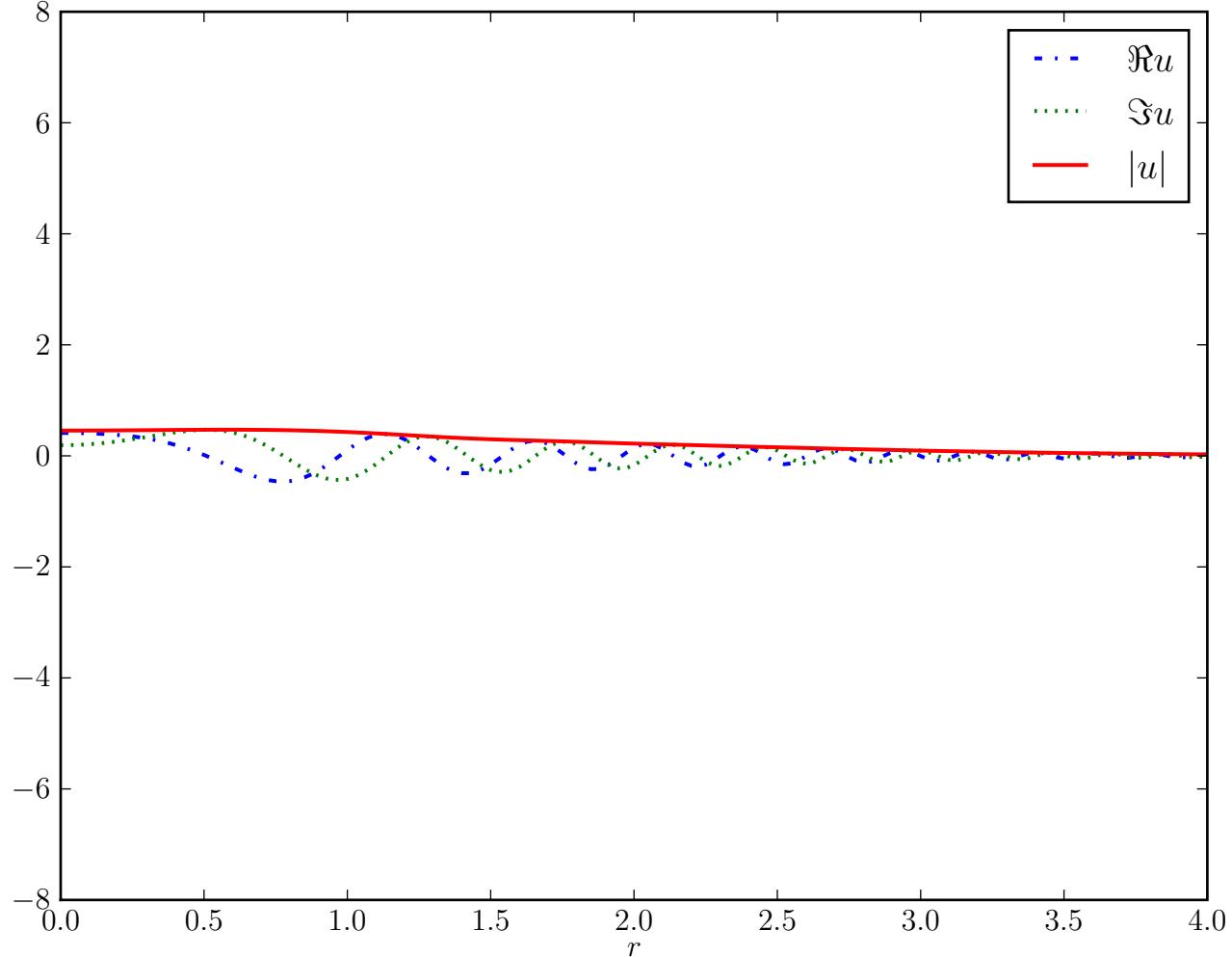
$t = 0.05$

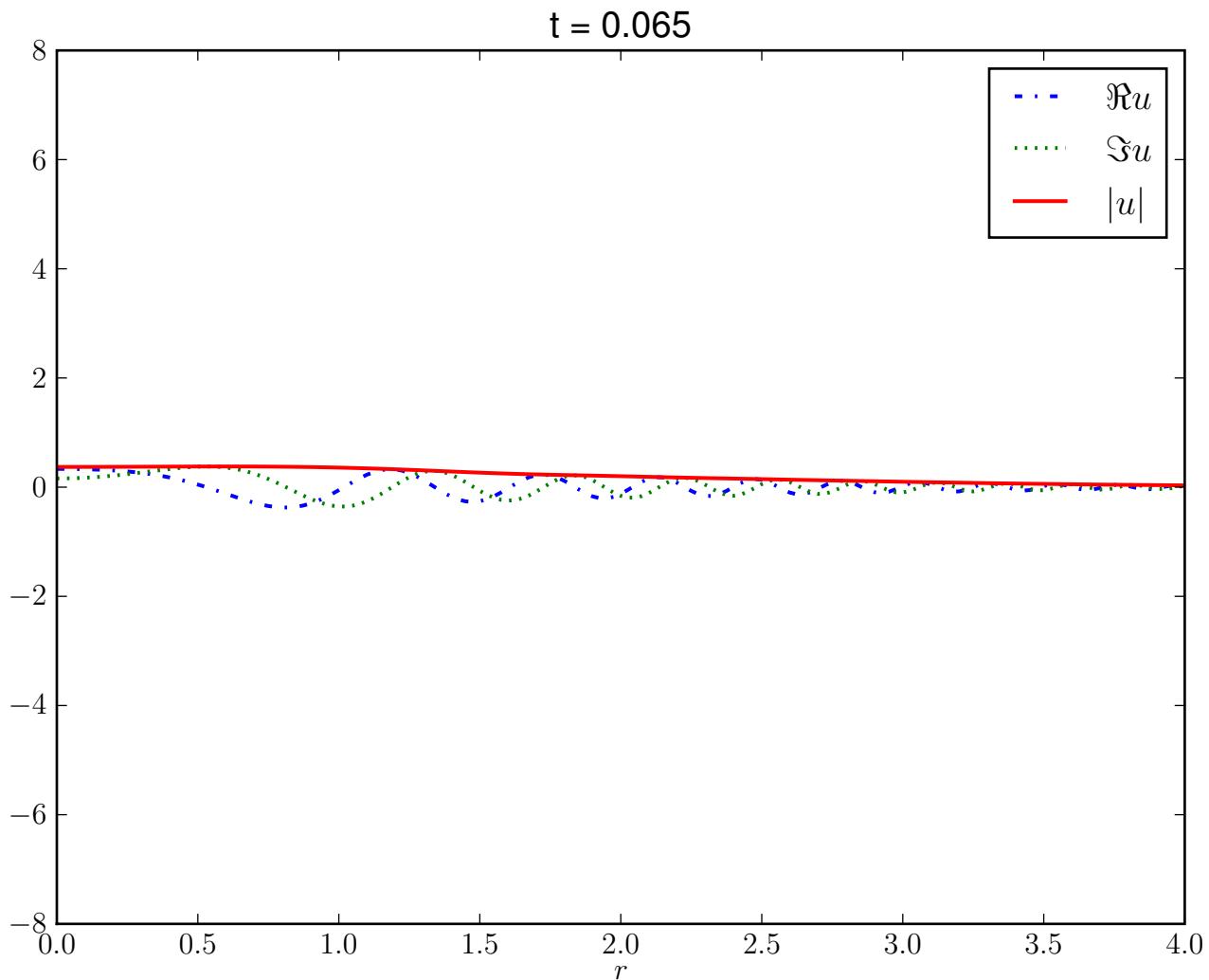


$t = 0.055$

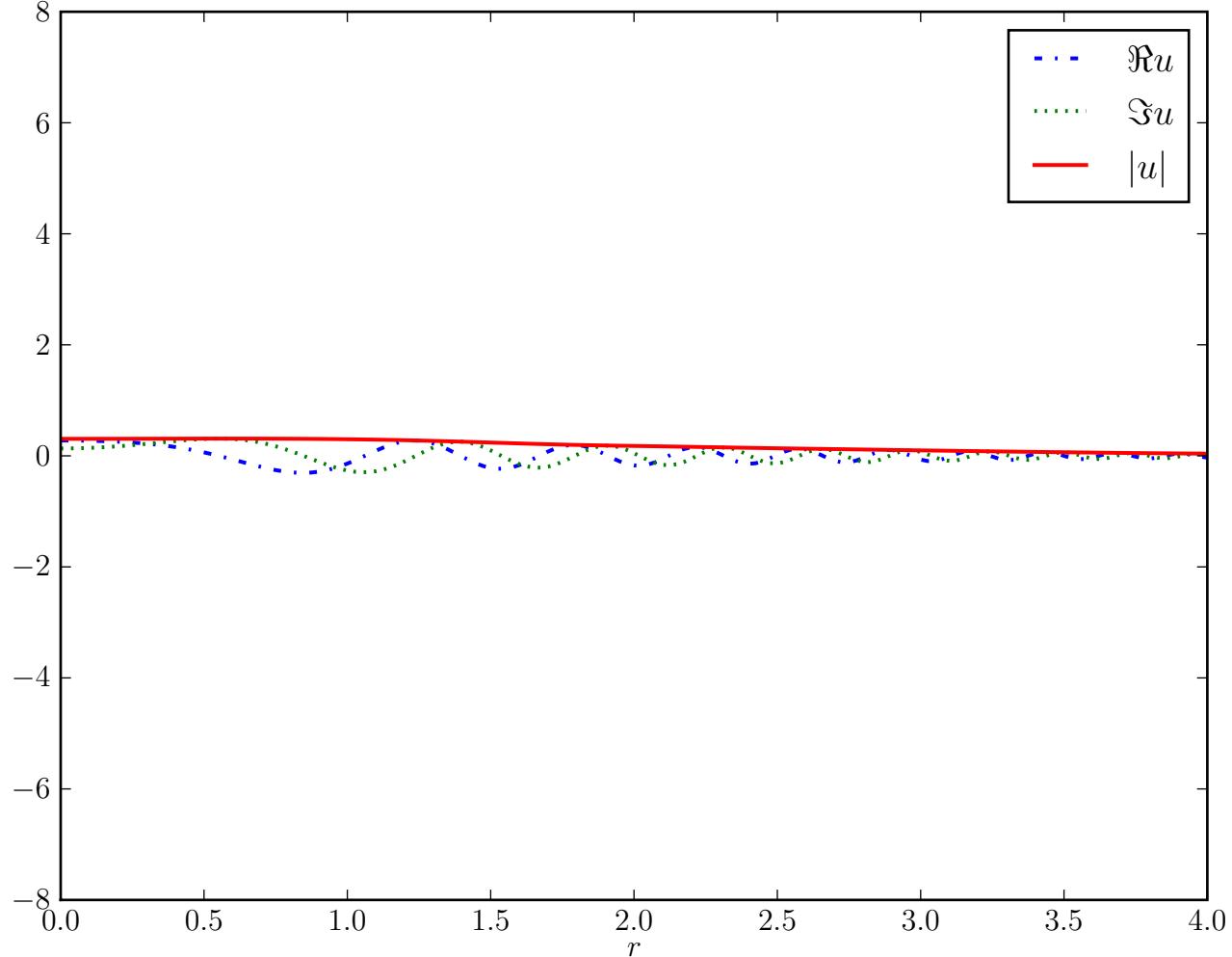


$t = 0.06$

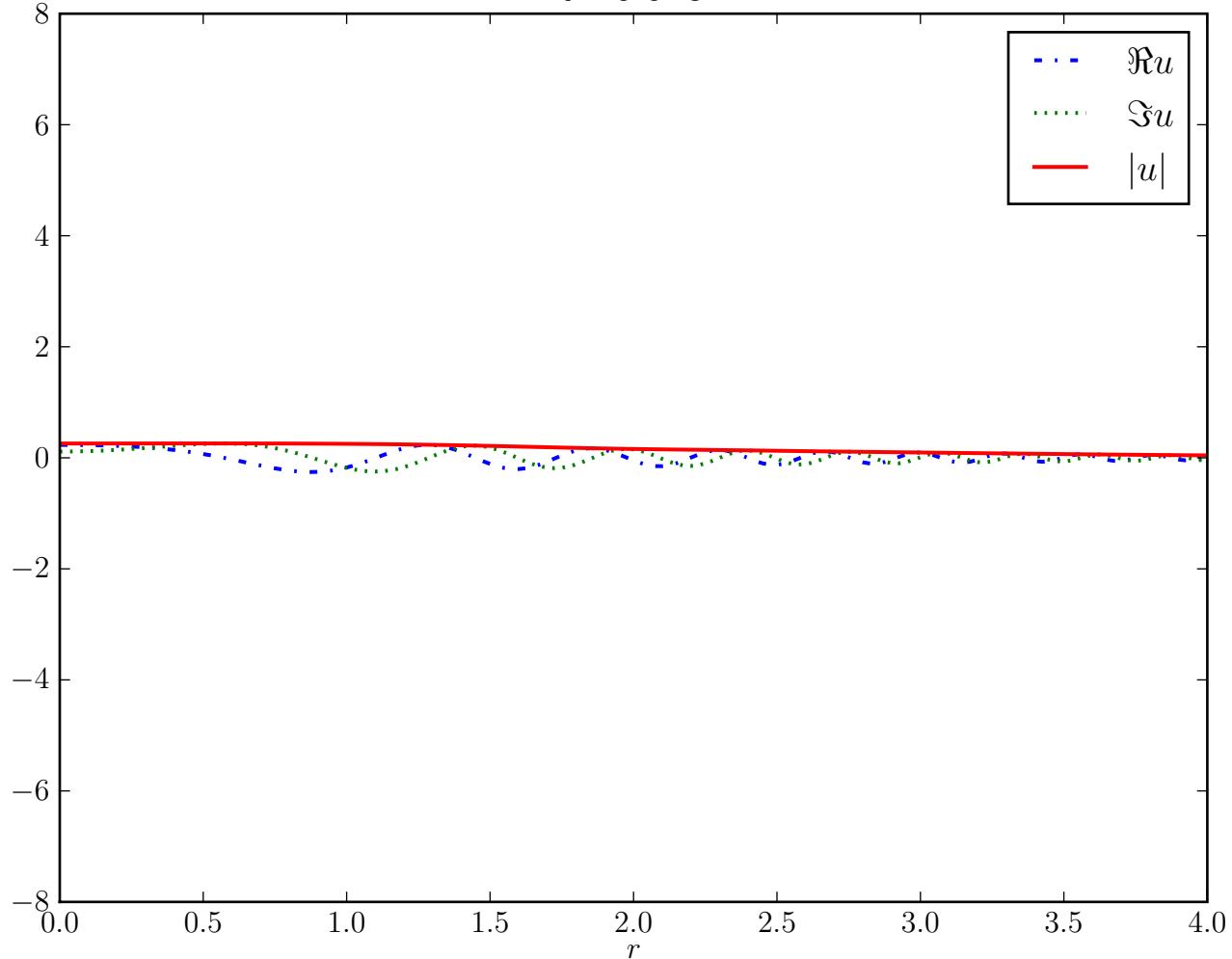




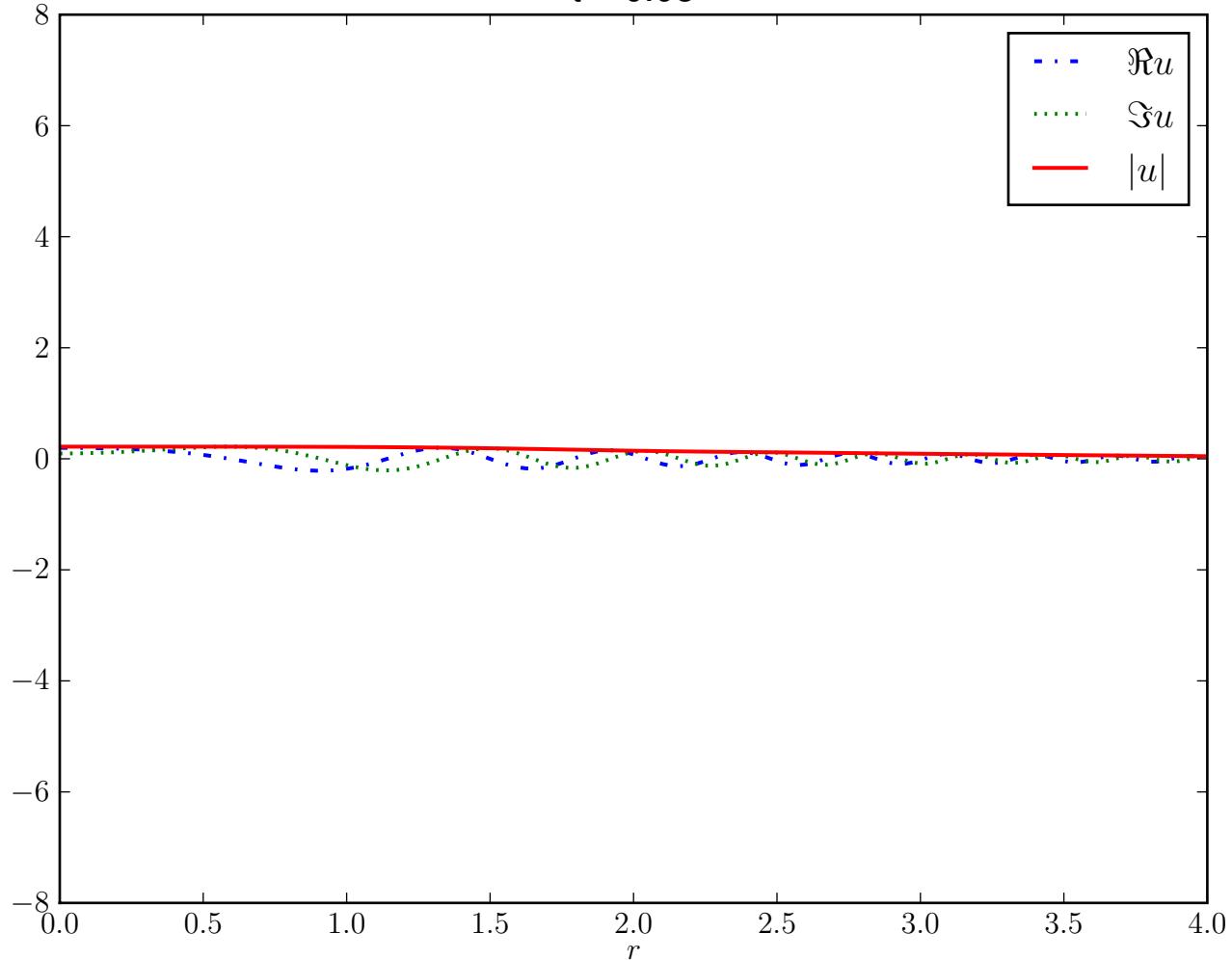
$t = 0.07$



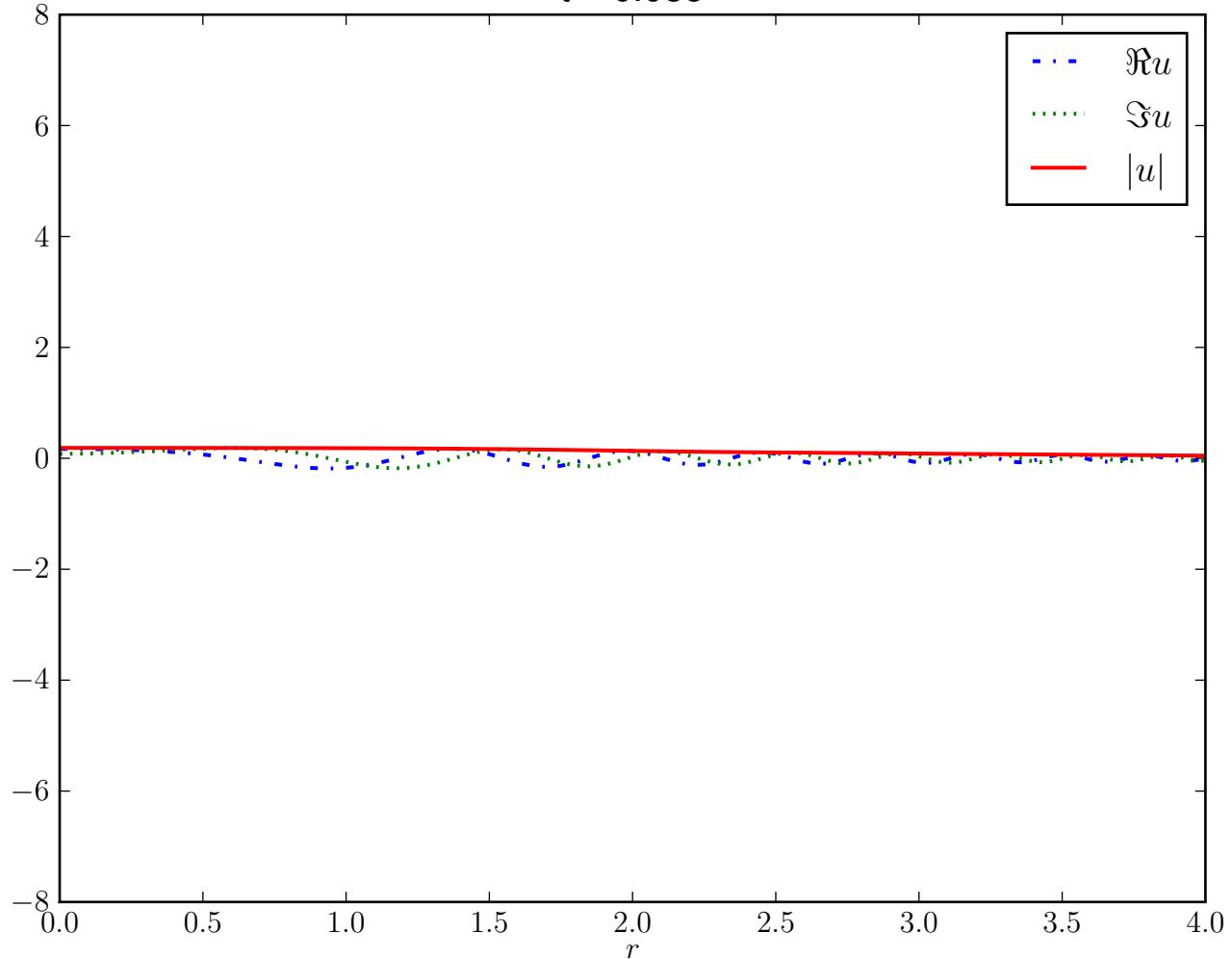
$t = 0.075$



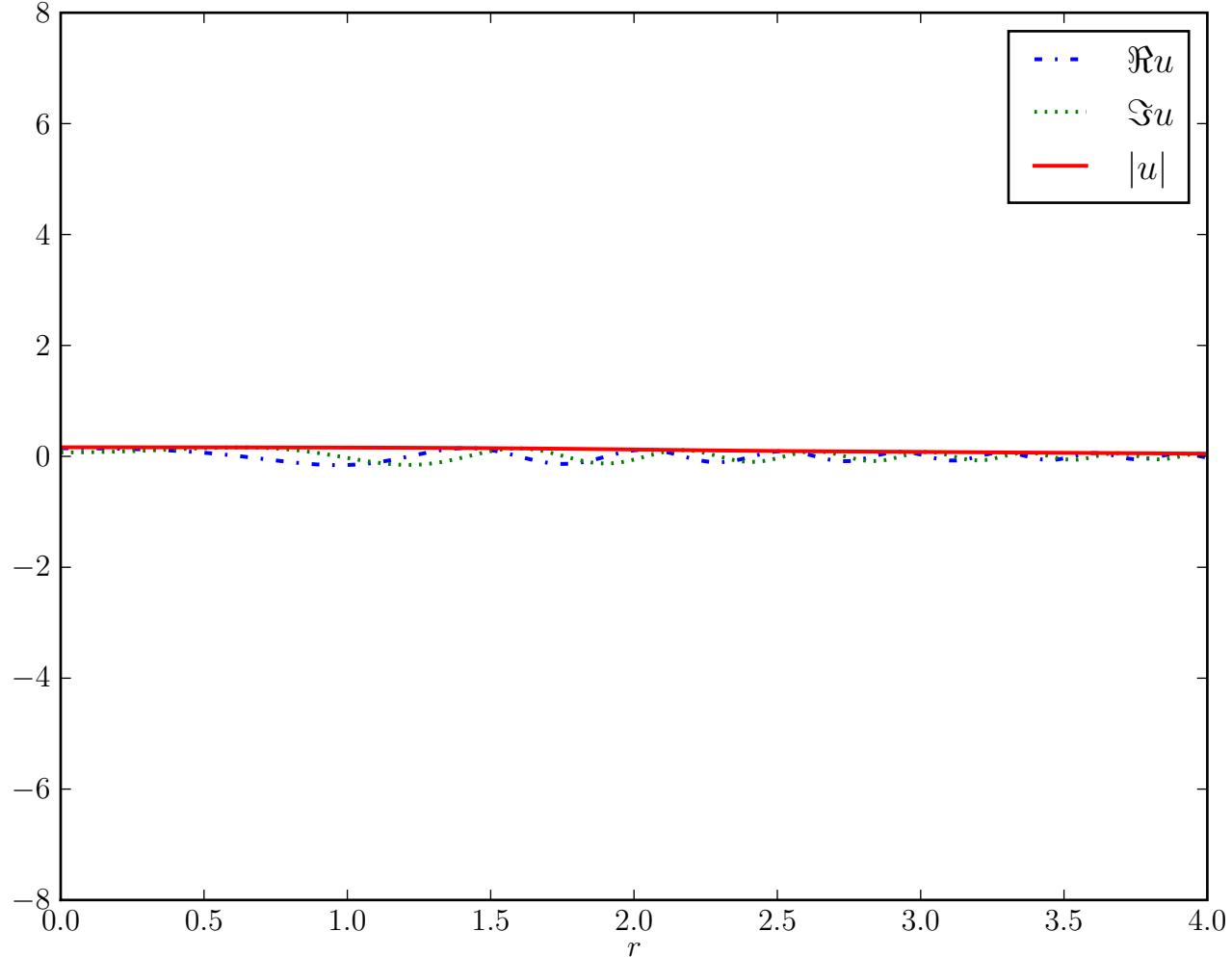
$t = 0.08$



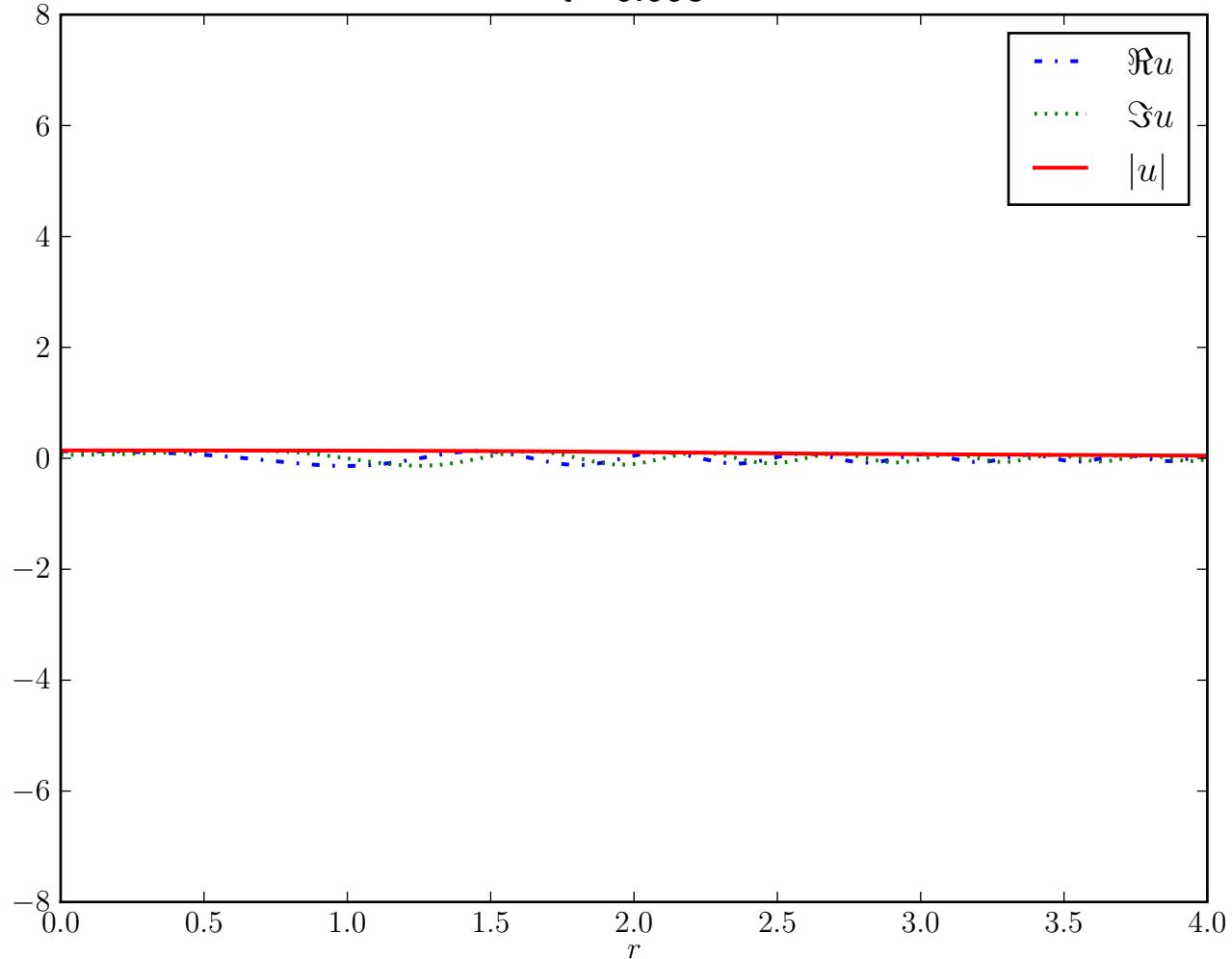
$t = 0.085$

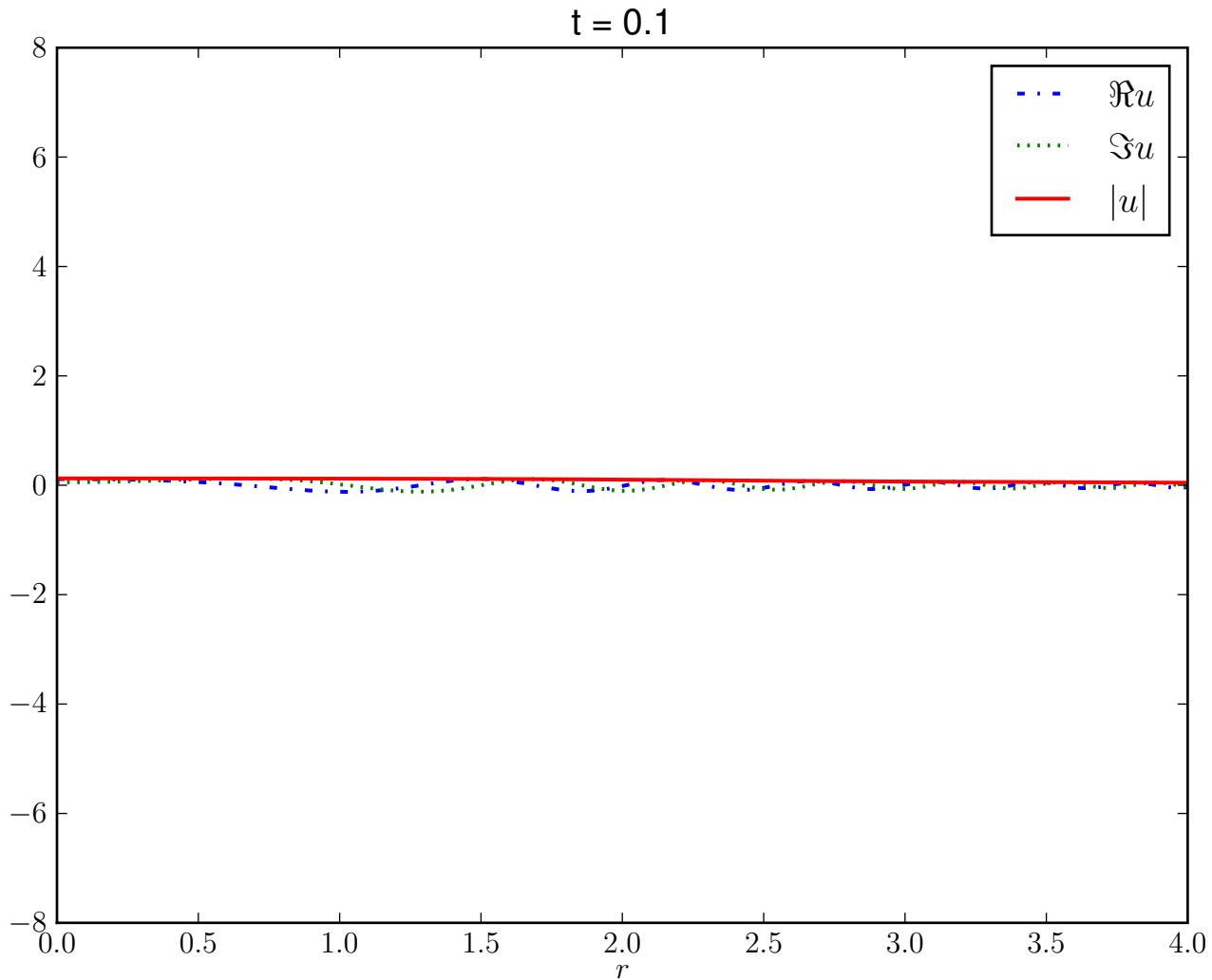


$t = 0.09$

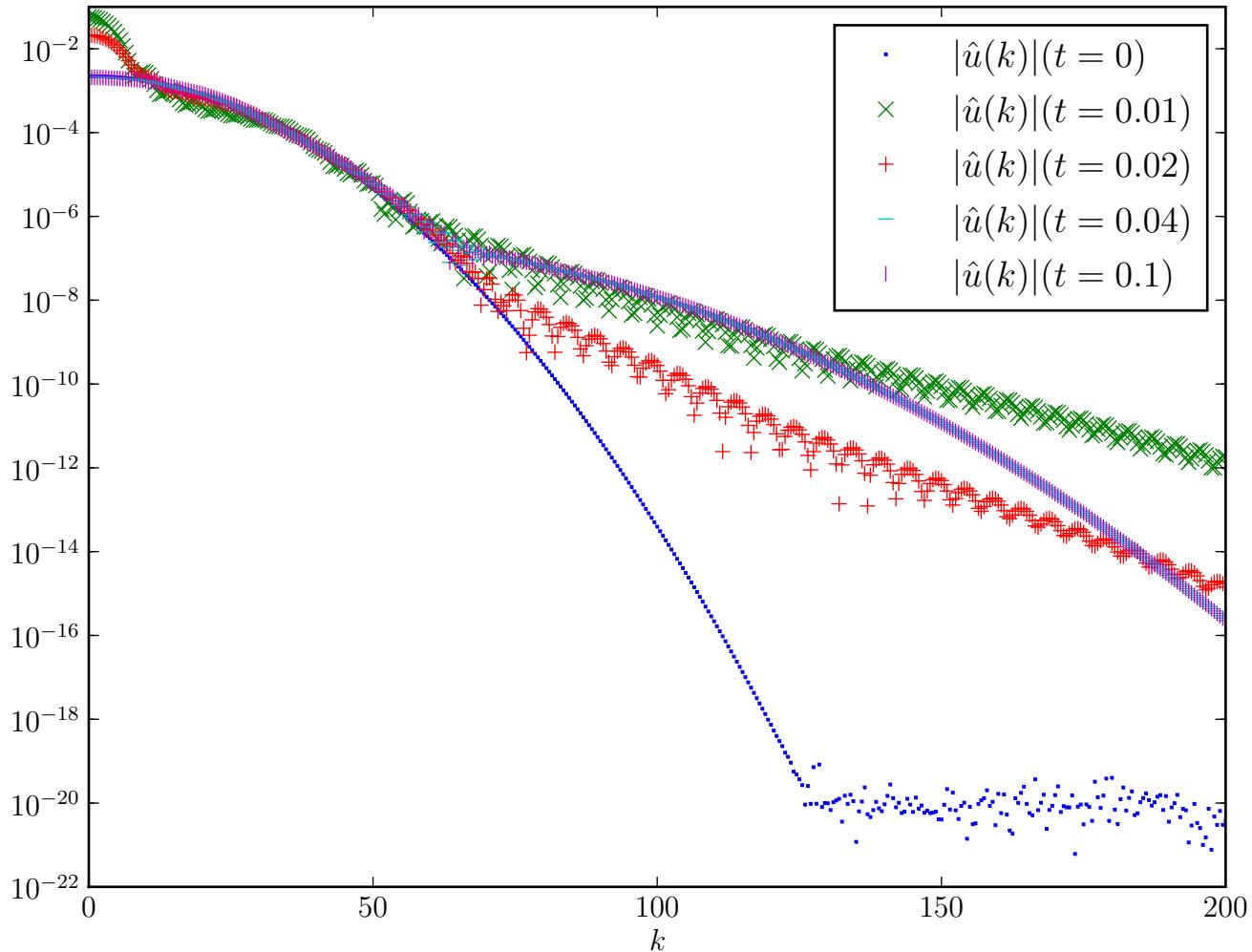


$t = 0.095$

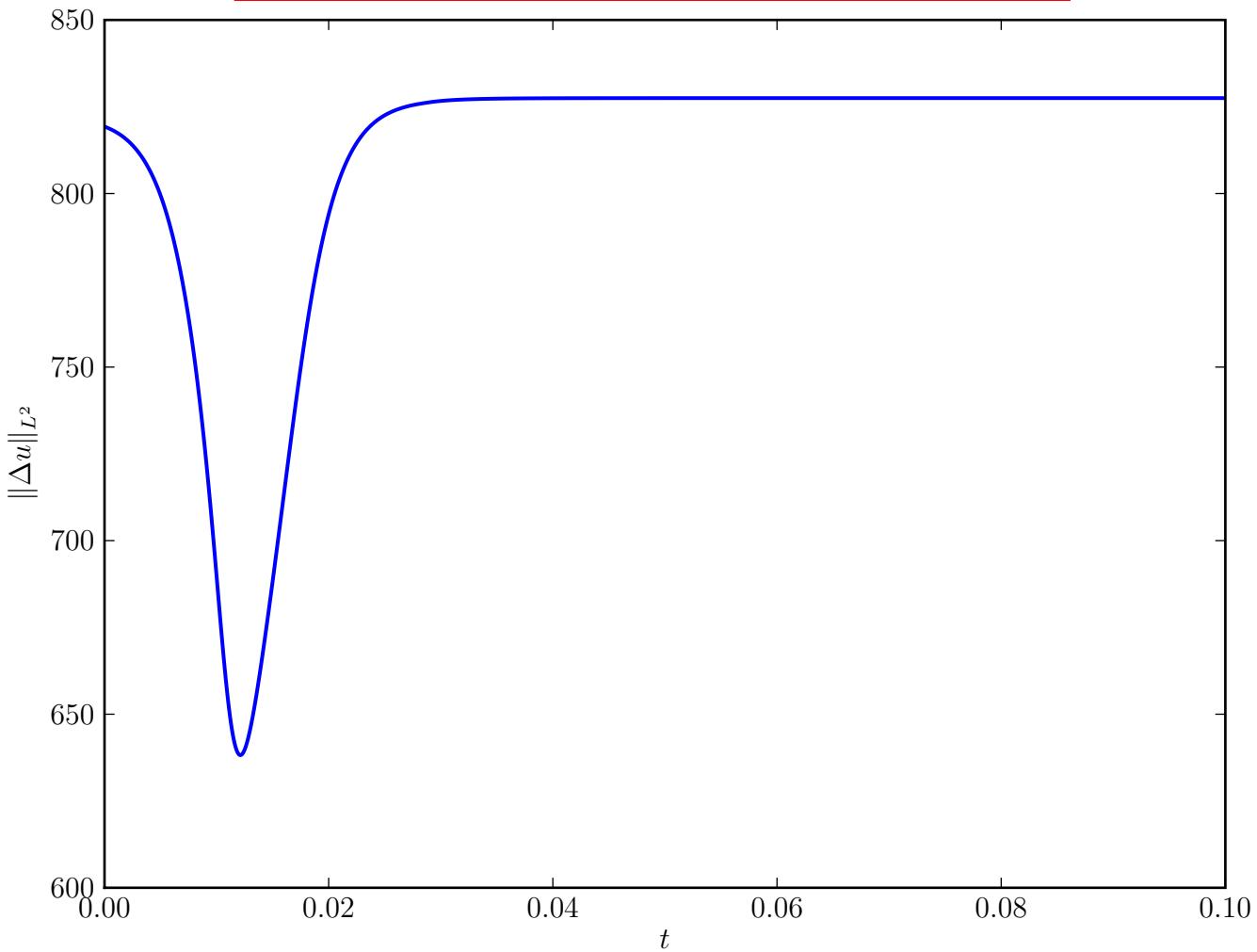




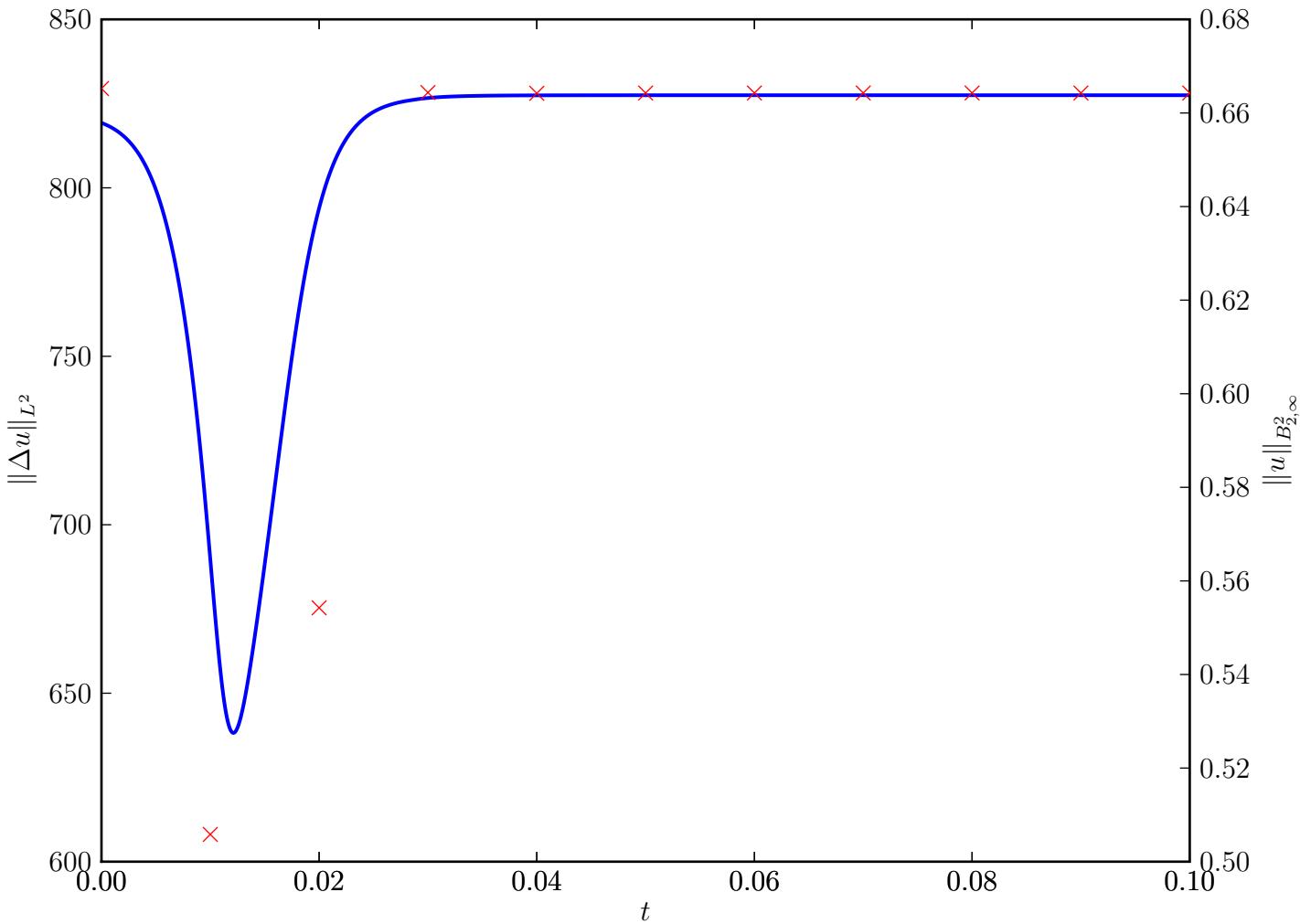
Phased Centered Gaussian Fourier transform snapshots along nonlinear flow



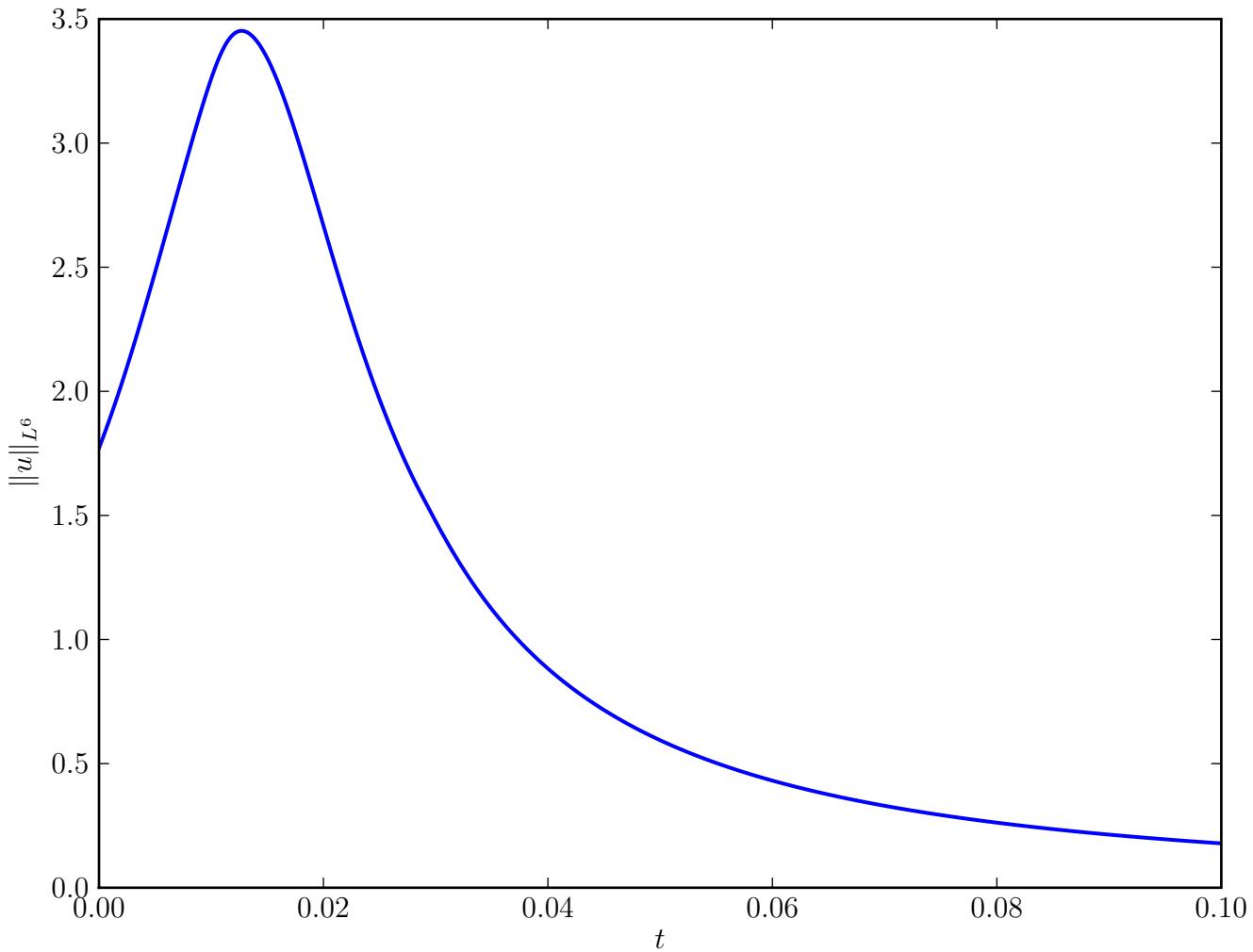
H^2 norm of Phased Centered Gaussian along nonlinear flow



Sobolev vs. Besov: Phased Centered Gaussian along nonlinear flow



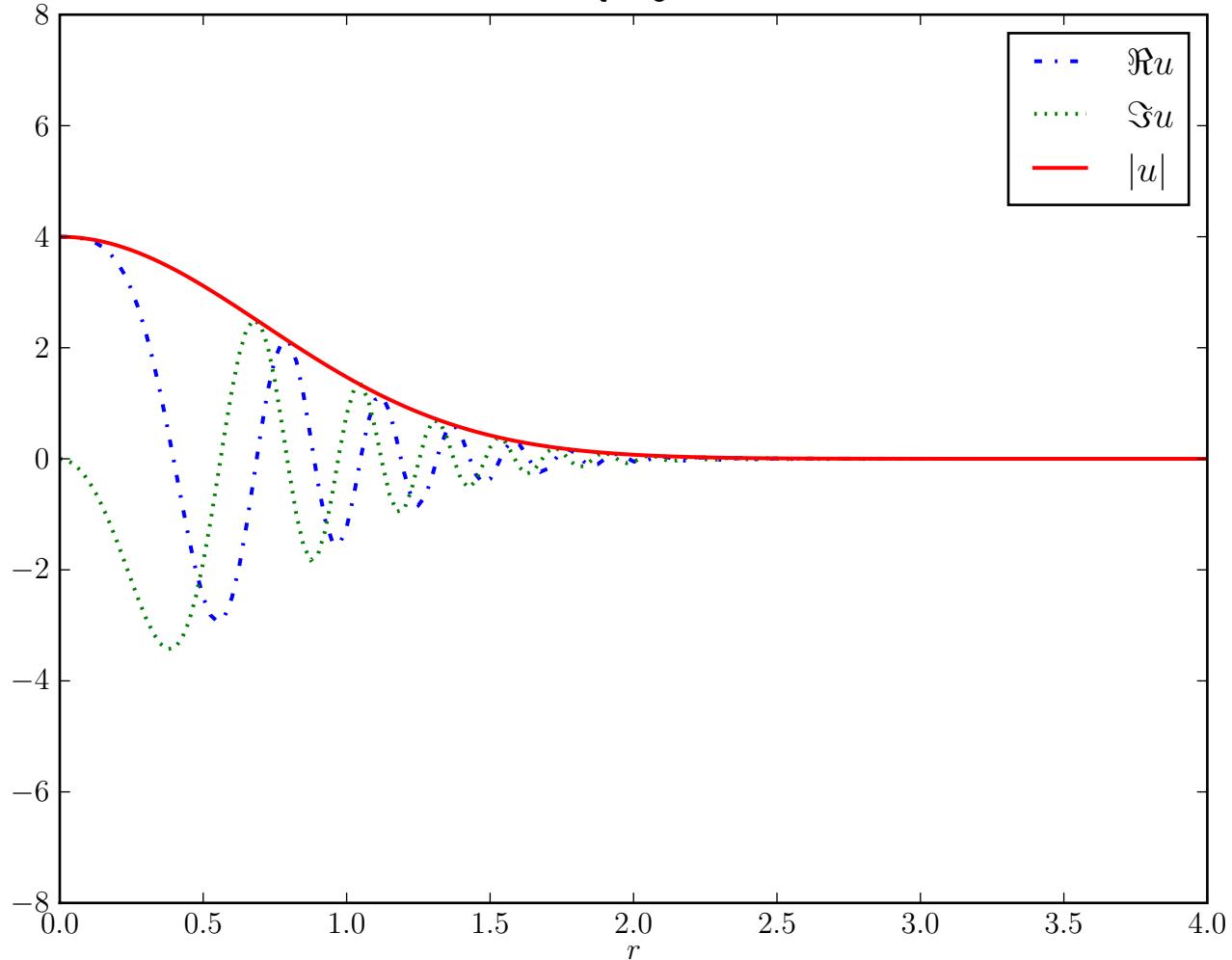
Potential Energy Norm Decay: Phased Centered Gaussian along nonlinear flow

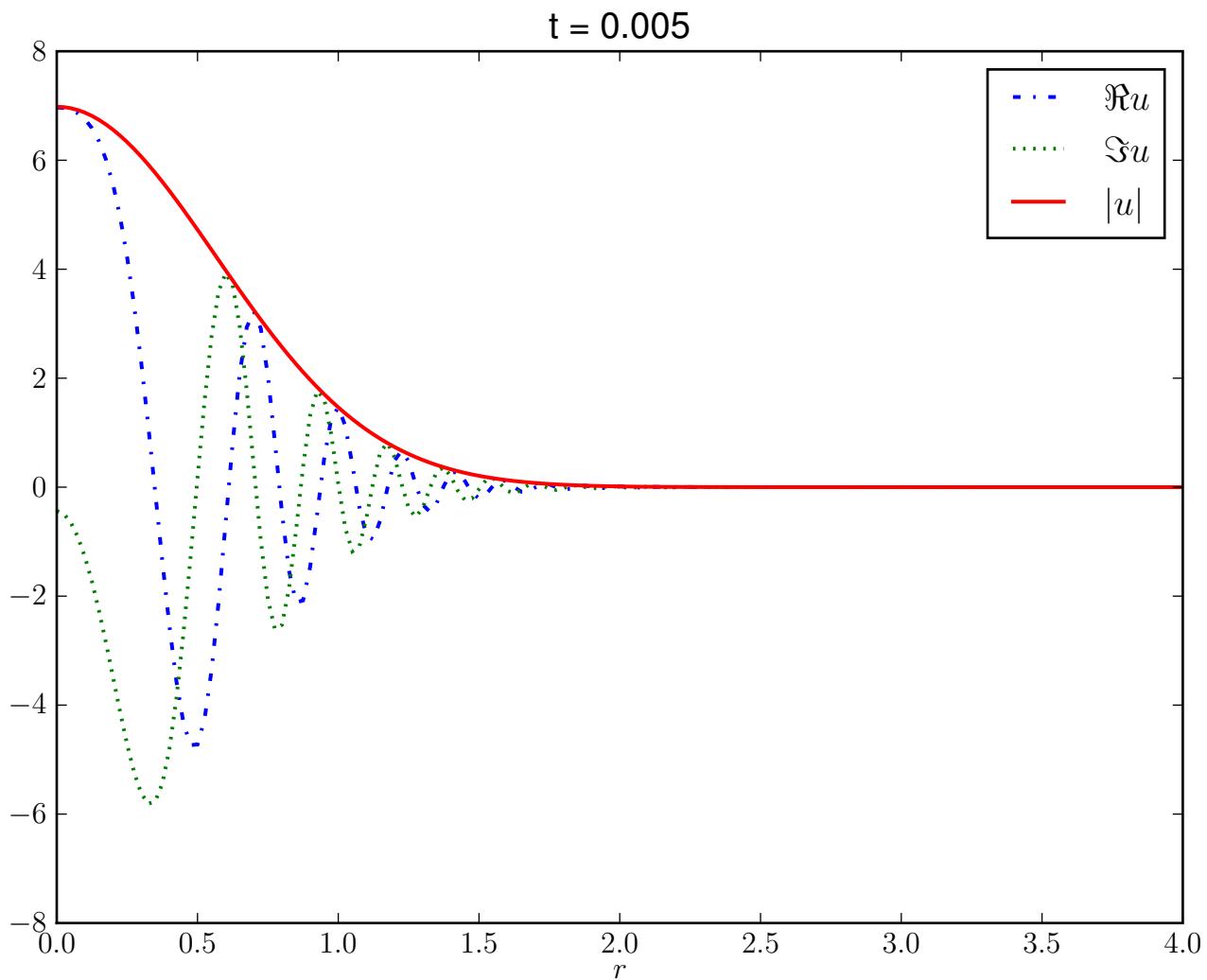


Phased Centered Gaussian Initial Data

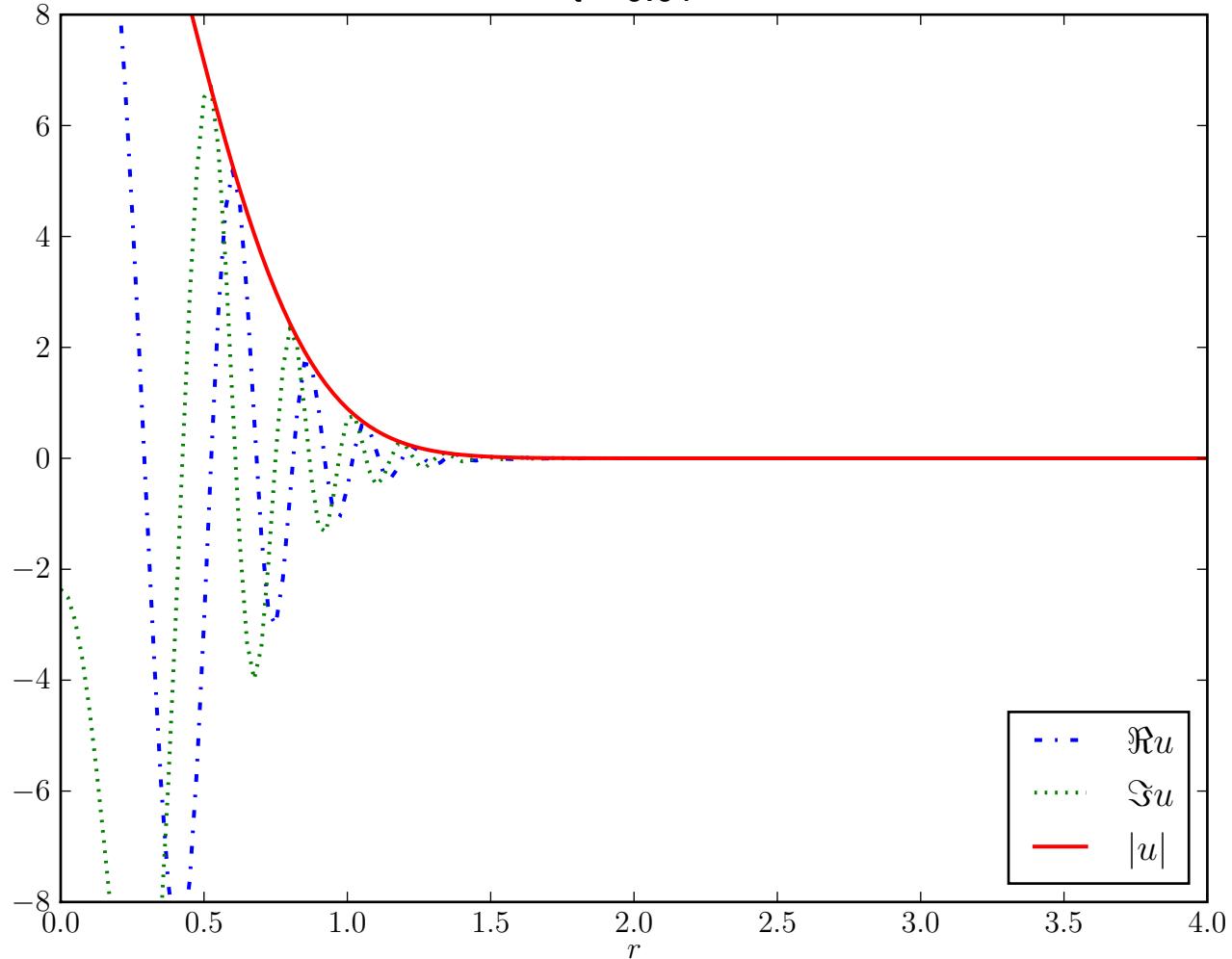
Linear Flow

$t = 0$

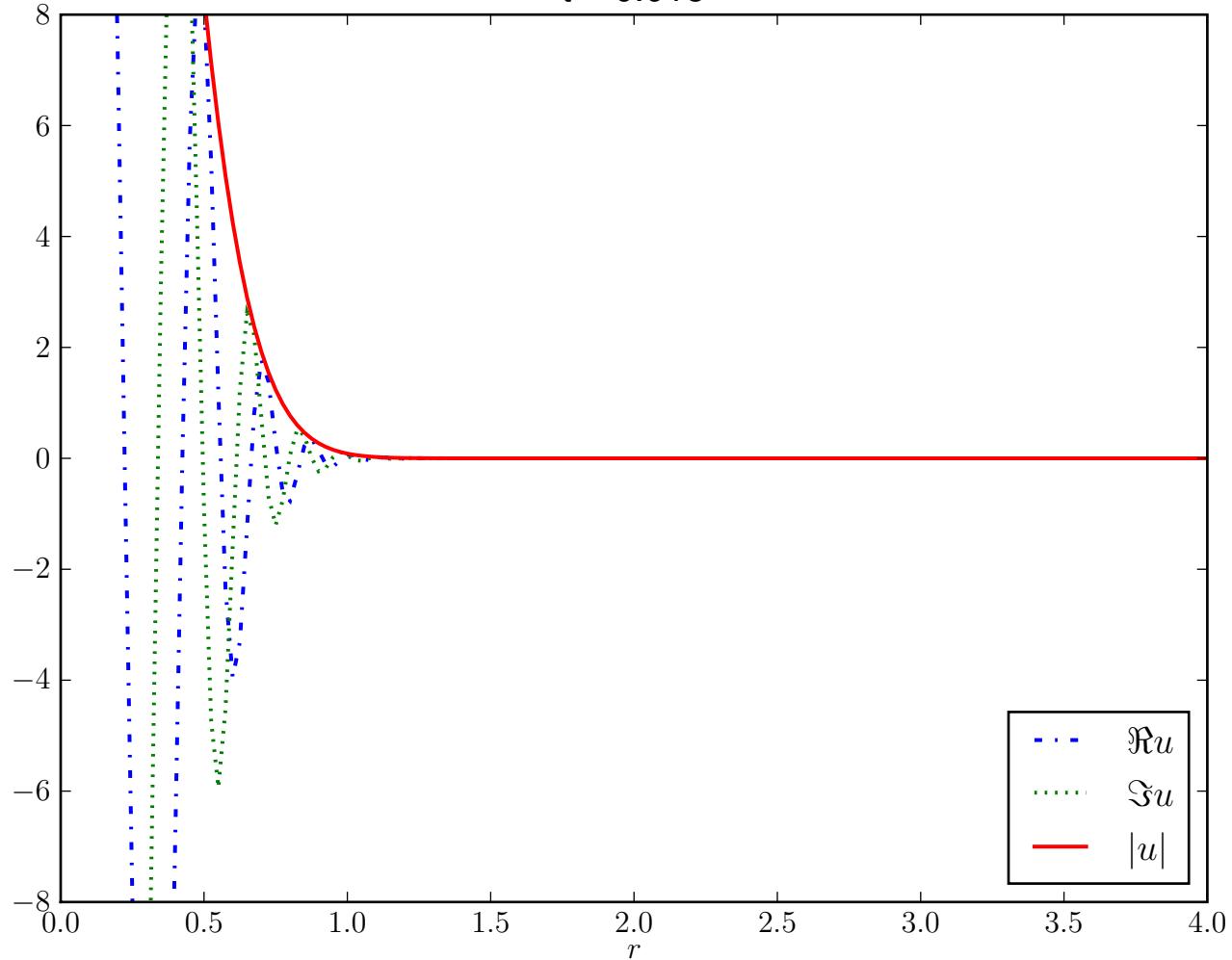




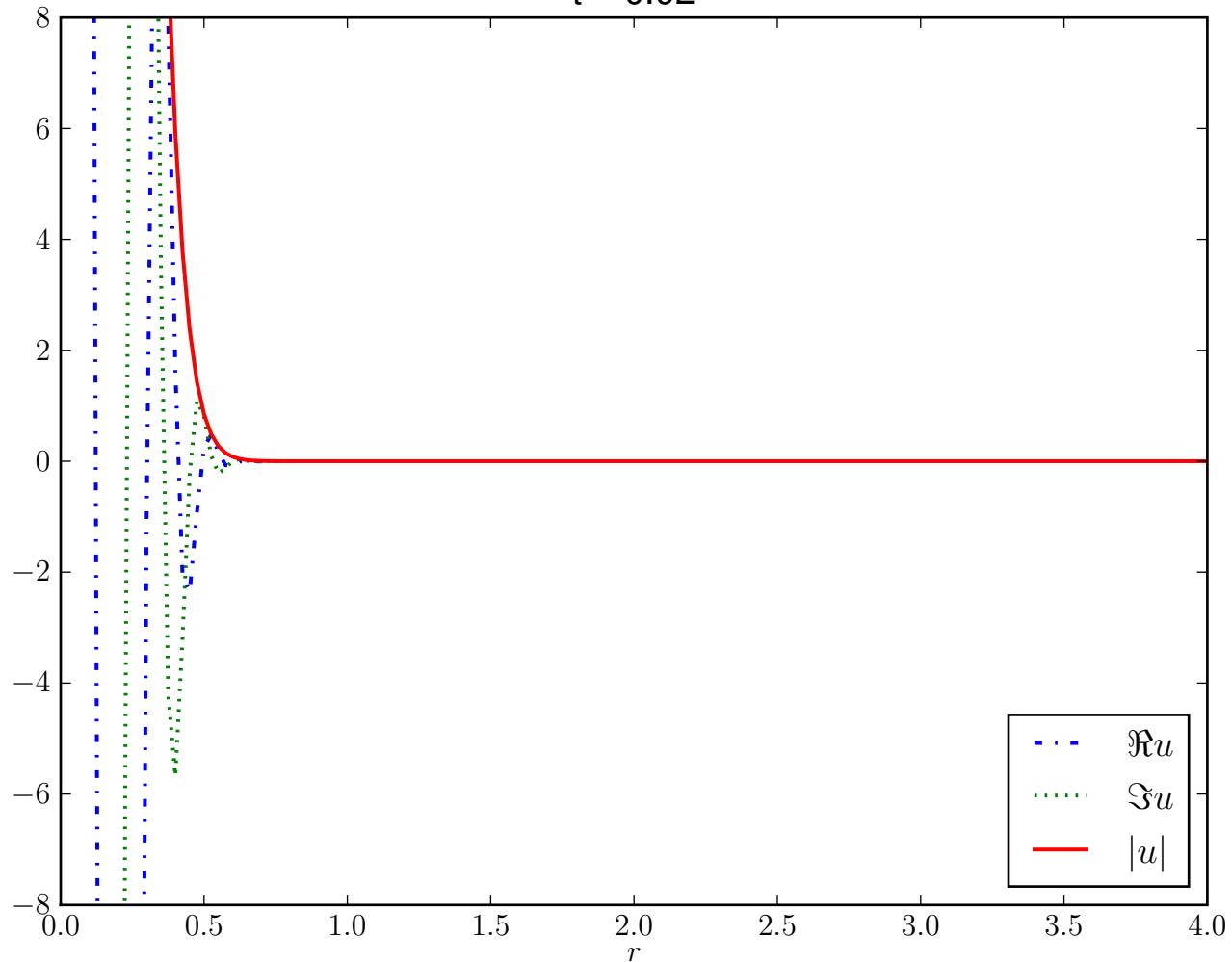
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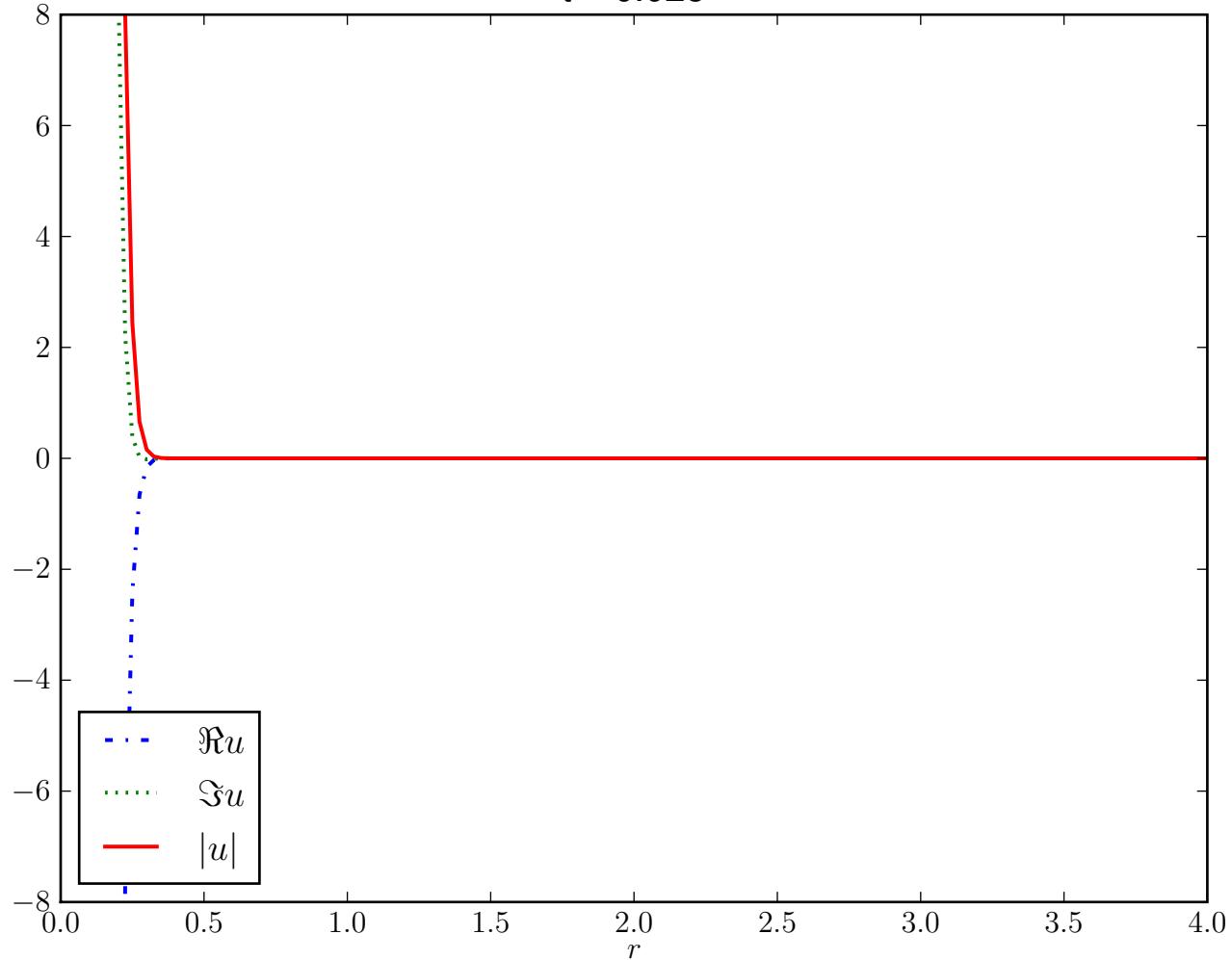
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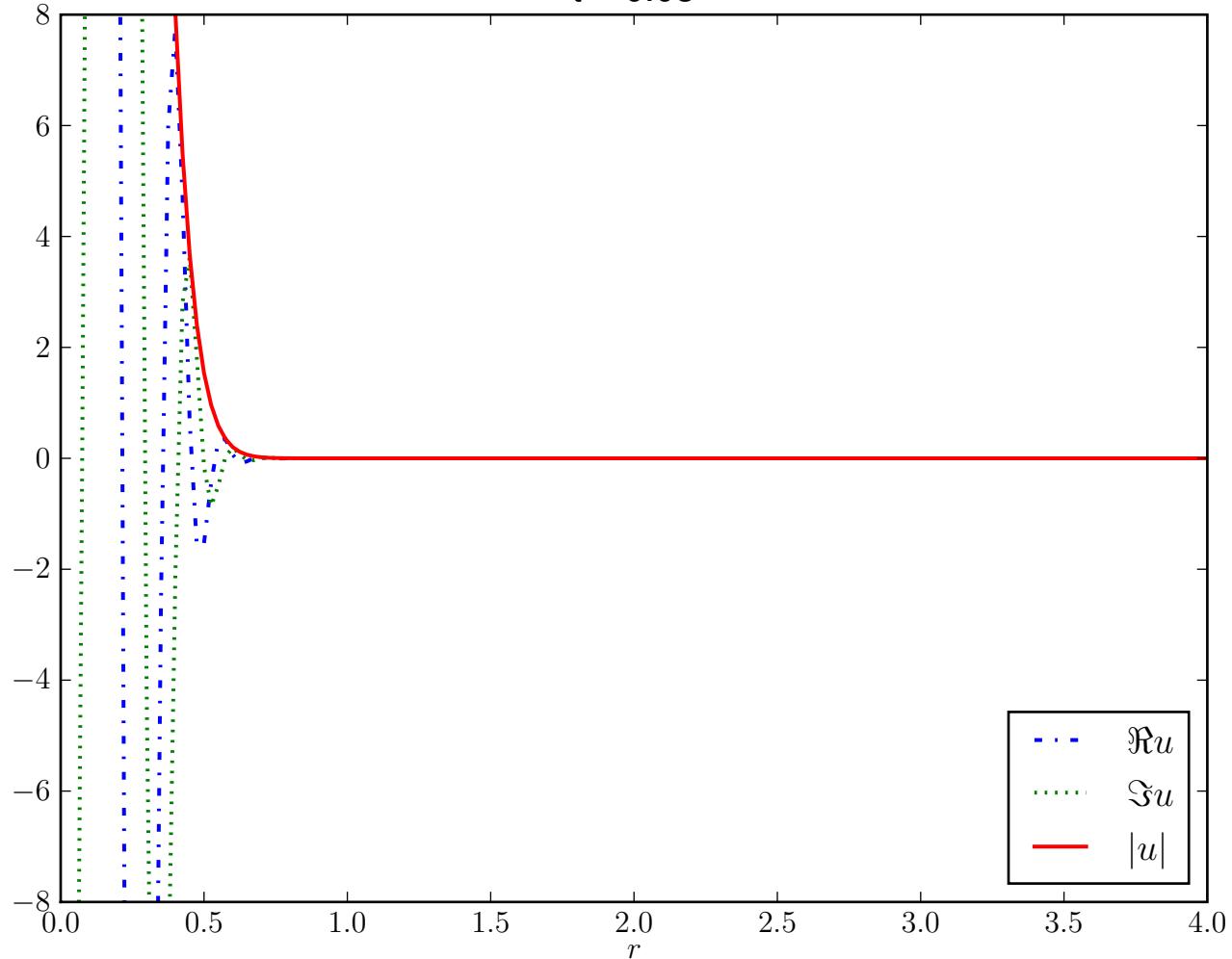
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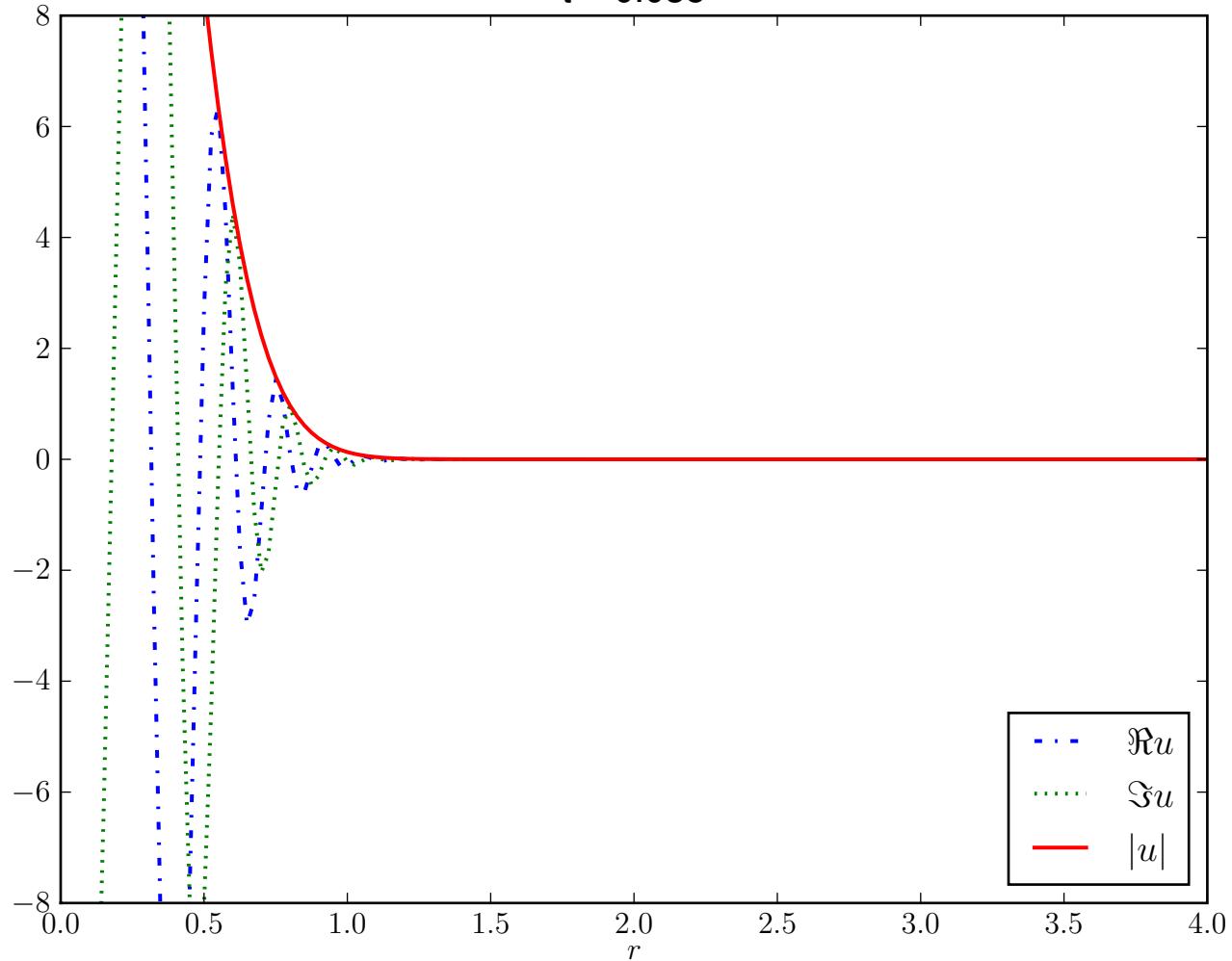
$t = 0.025$



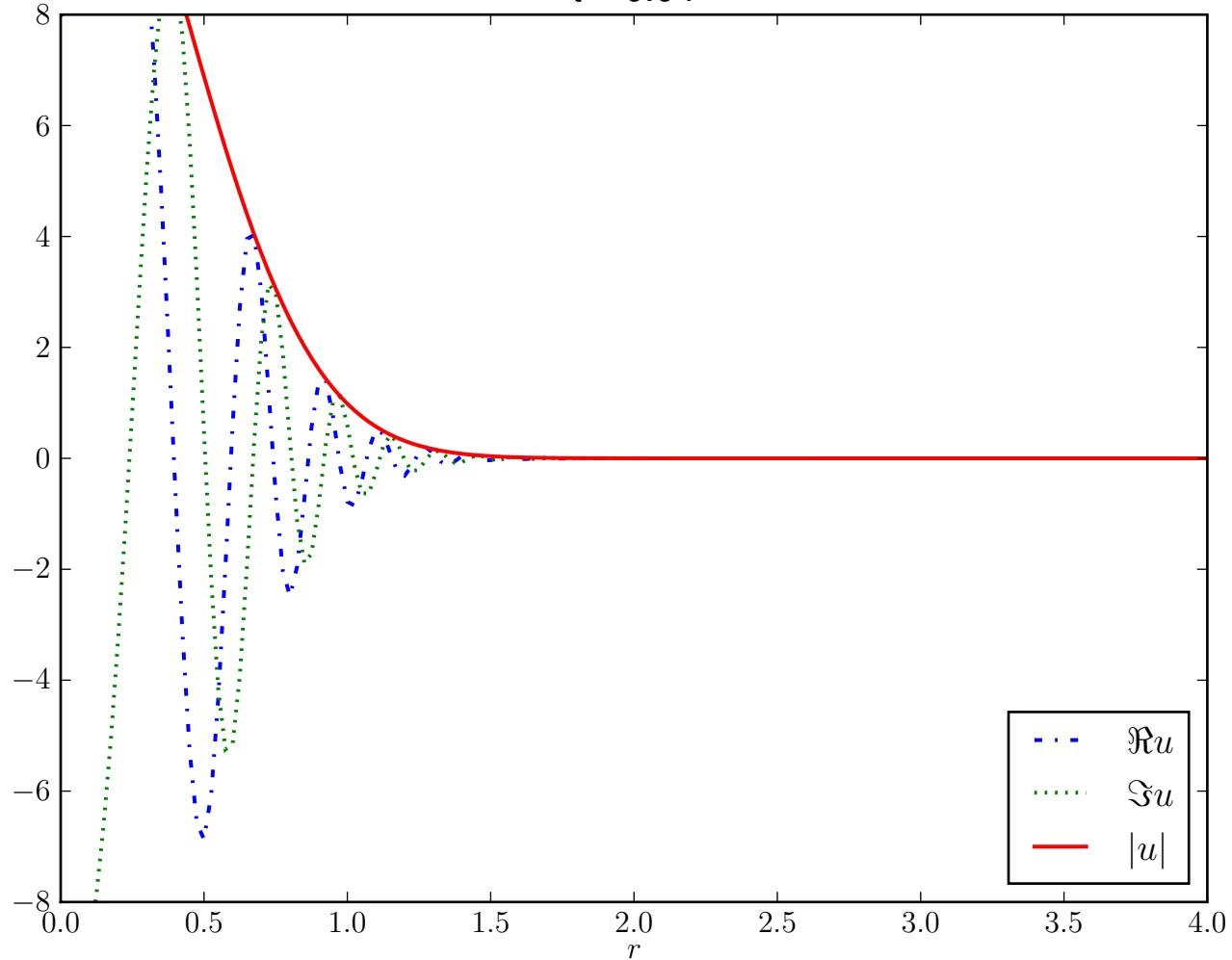
$t = 0.03$



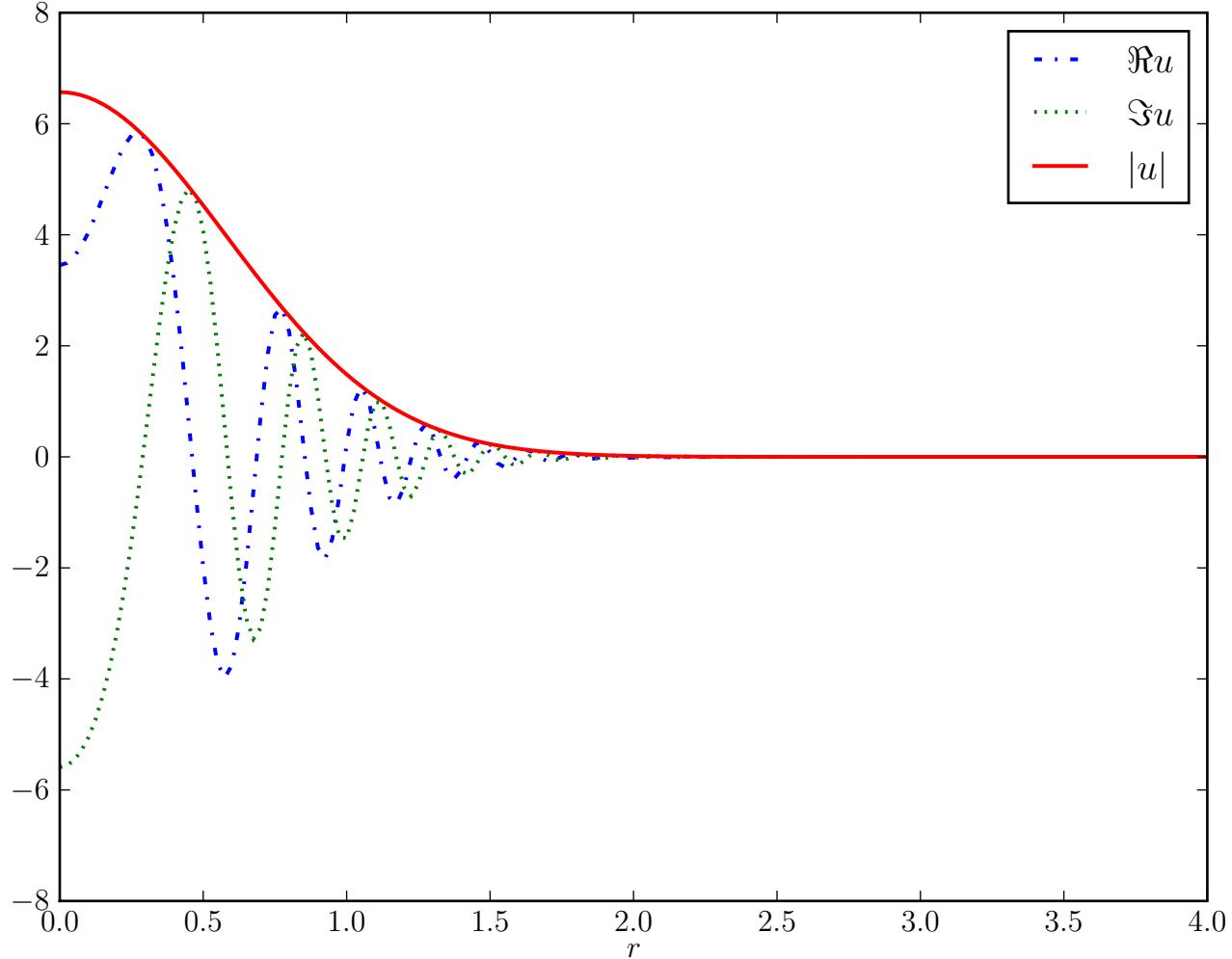
$t = 0.035$



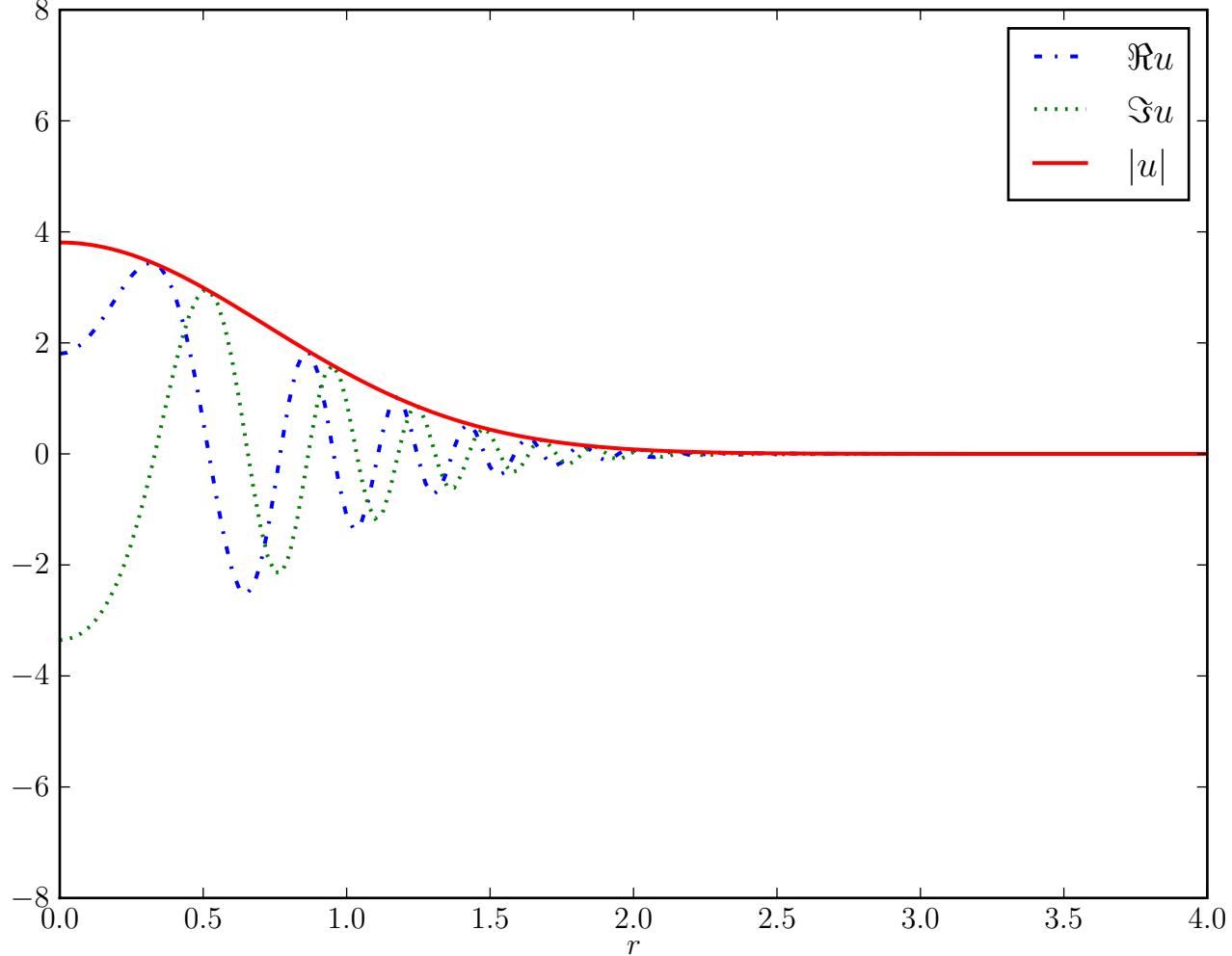
$t = 0.04$



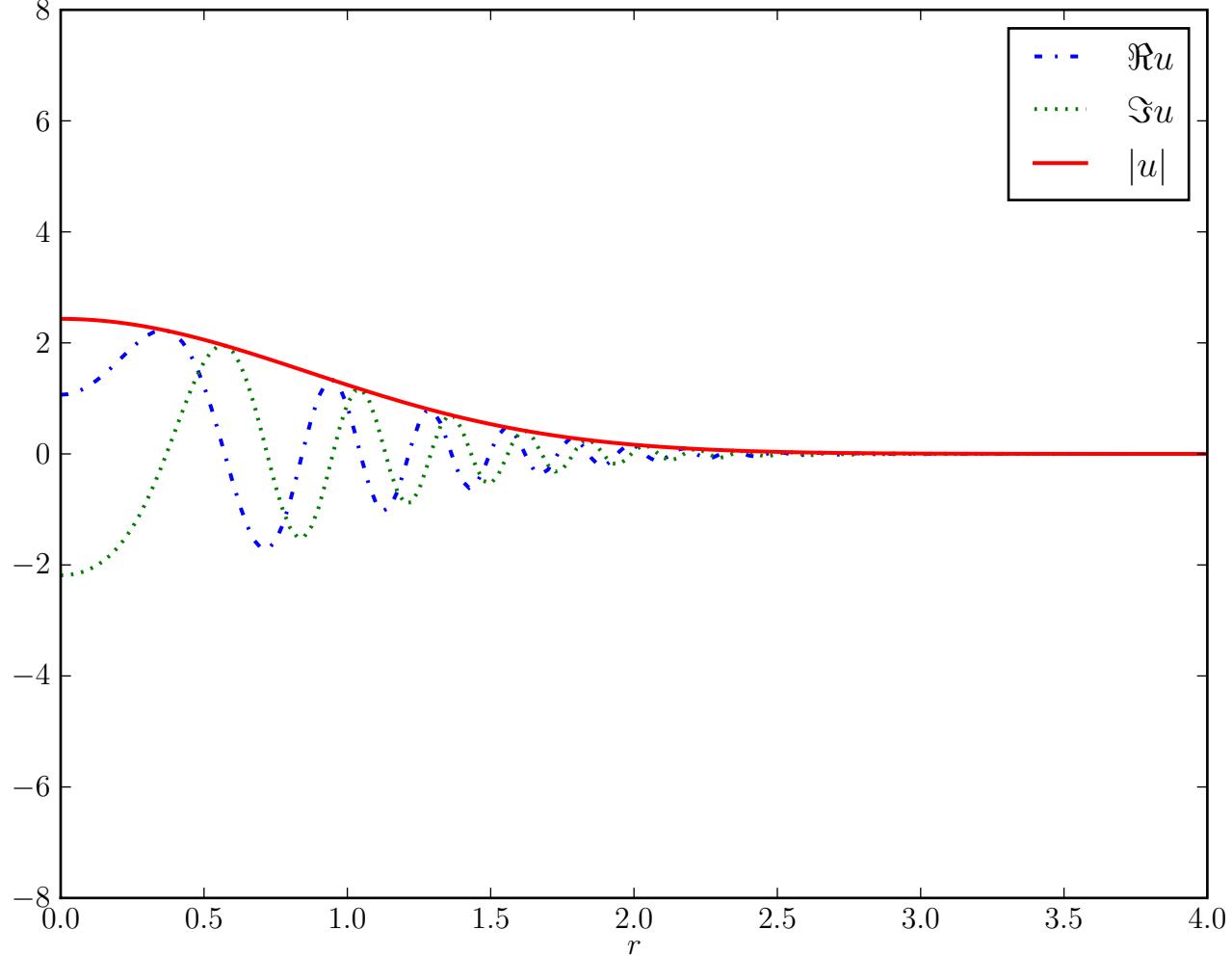
$t = 0.045$



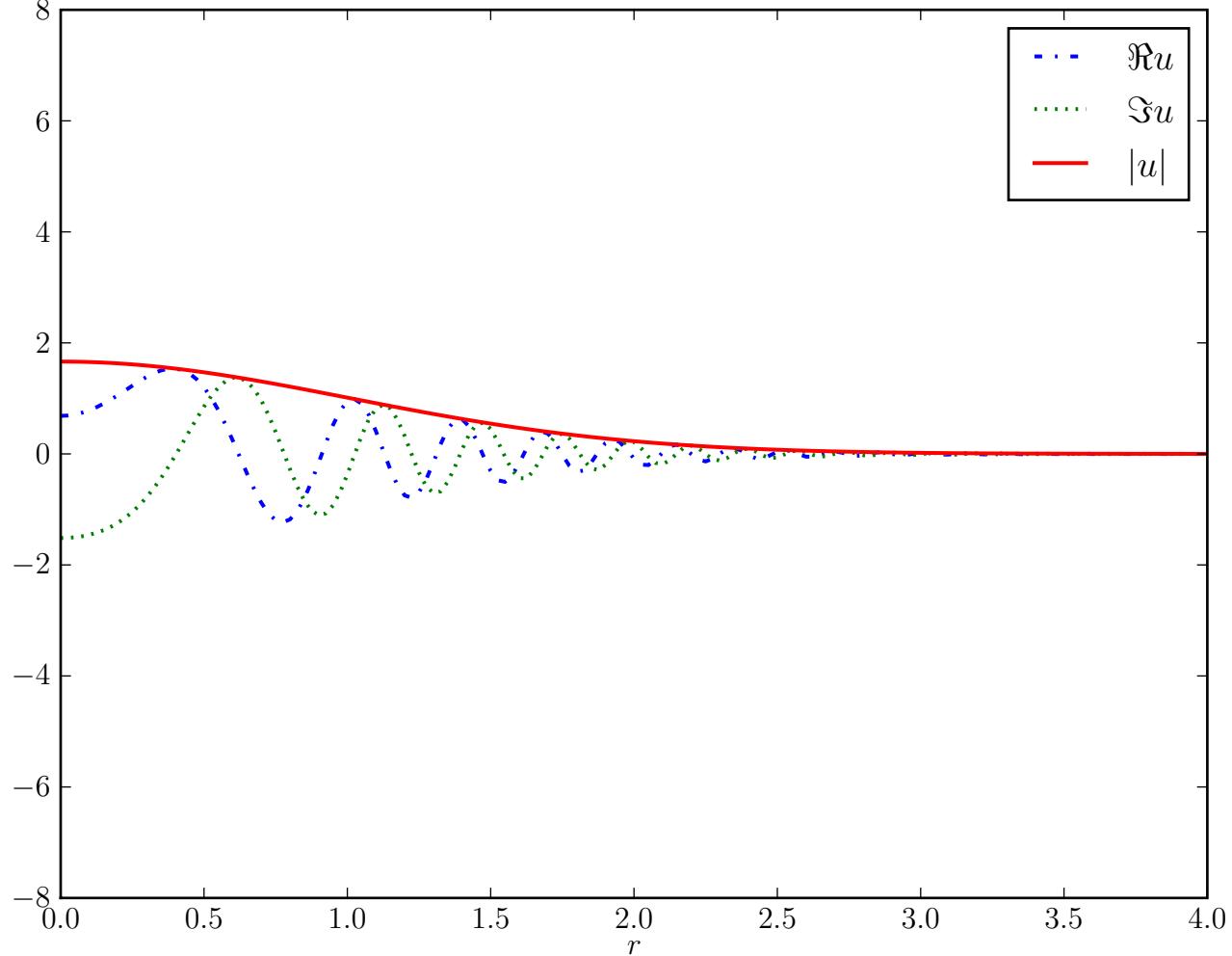
$t = 0.05$



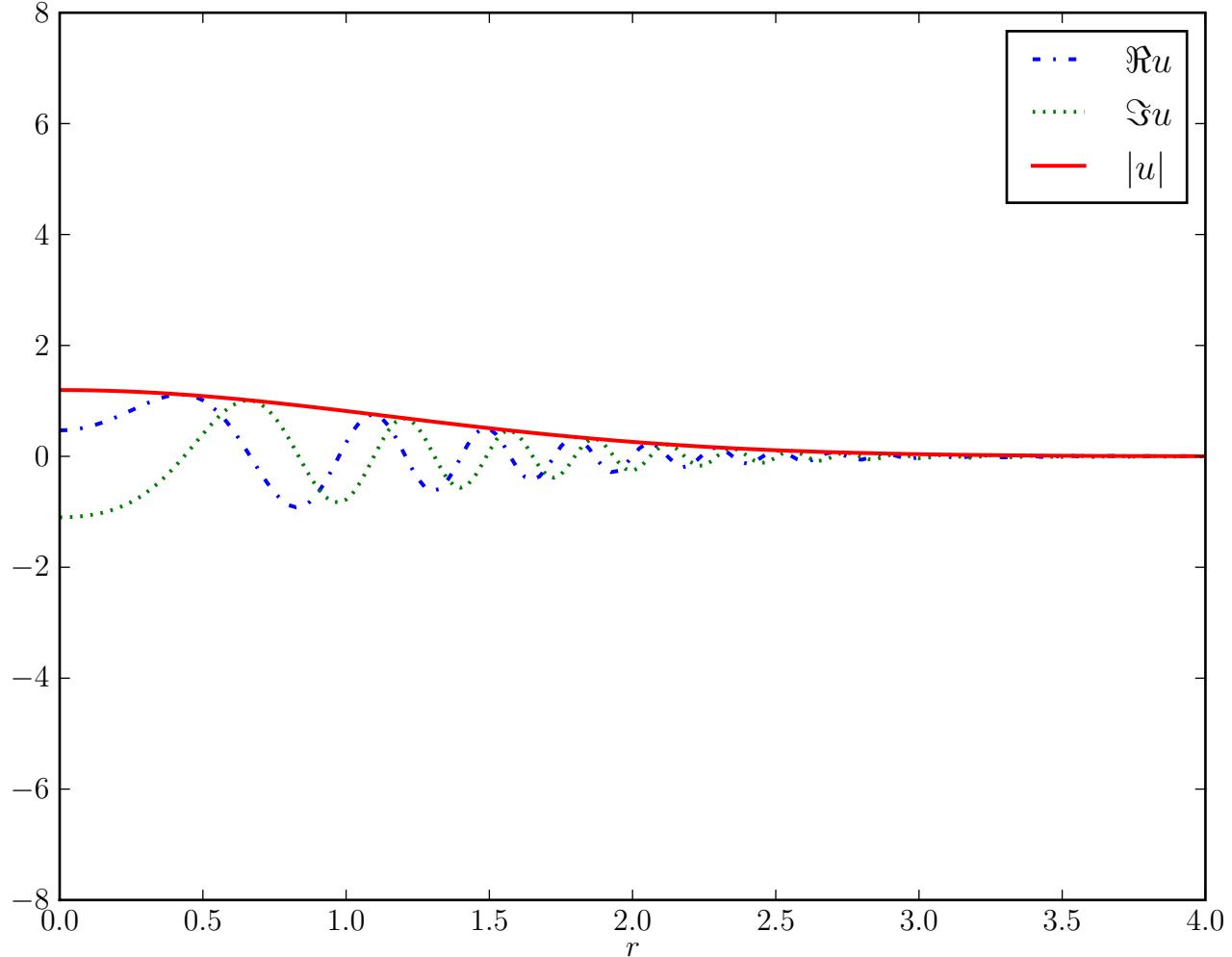
$t = 0.055$



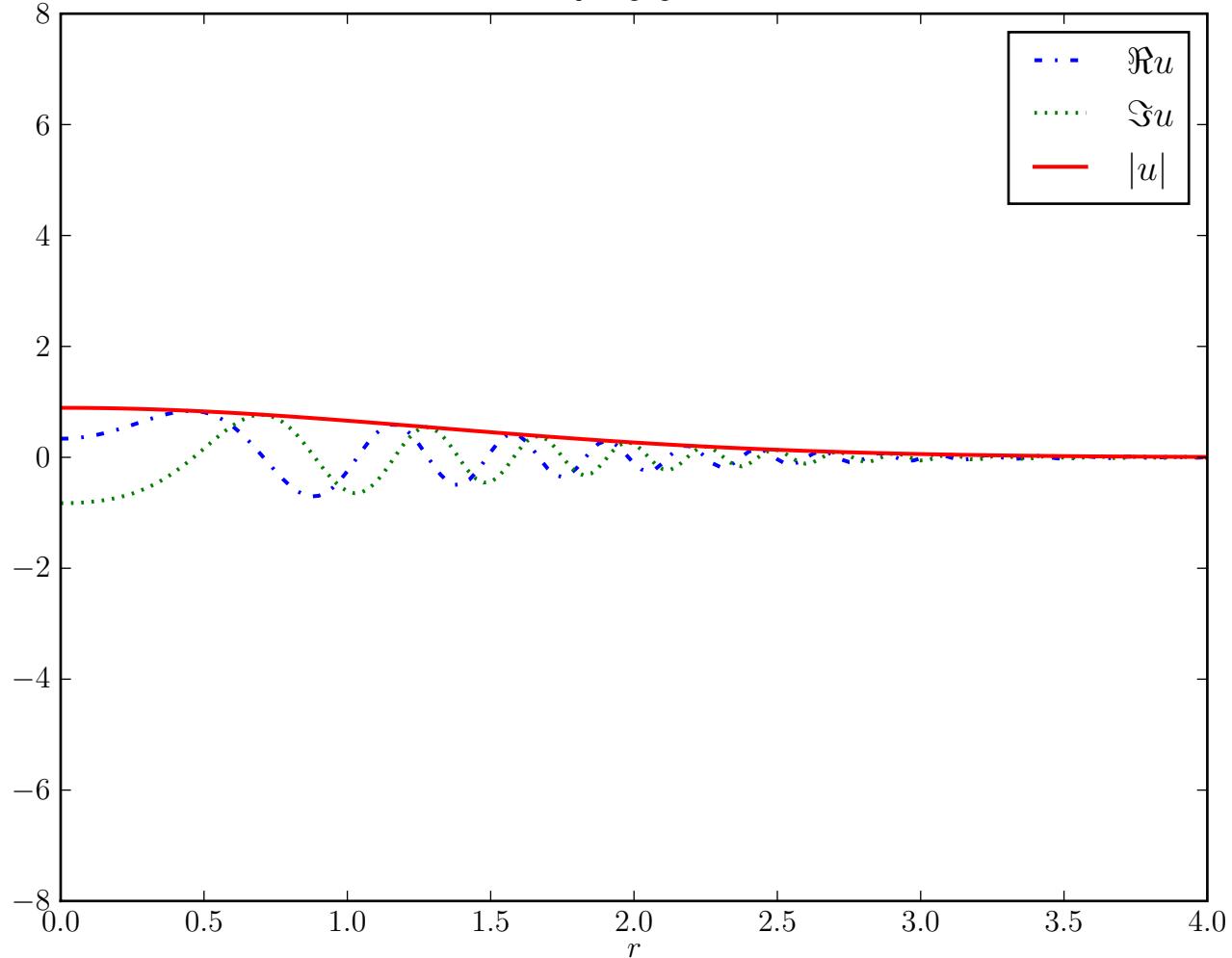
$t = 0.06$

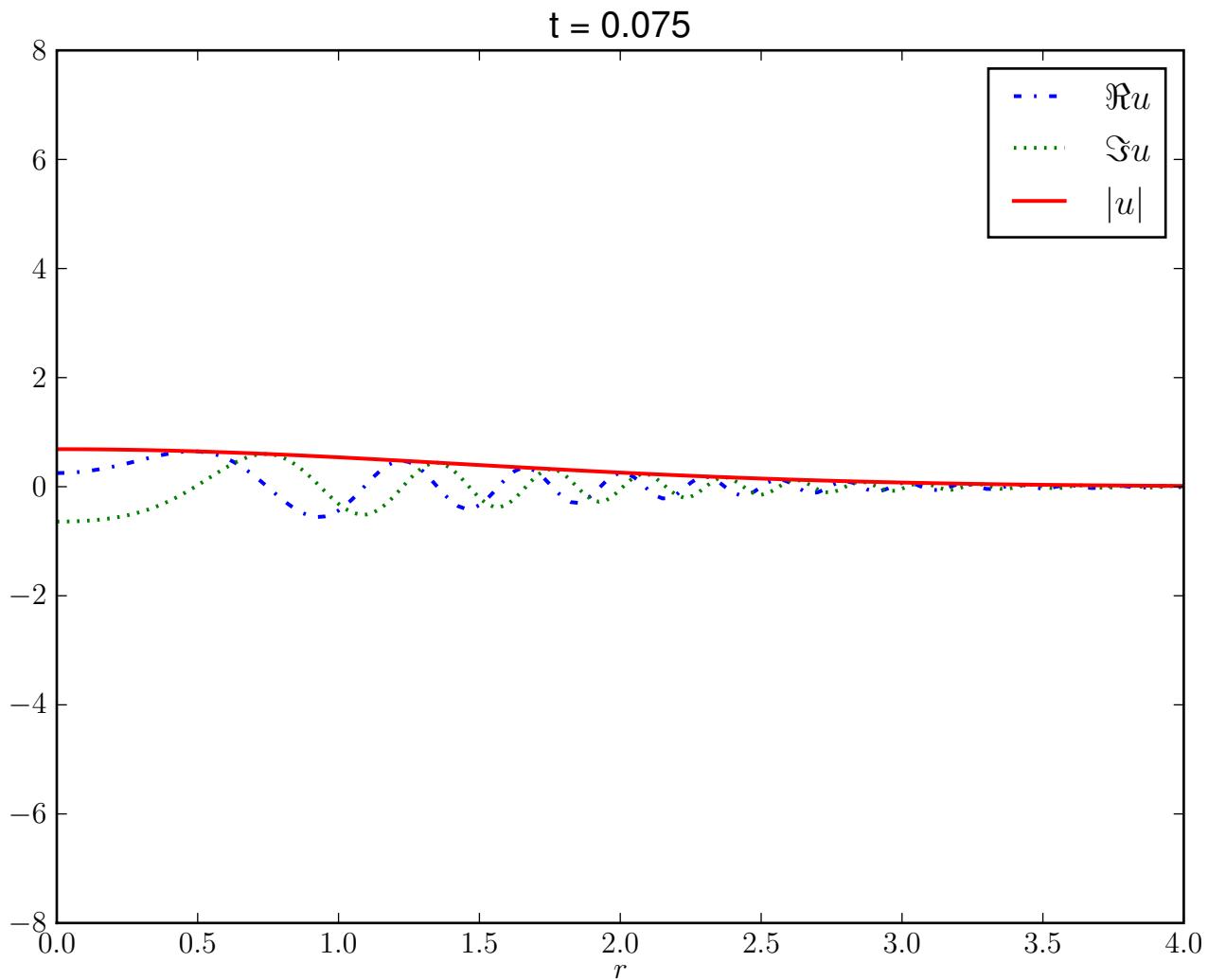


$t = 0.065$

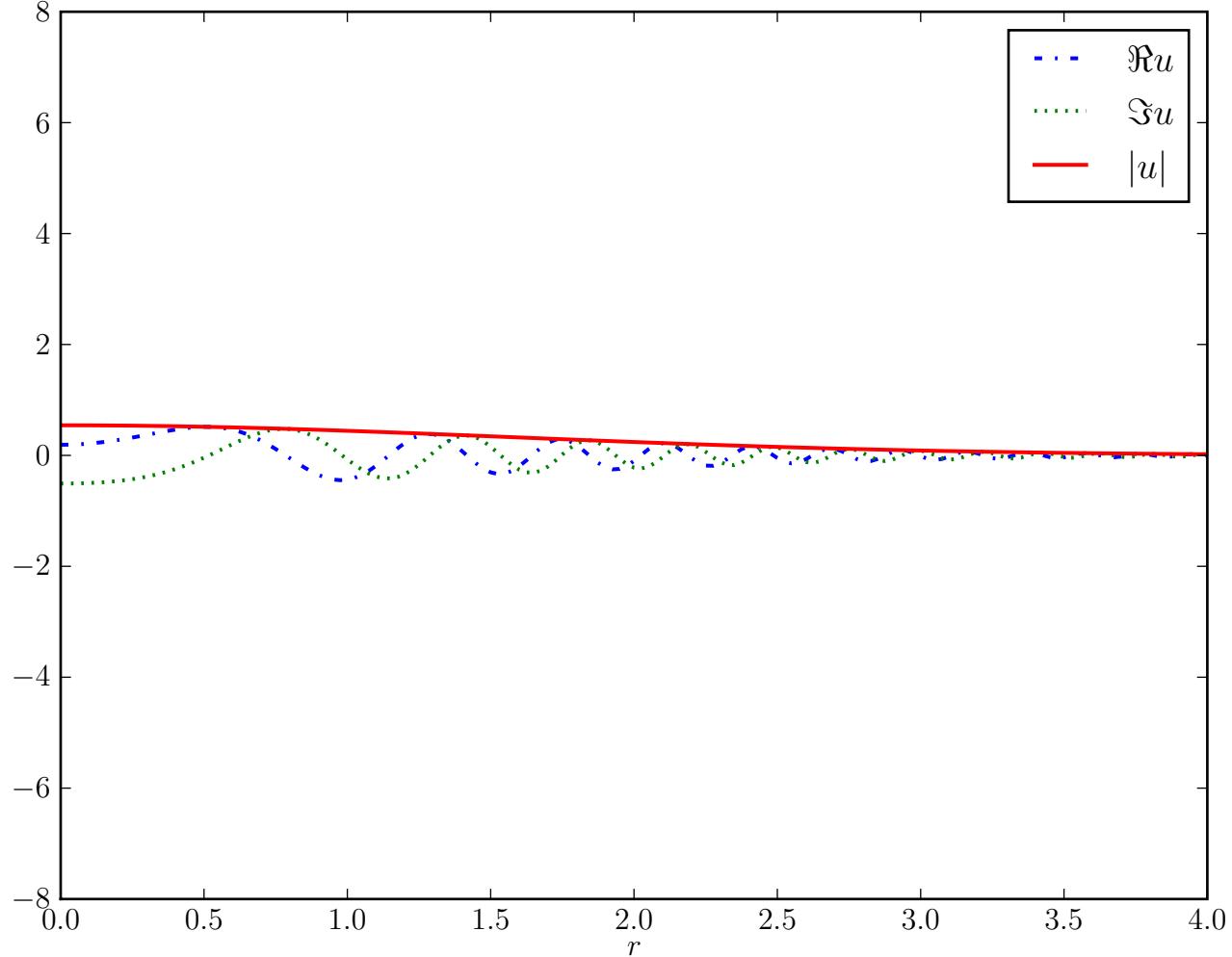


$t = 0.07$

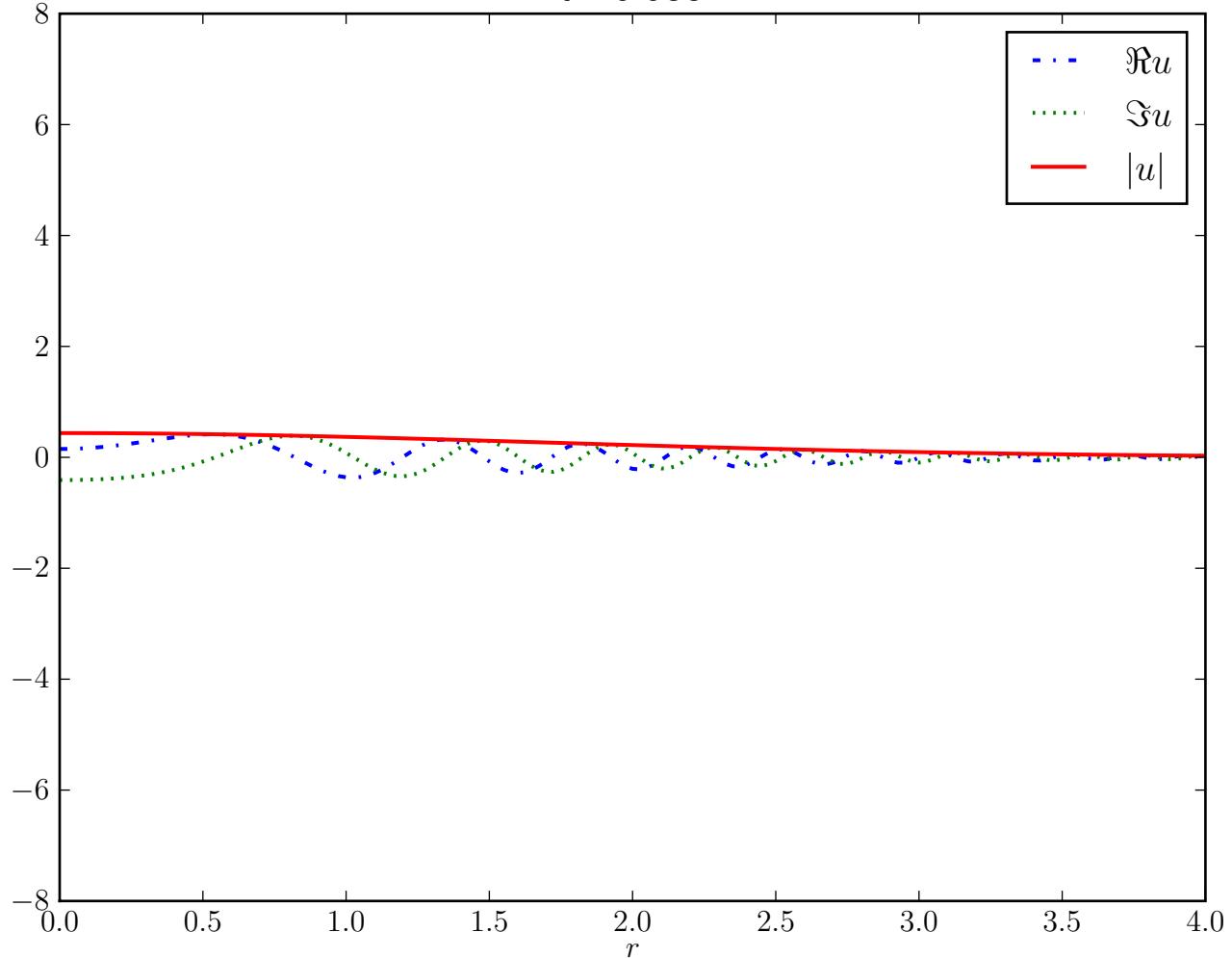


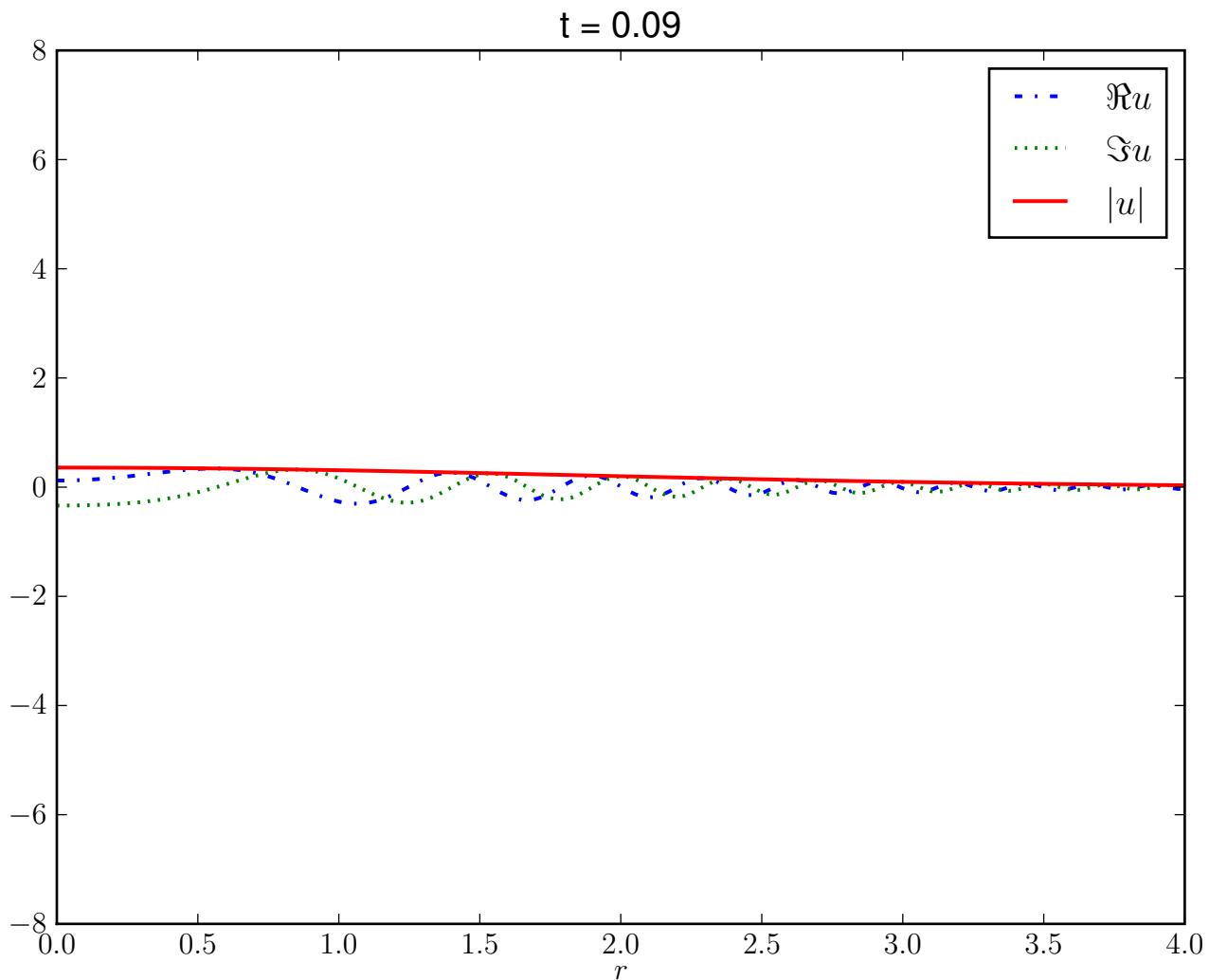


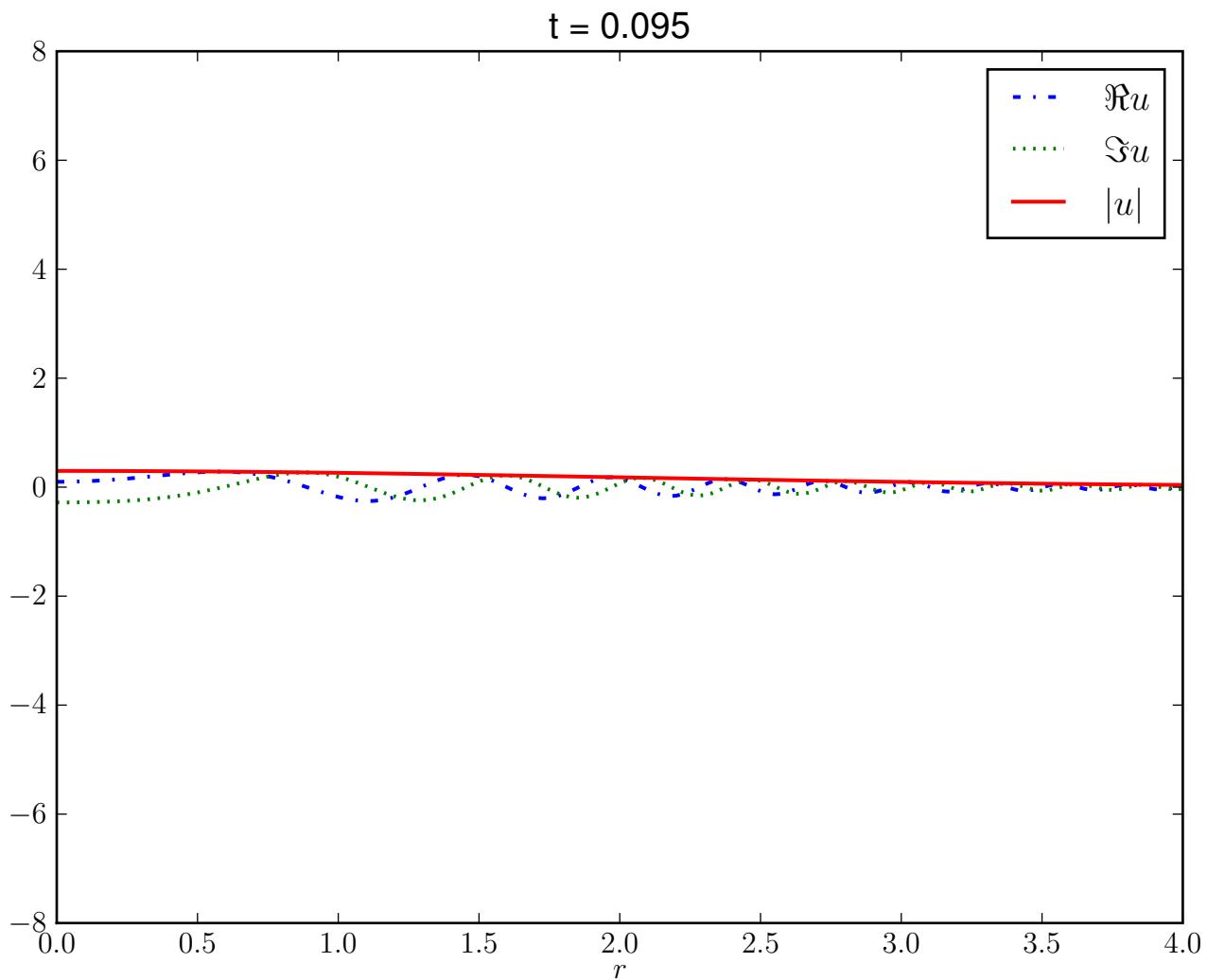
$t = 0.08$

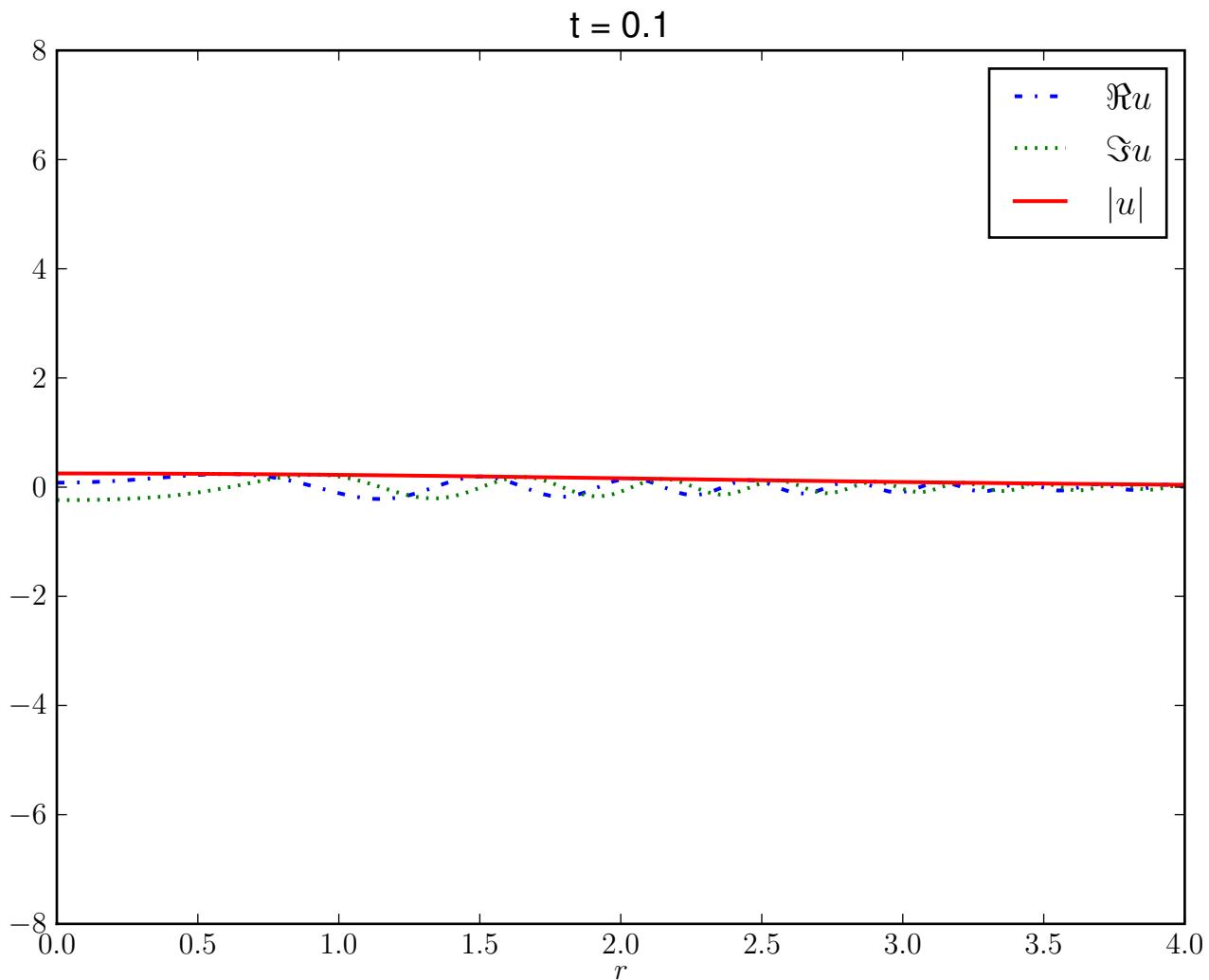


$t = 0.085$

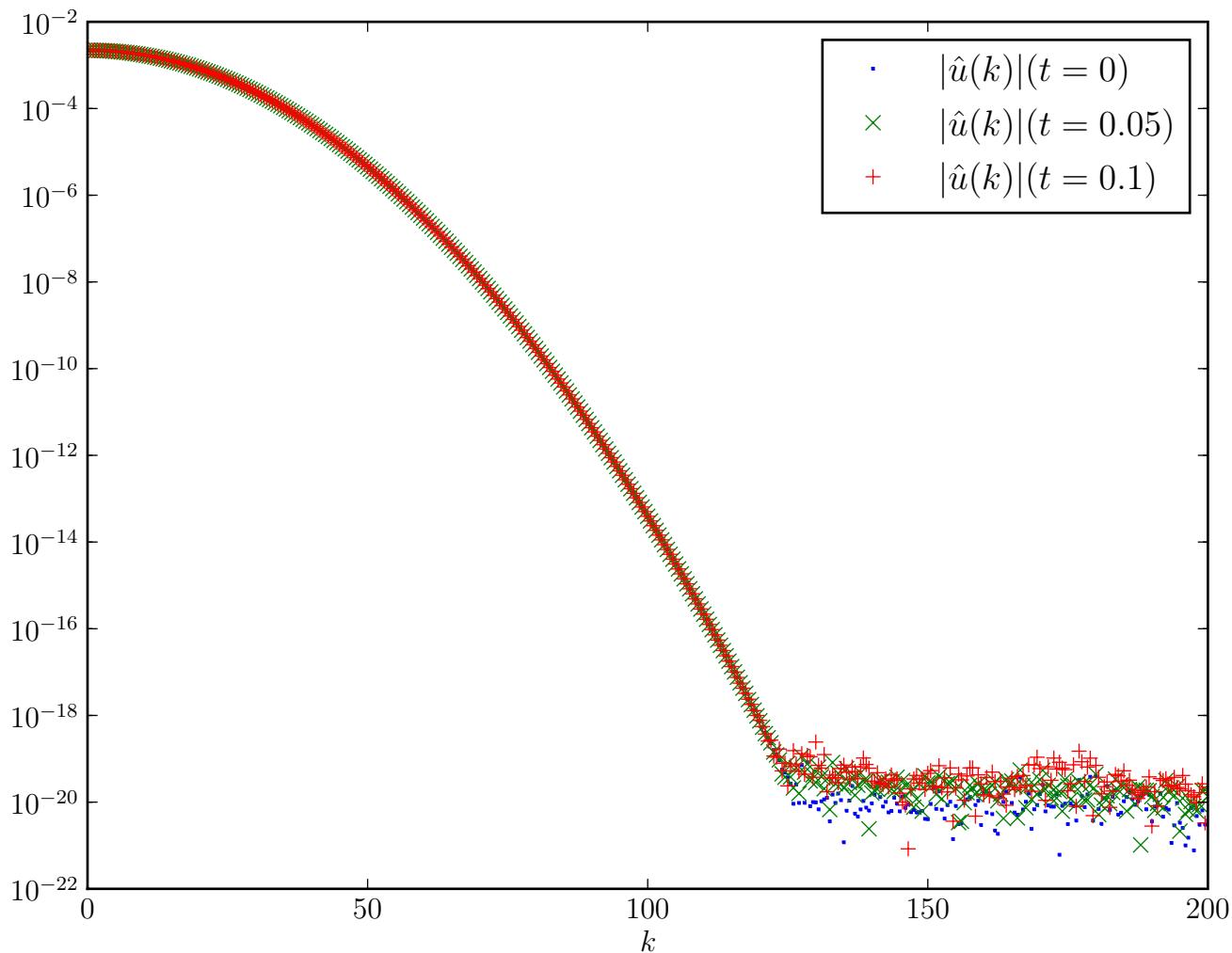




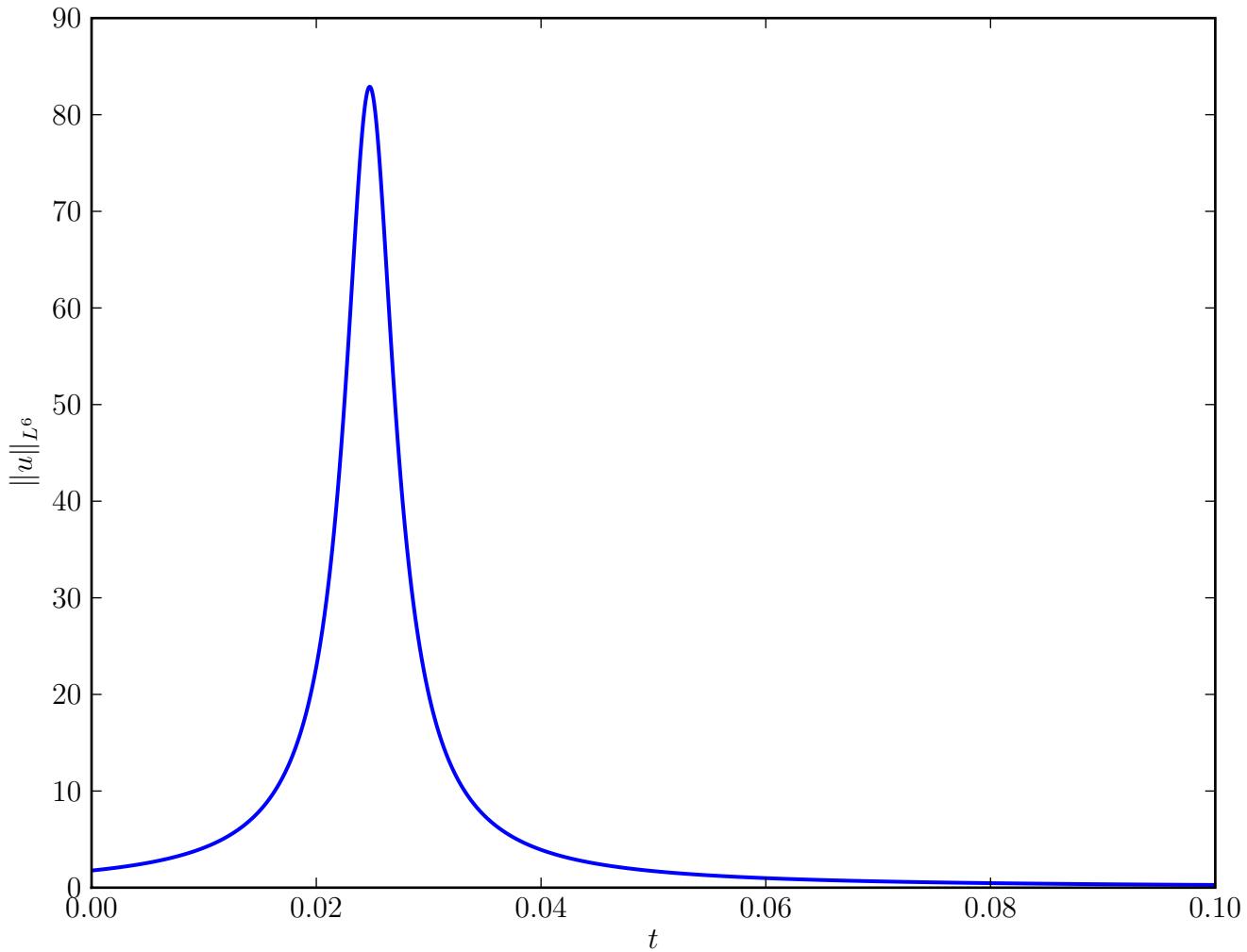




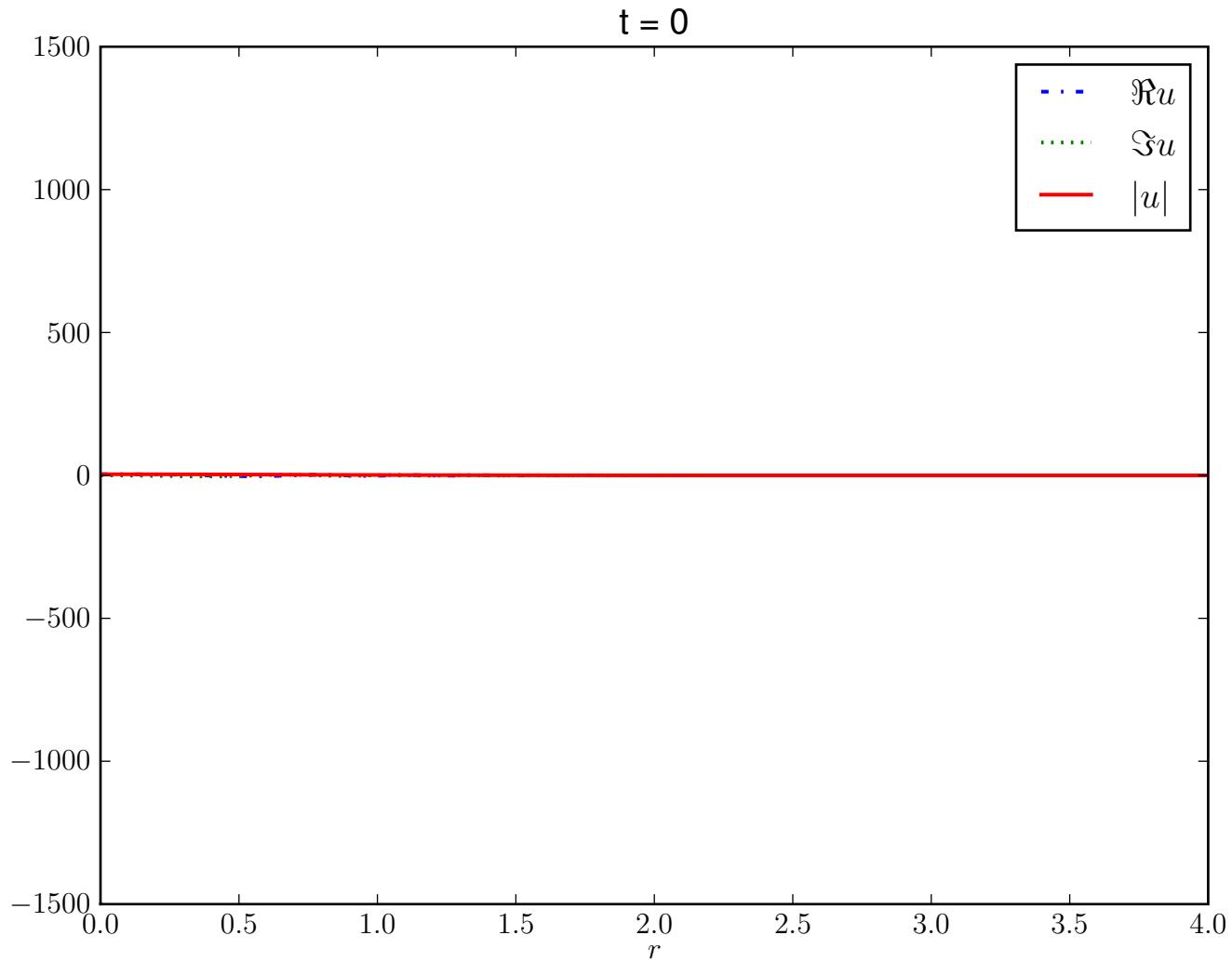
Phased Centered Gaussian Fourier transform snapshots along linear flow



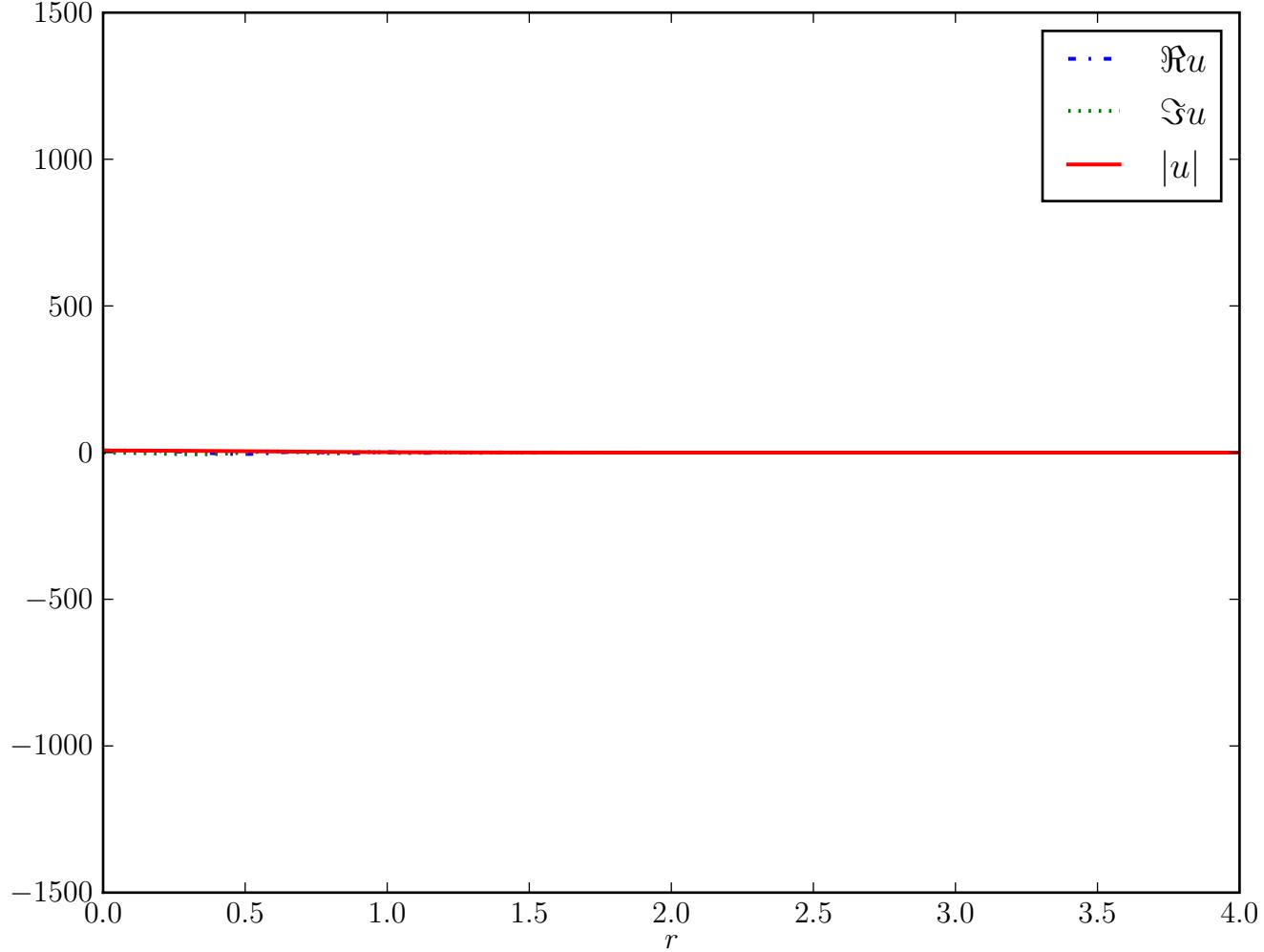
Potential Energy Norm under linear flow



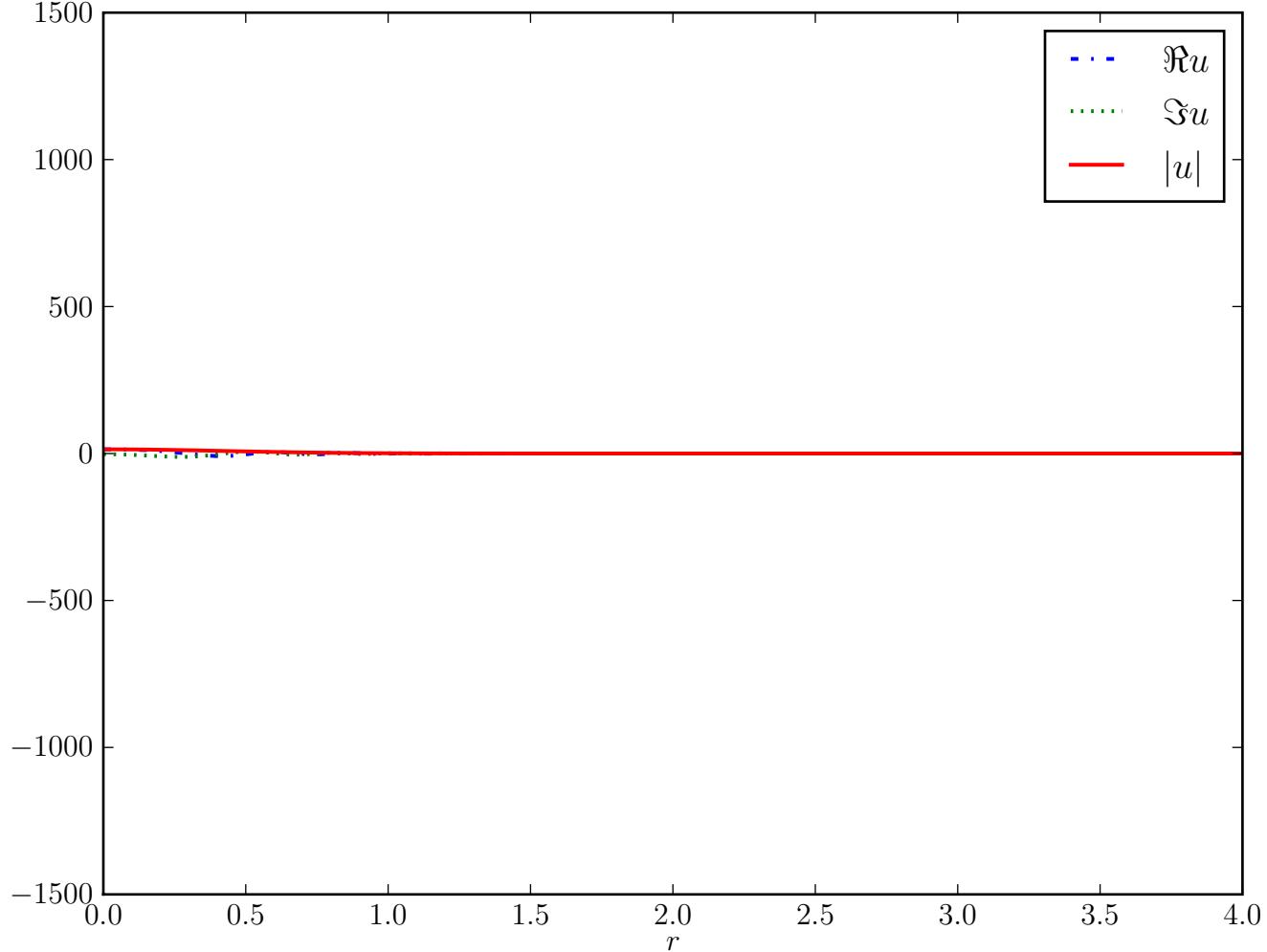
Phased Centered Gaussian under linear flow; bigger vertical axis



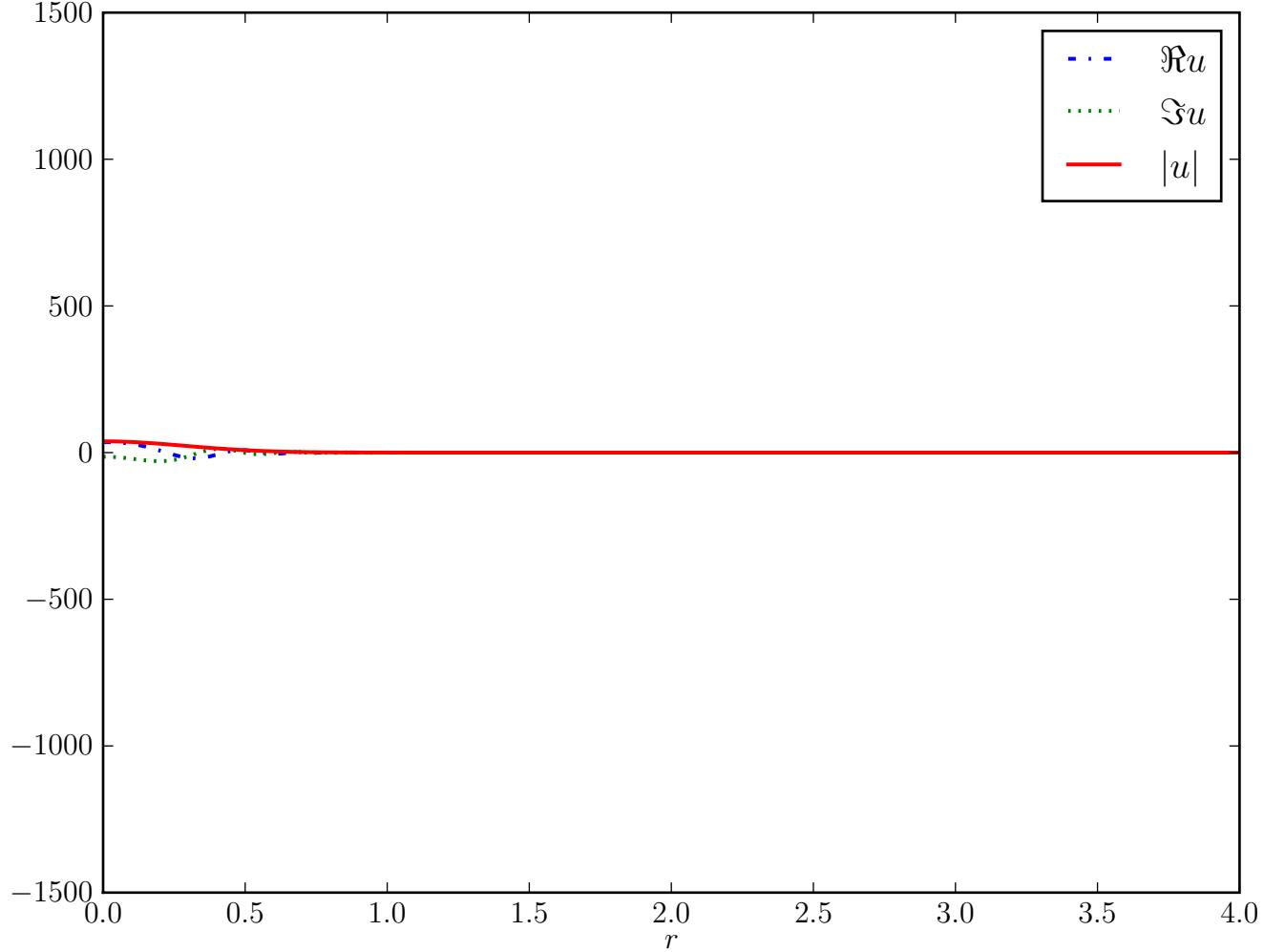
$t = 0.005$



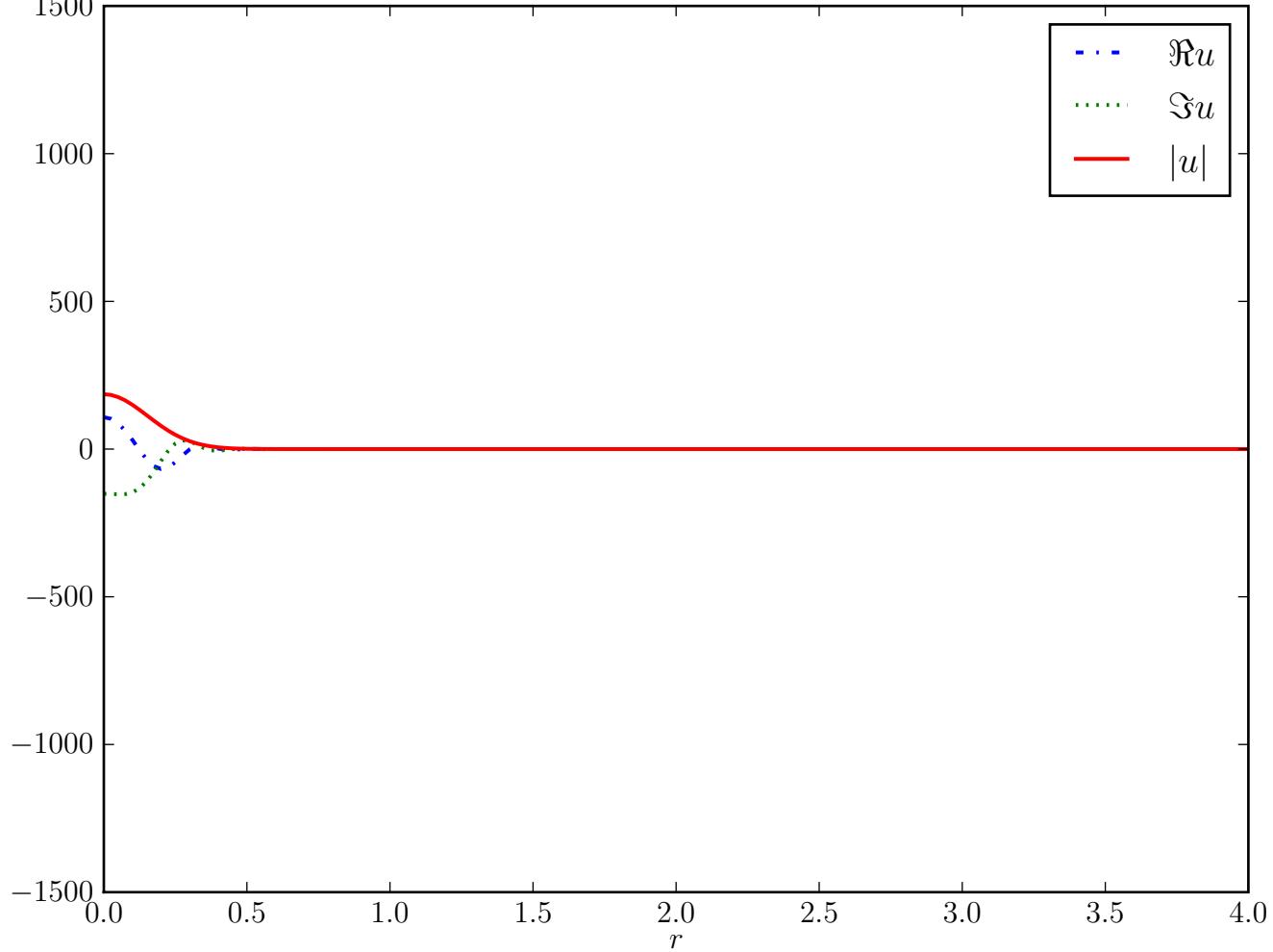
$t = 0.01$



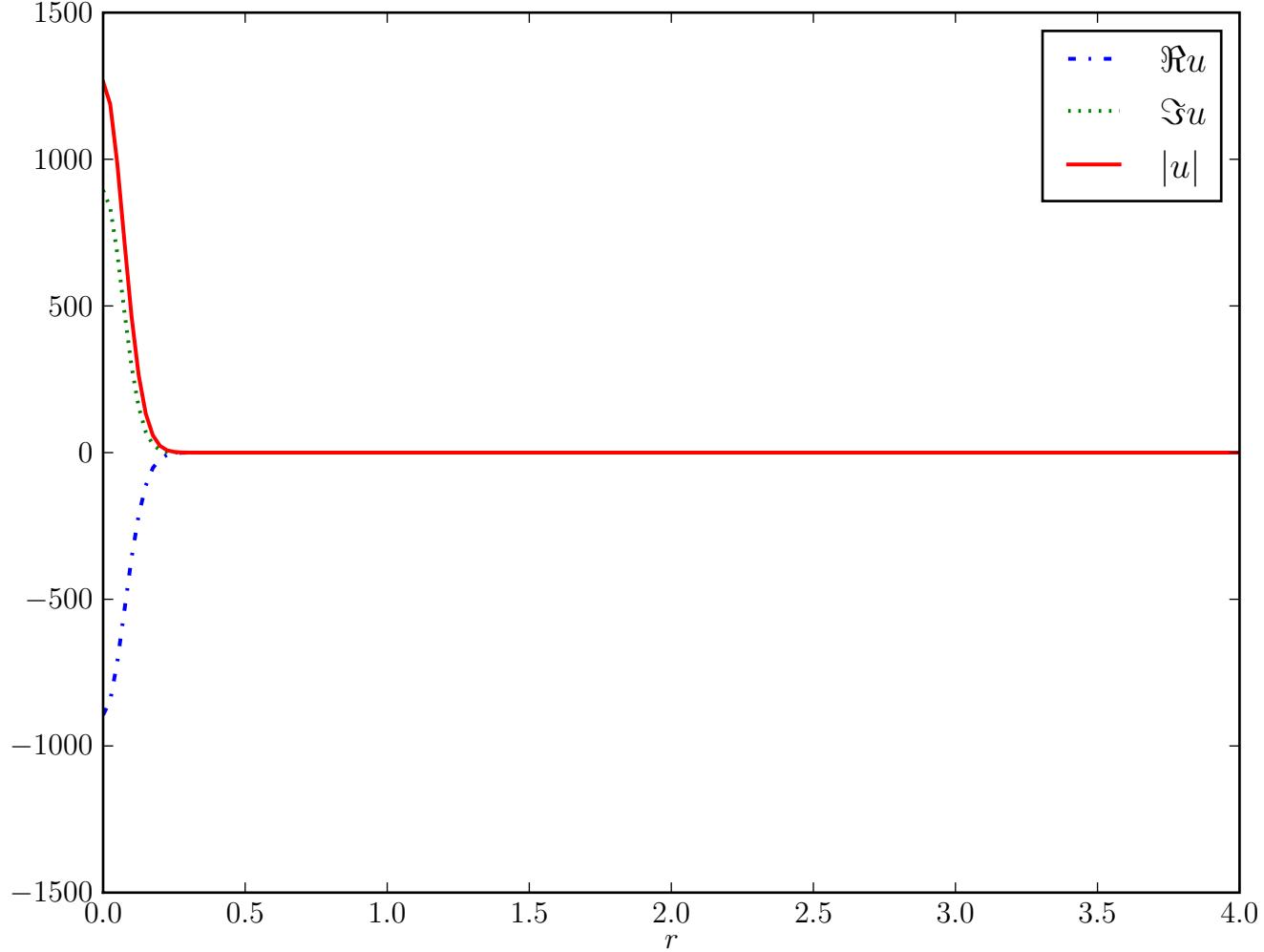
$t = 0.015$



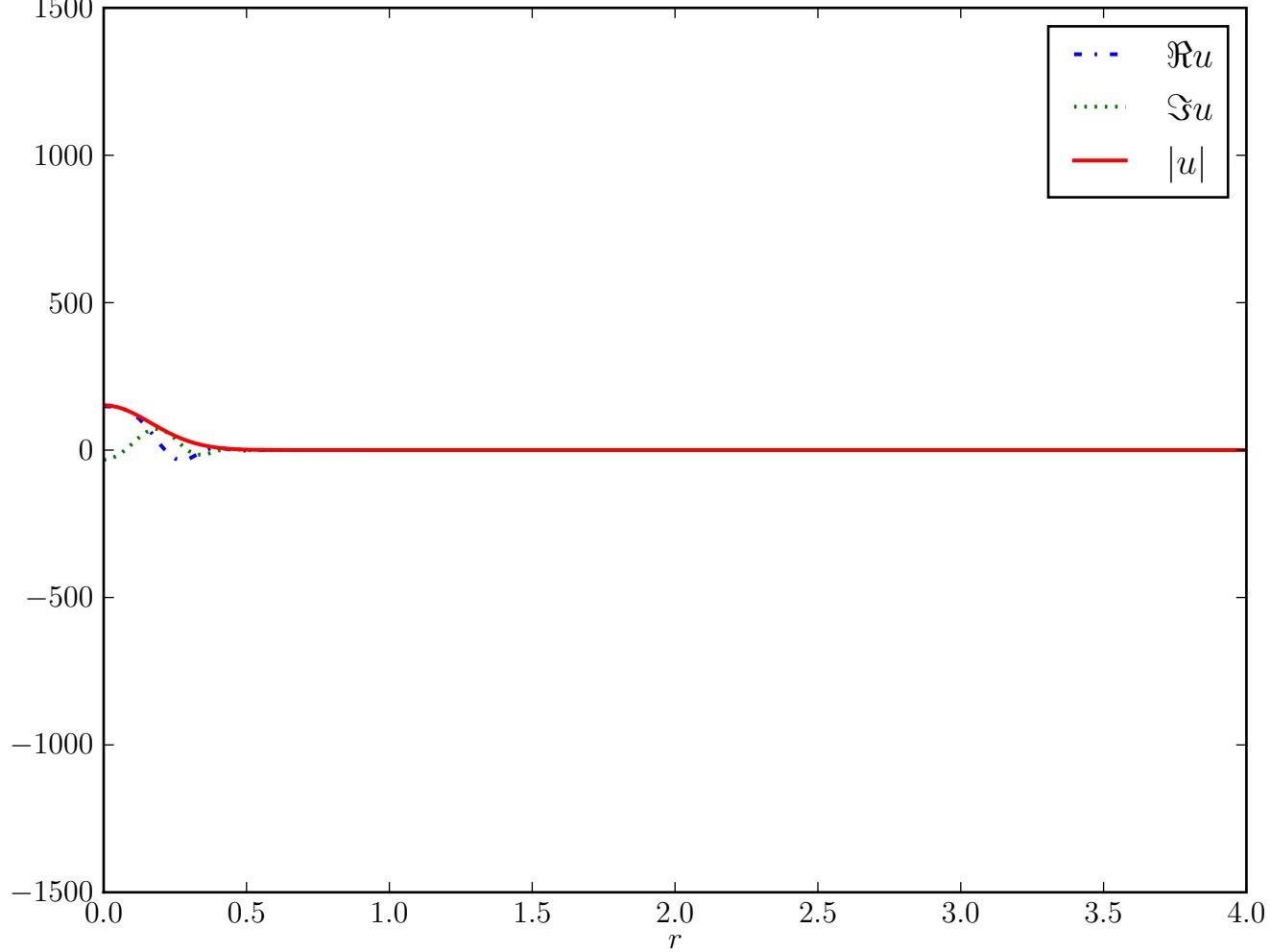
$t = 0.02$



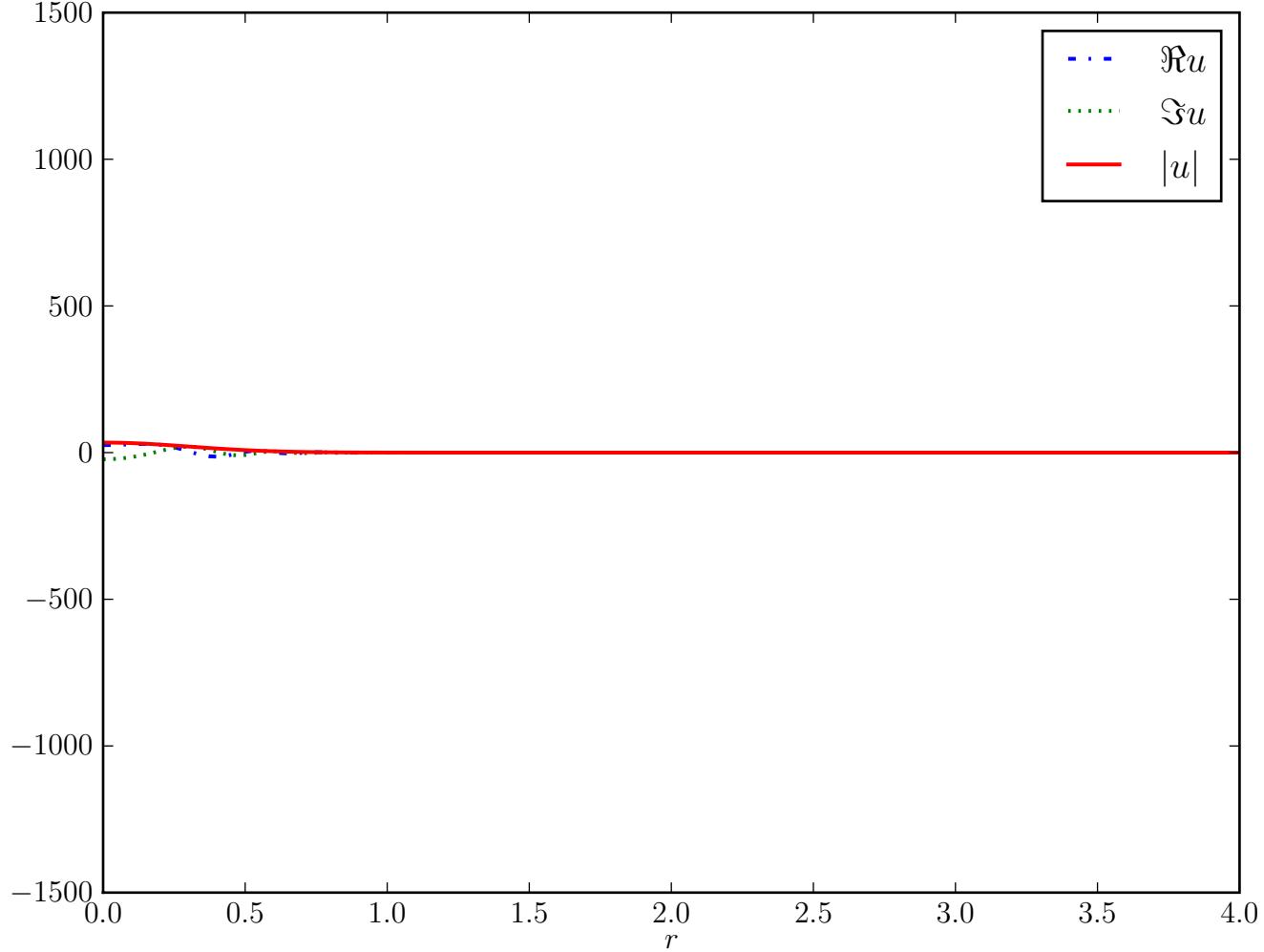
$t = 0.025$



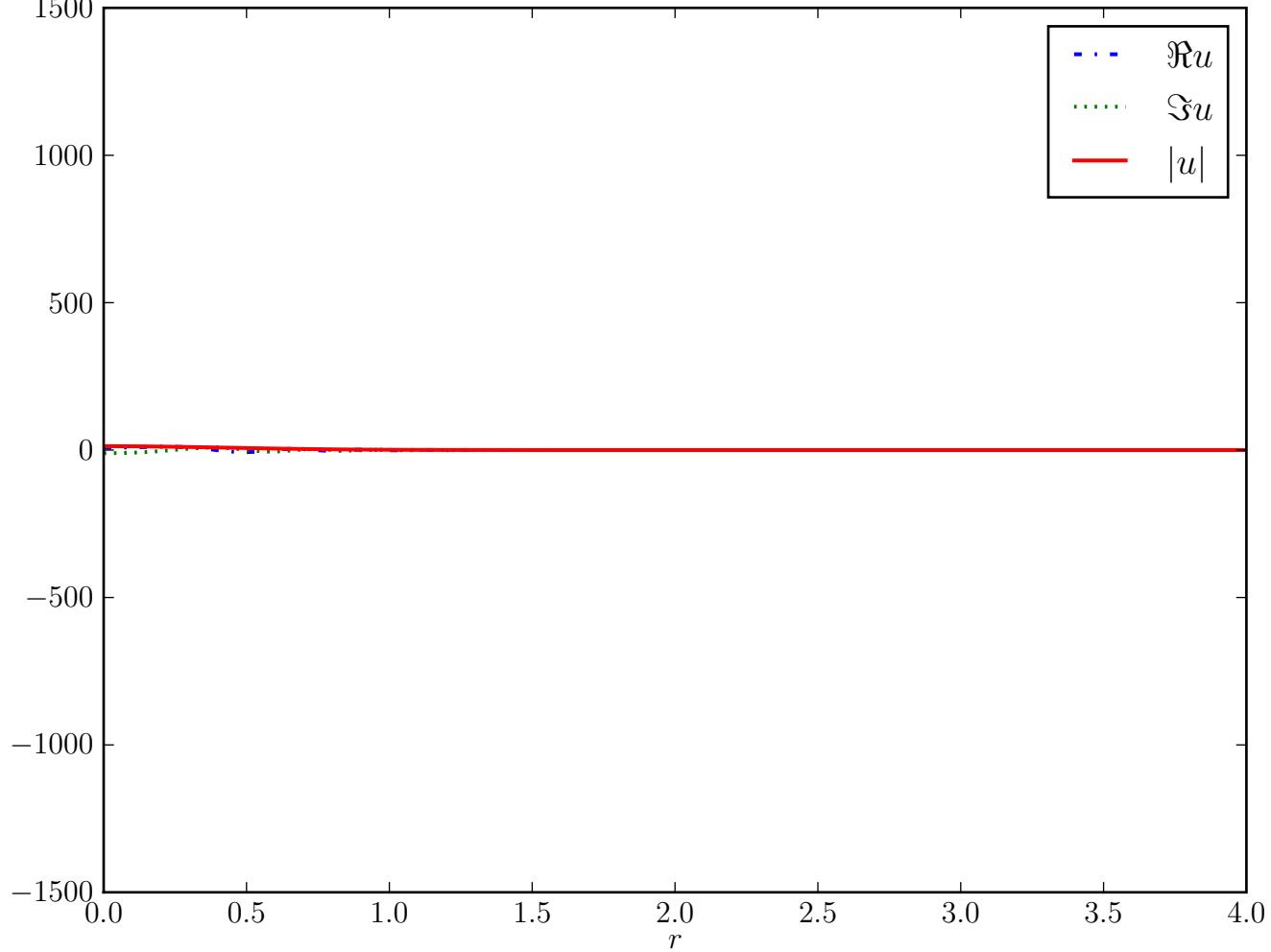
$t = 0.03$



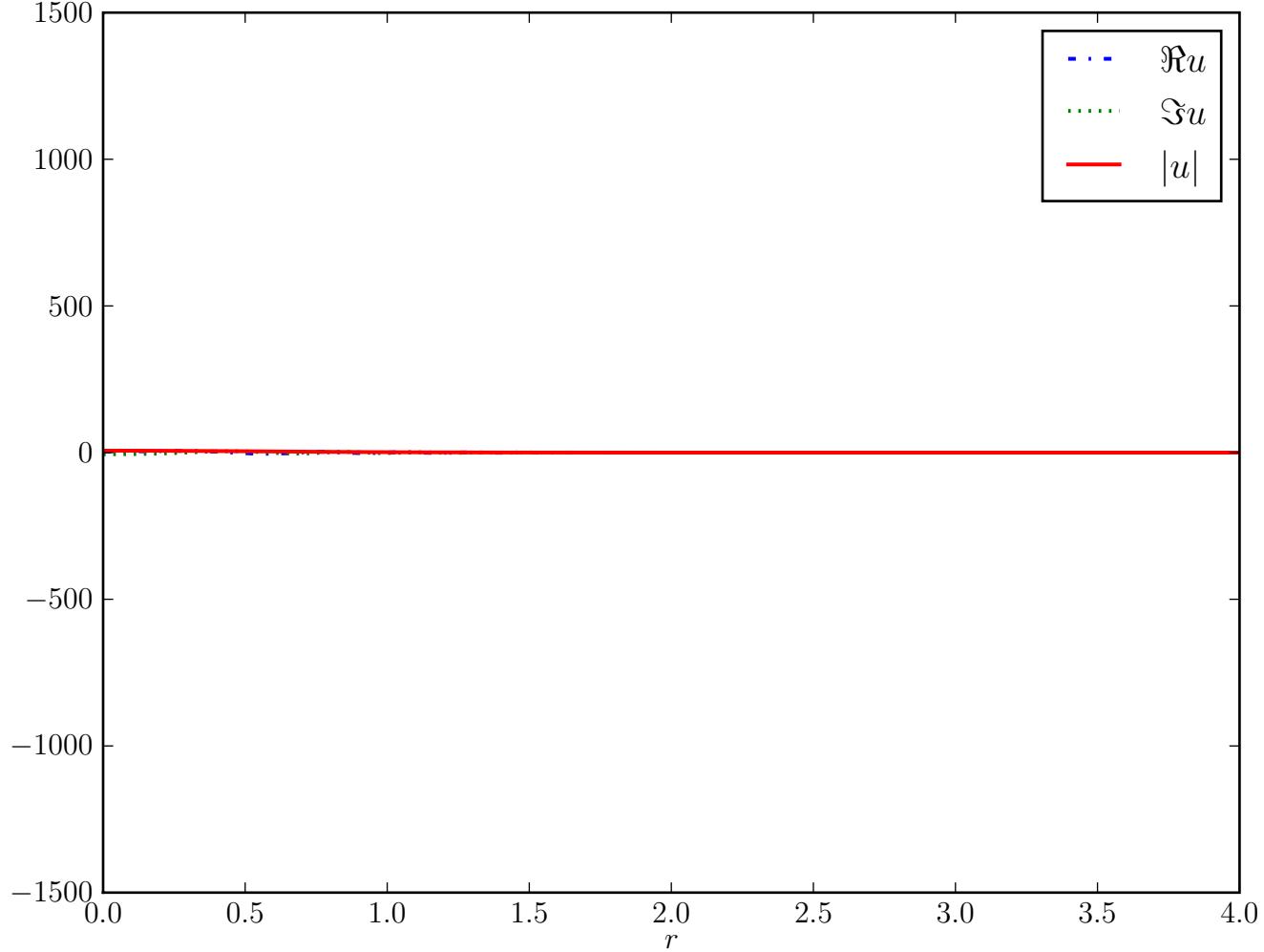
$t = 0.035$



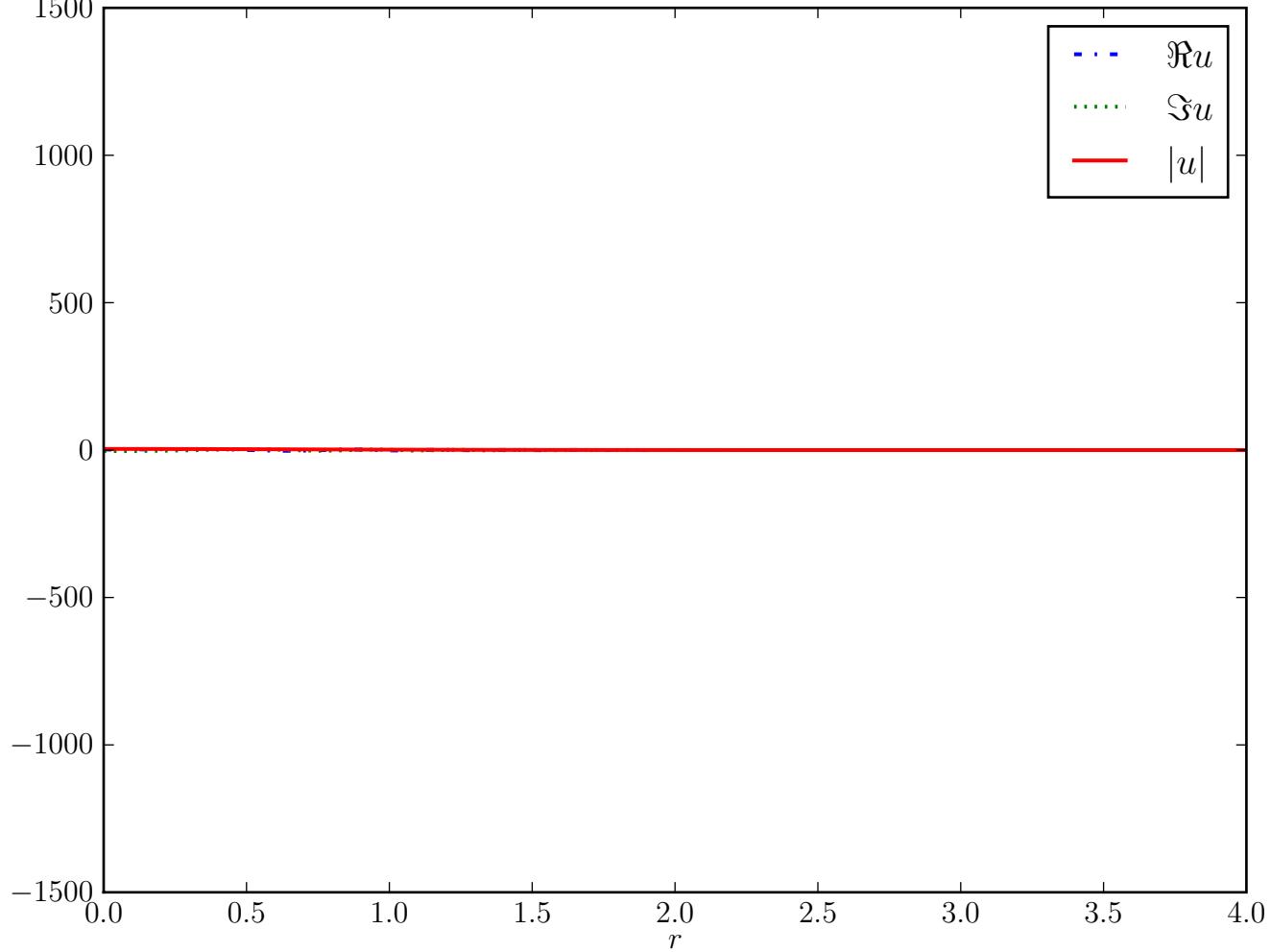
$t = 0.04$



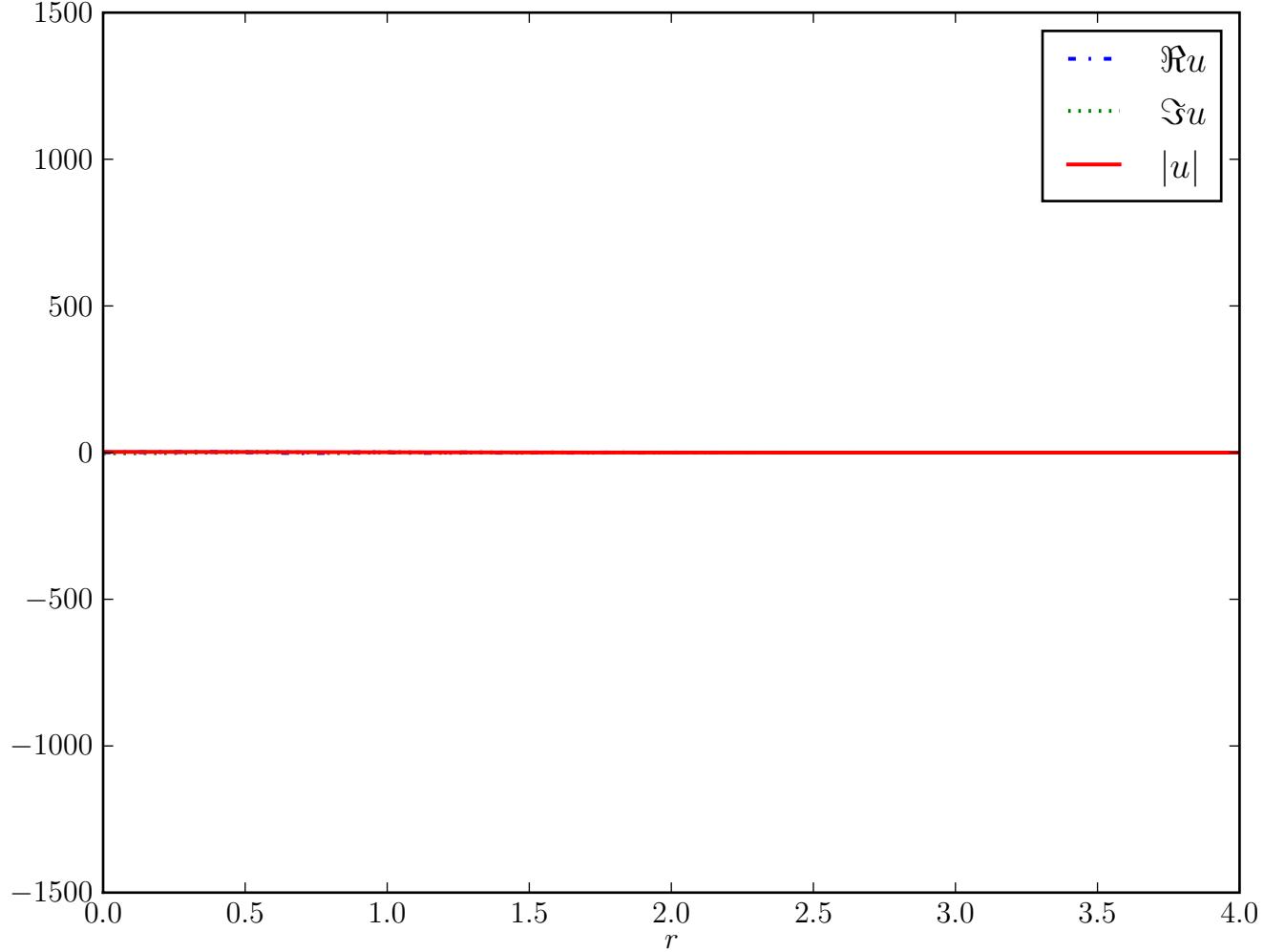
$t = 0.045$



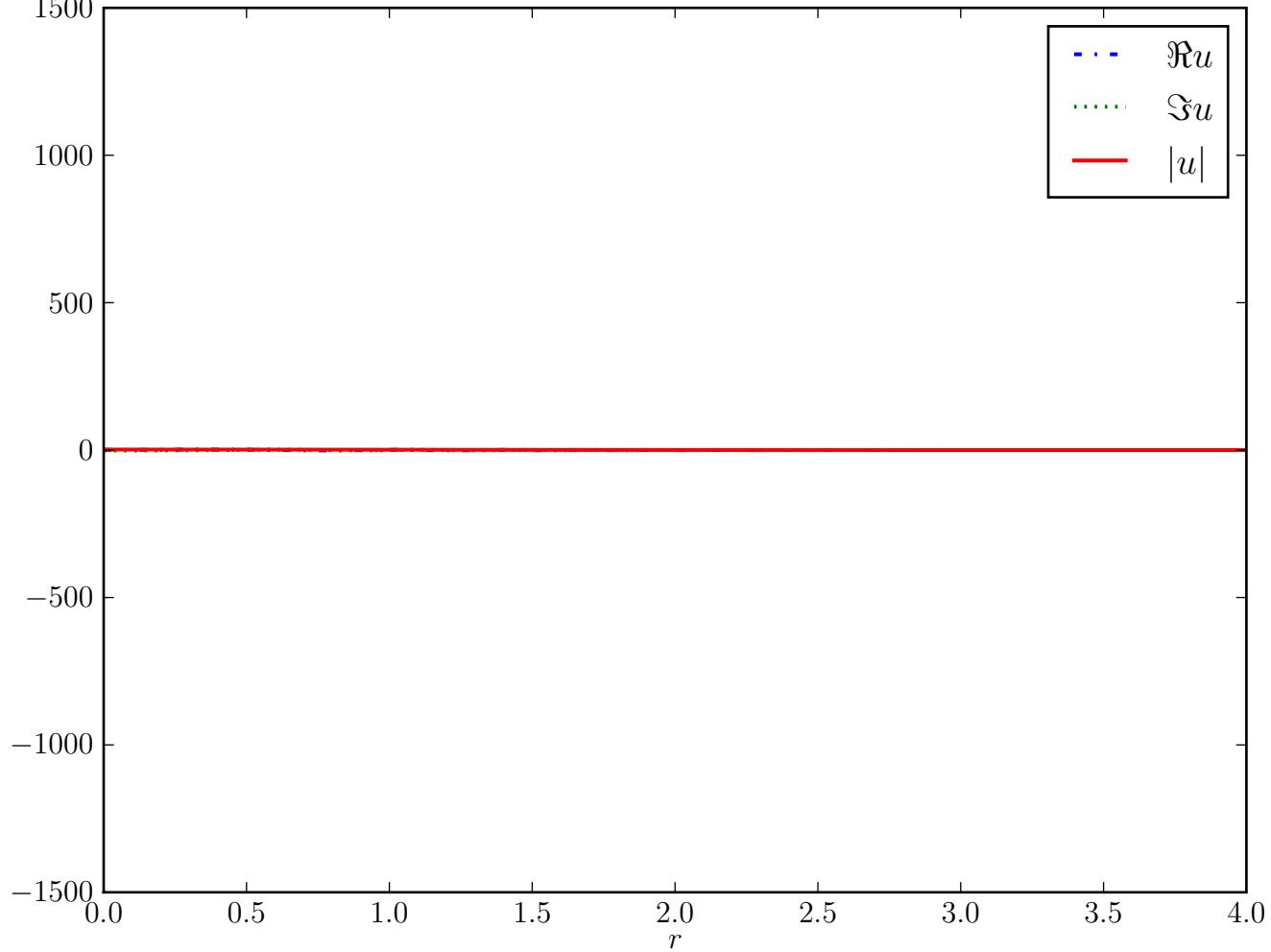
$t = 0.05$



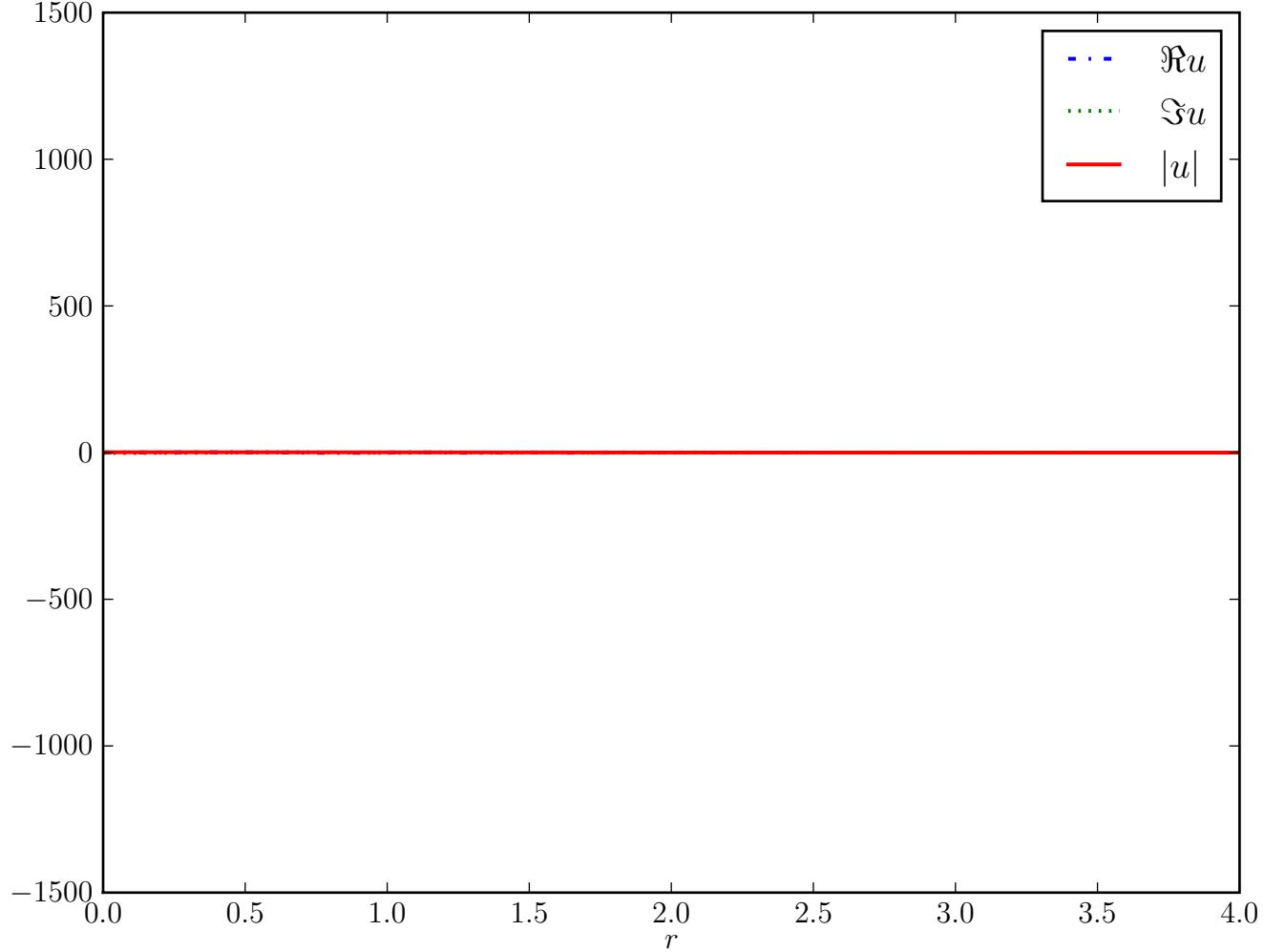
$t = 0.055$



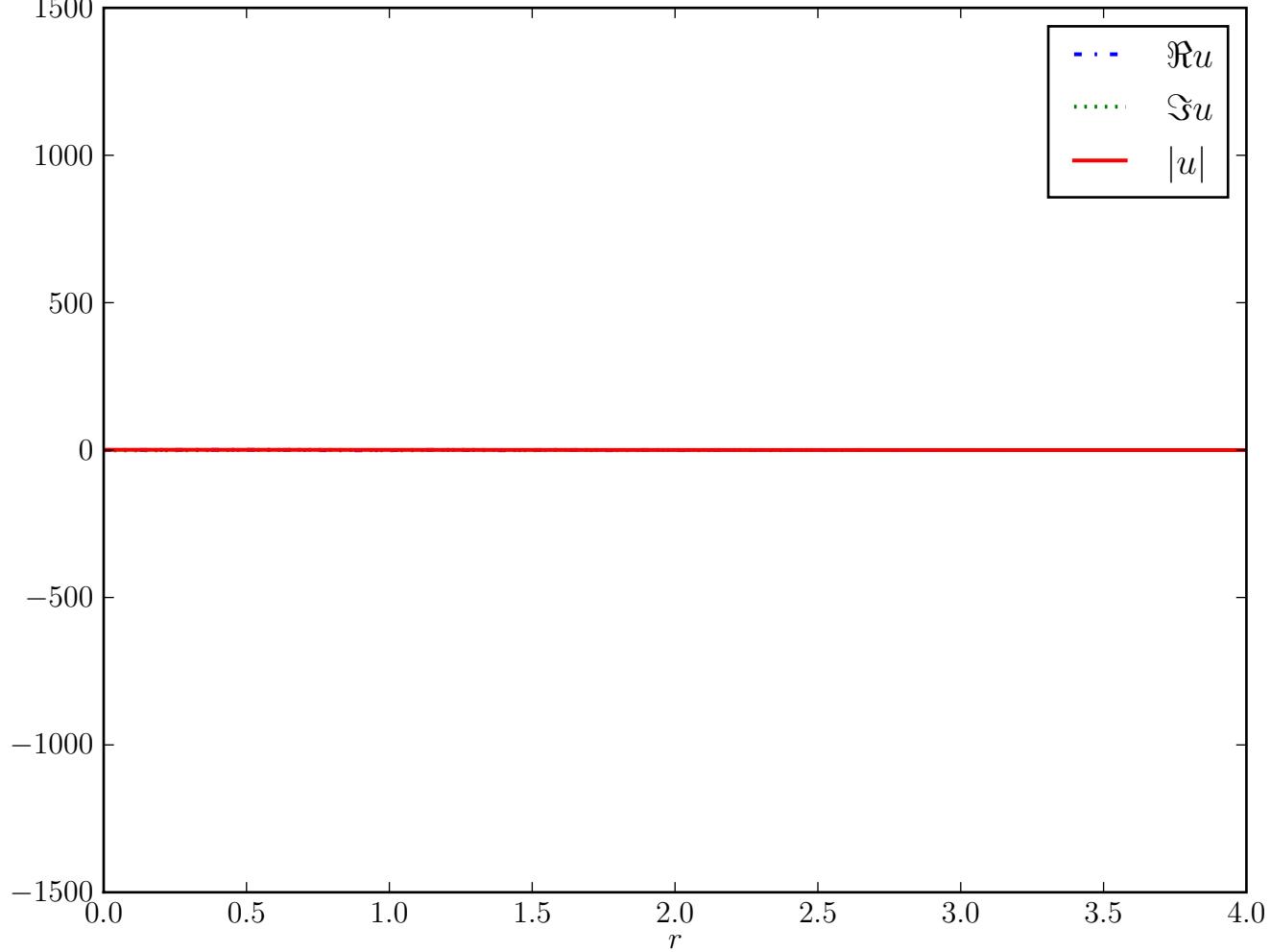
$t = 0.06$



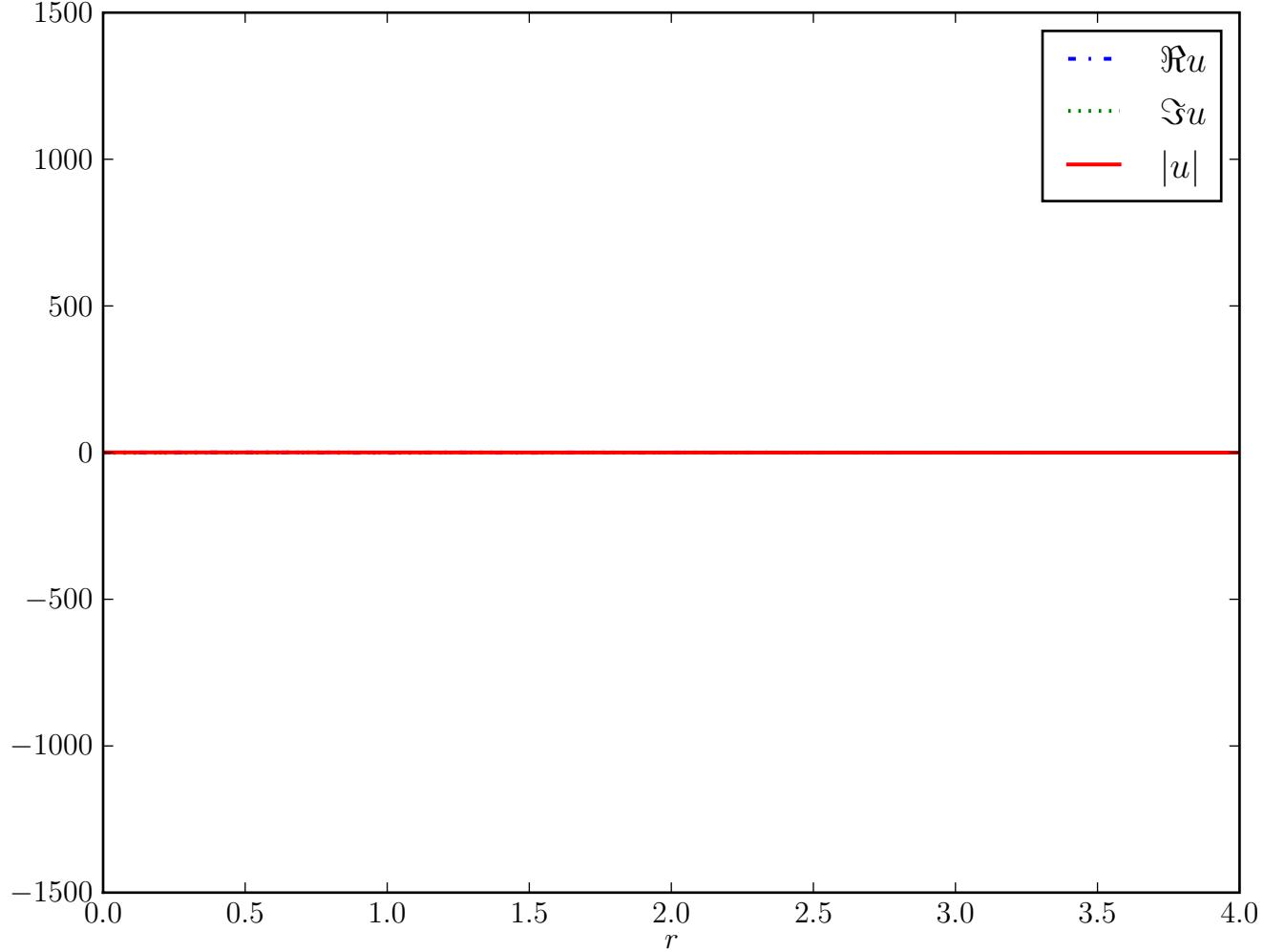
$t = 0.065$



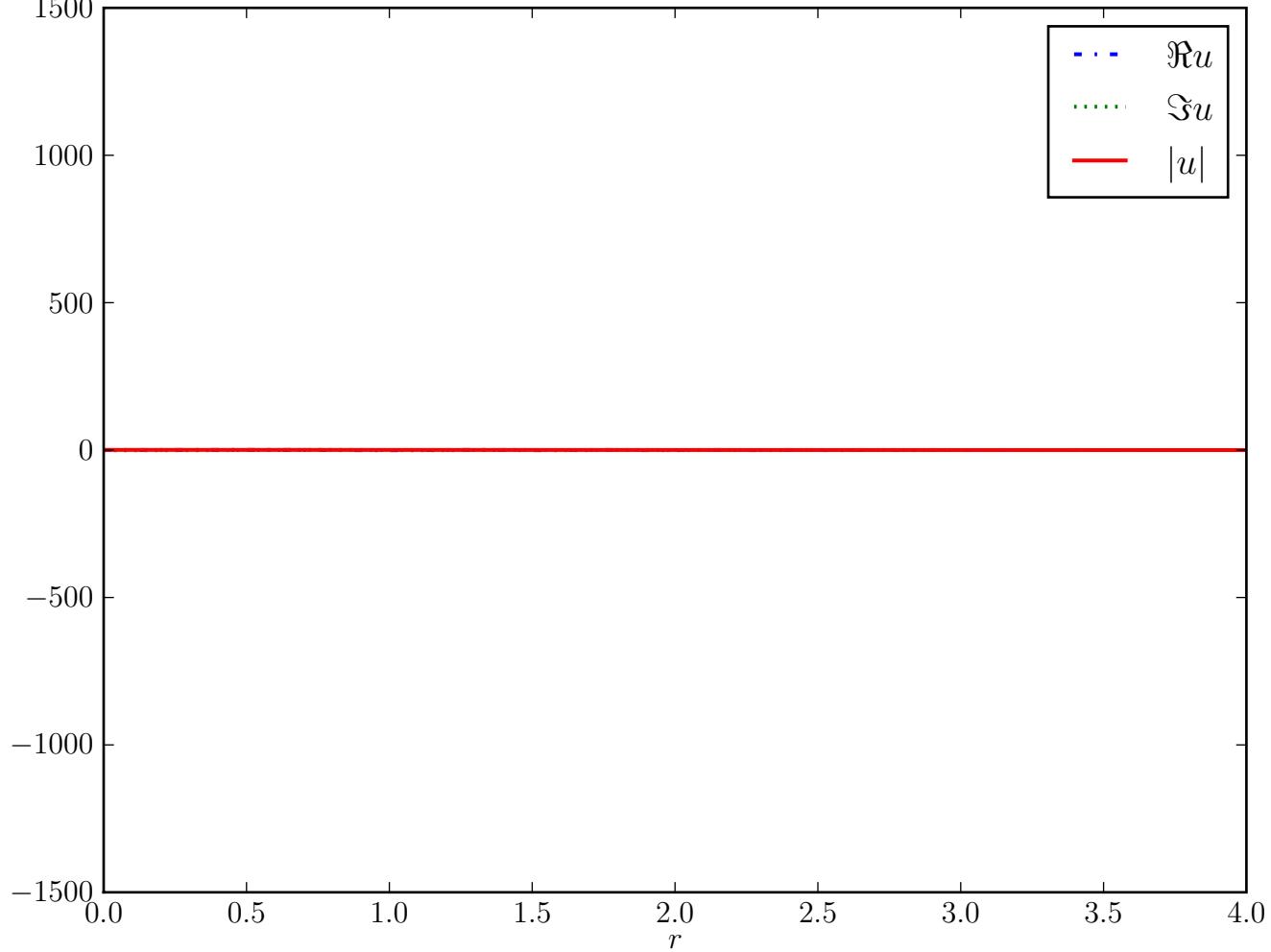
$t = 0.07$



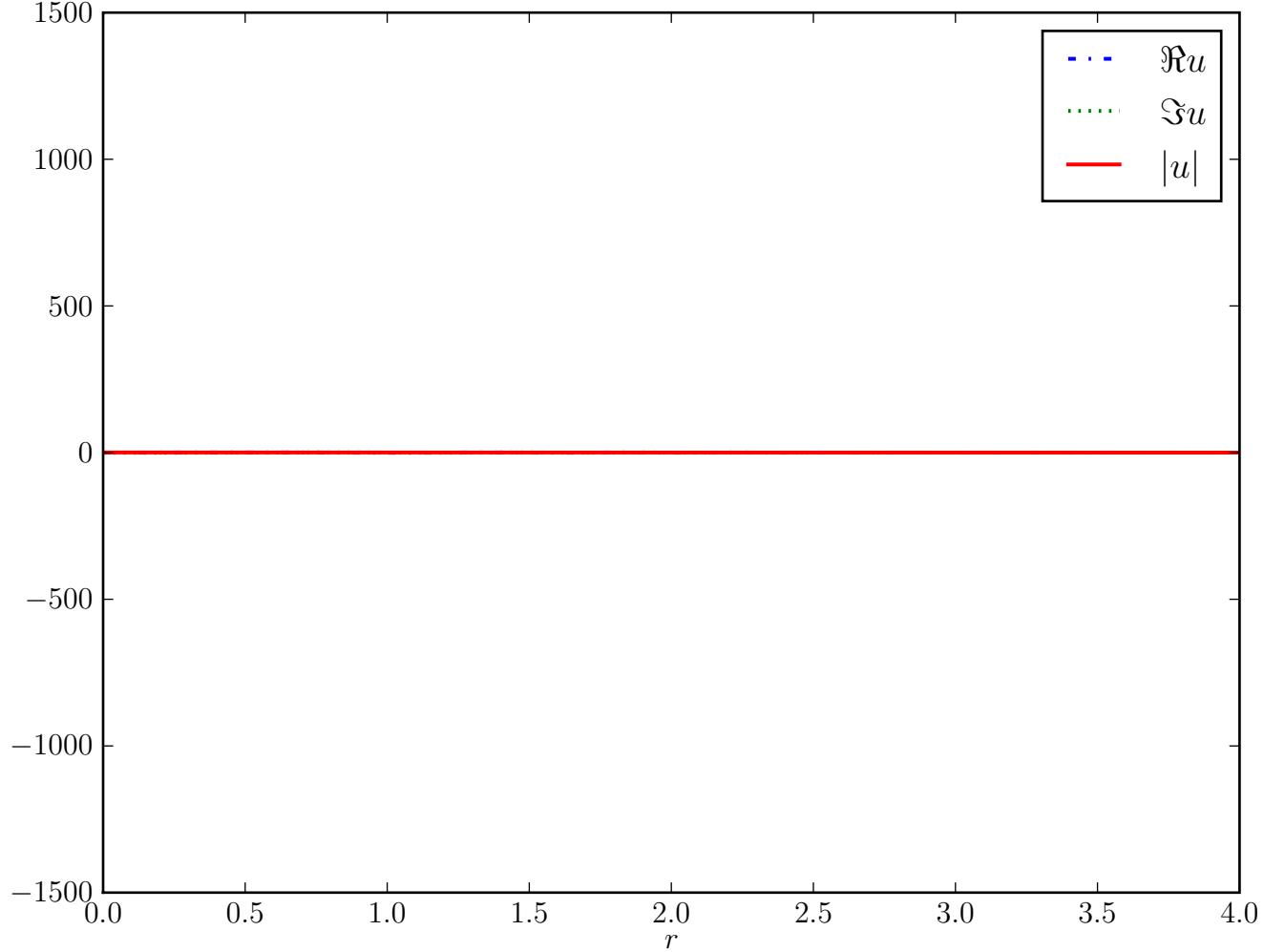
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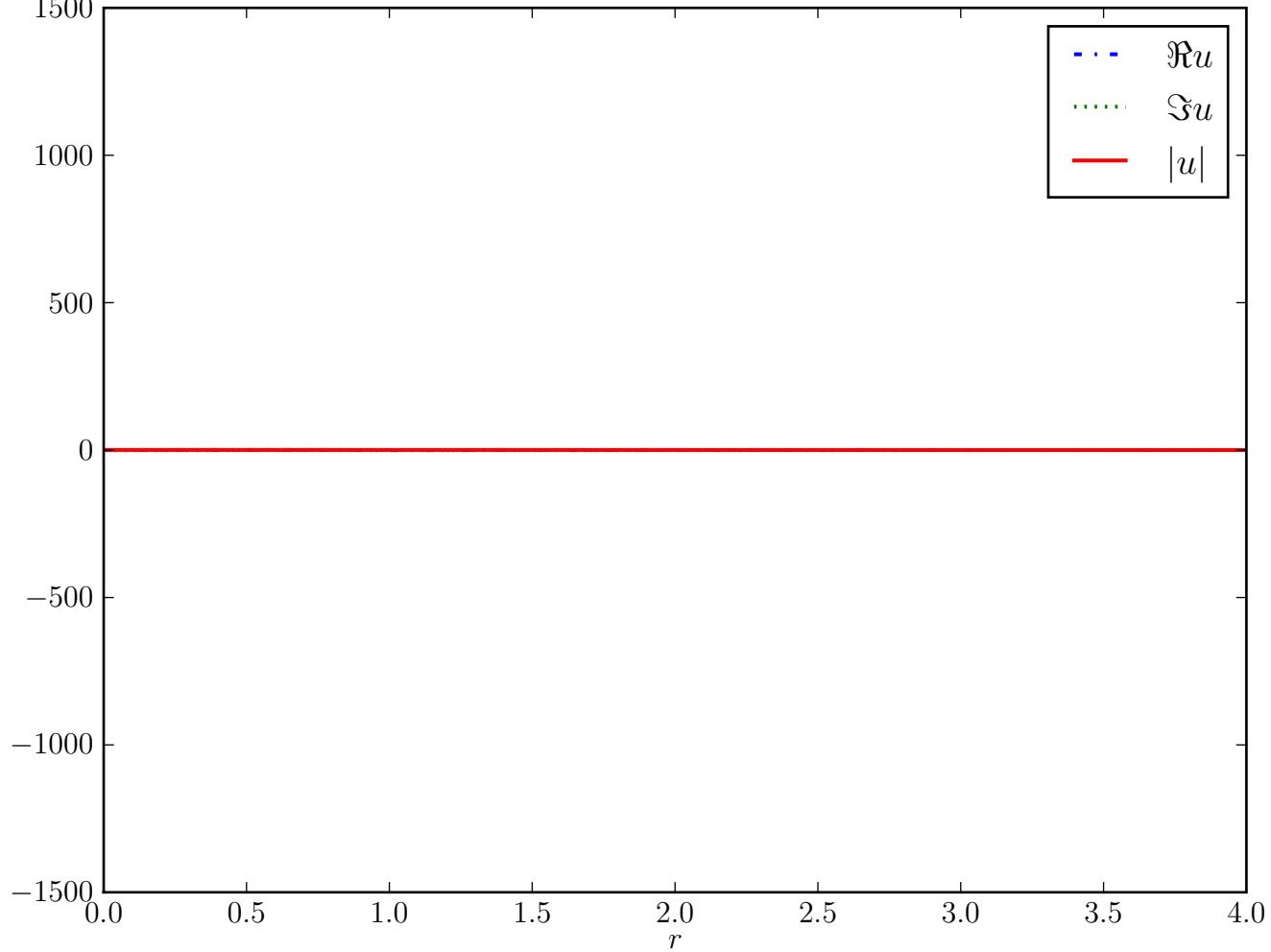
$t = 0.08$



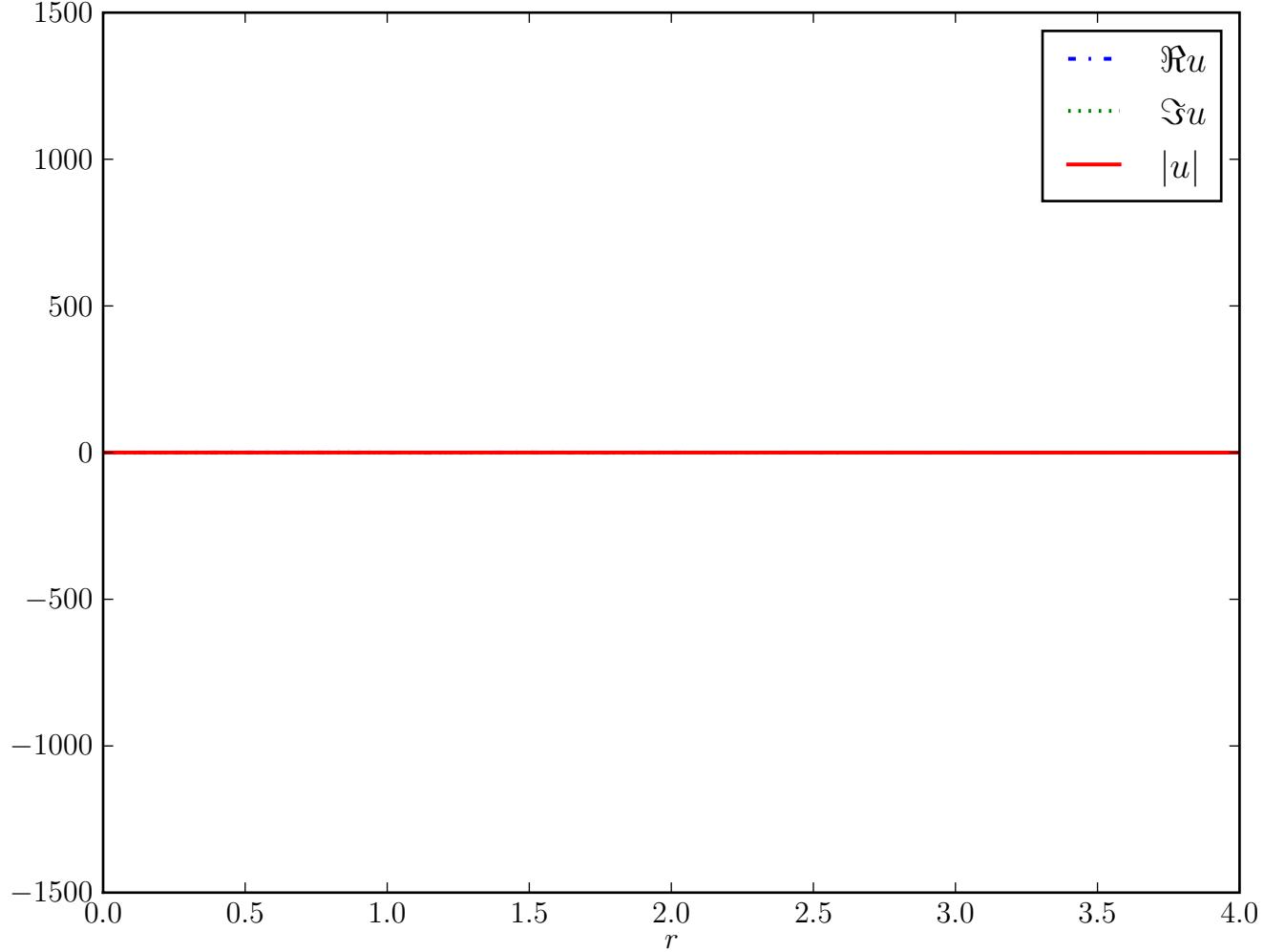
$t = 0.085$



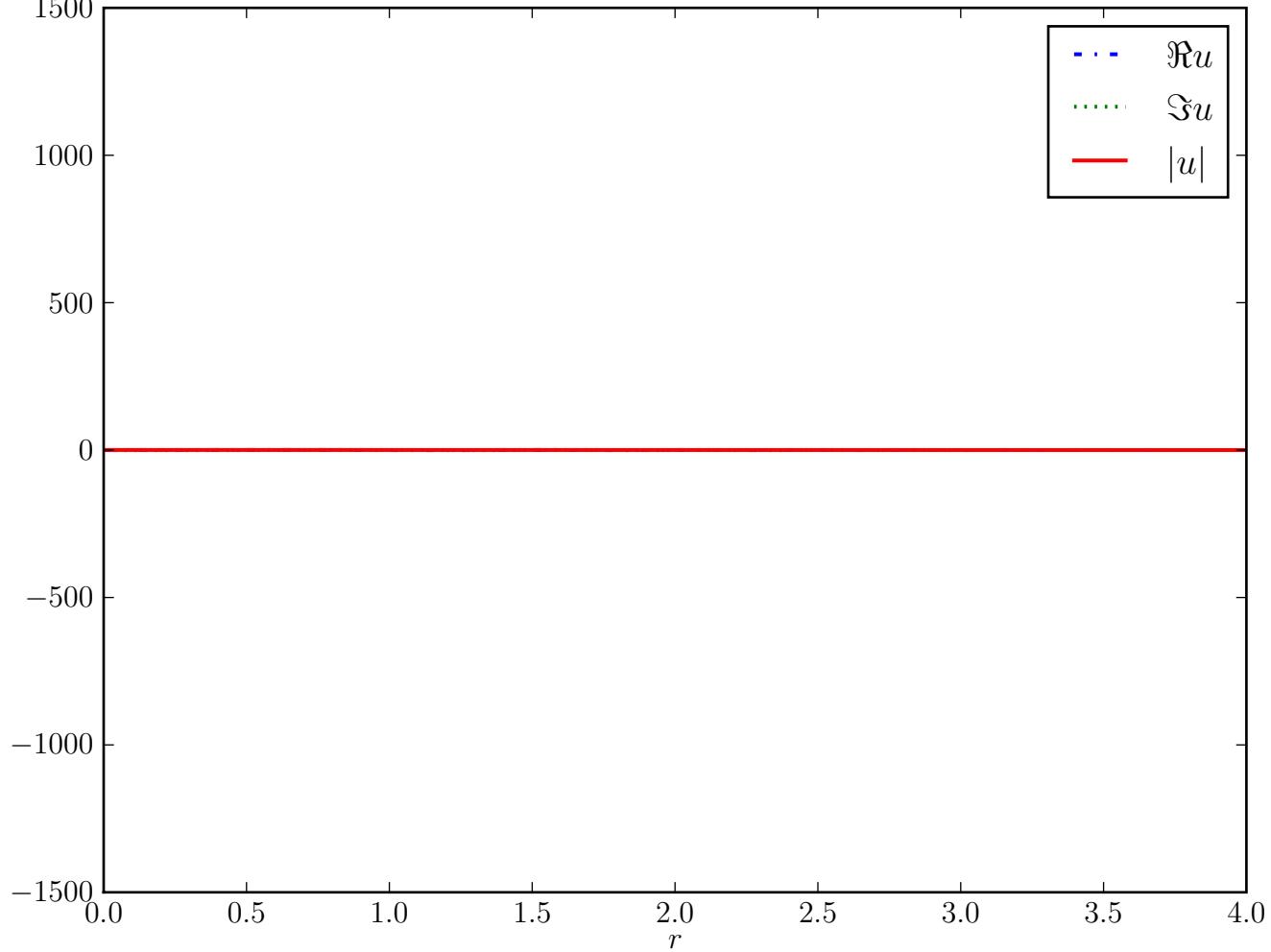
$t = 0.09$



$t = 0.095$

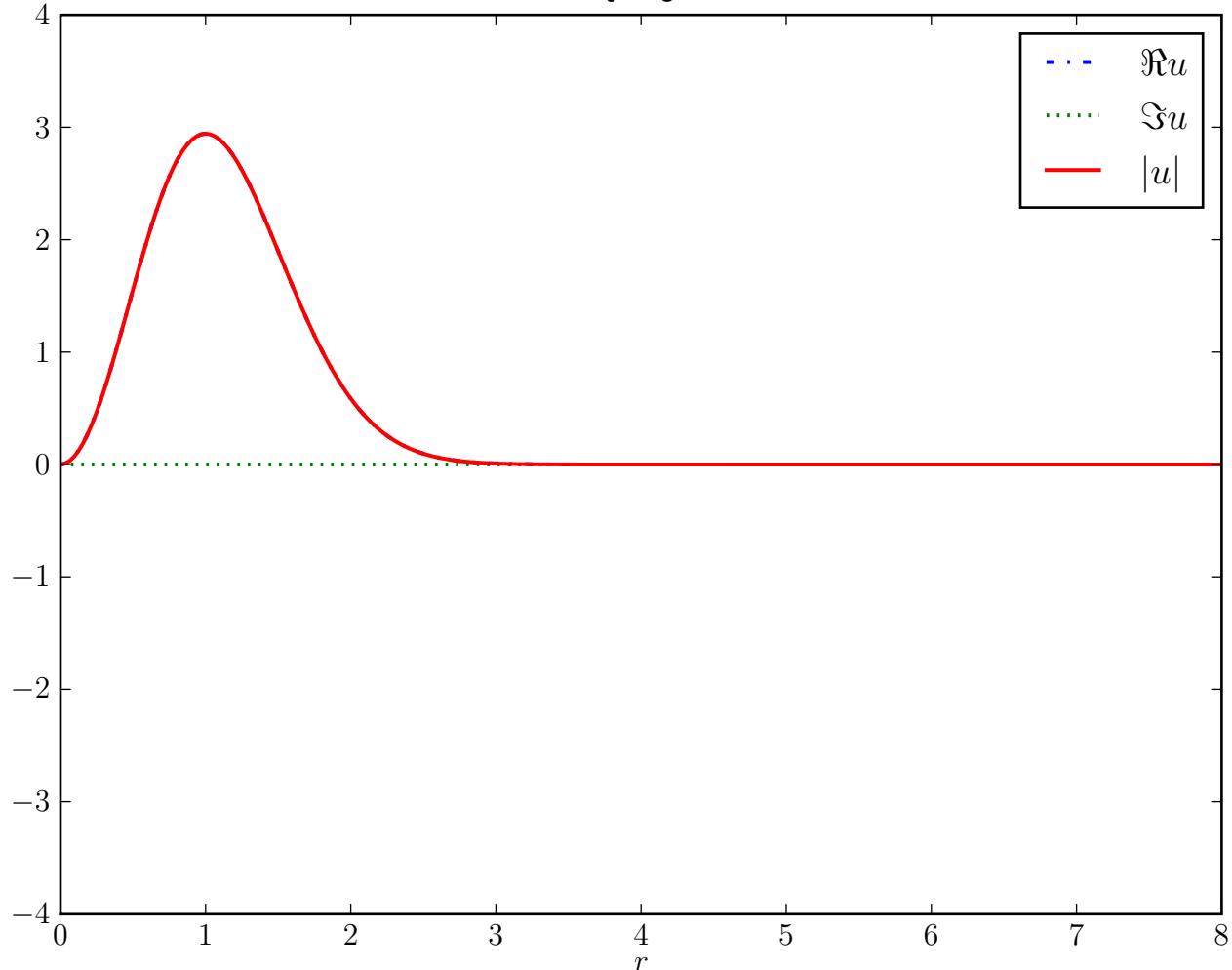


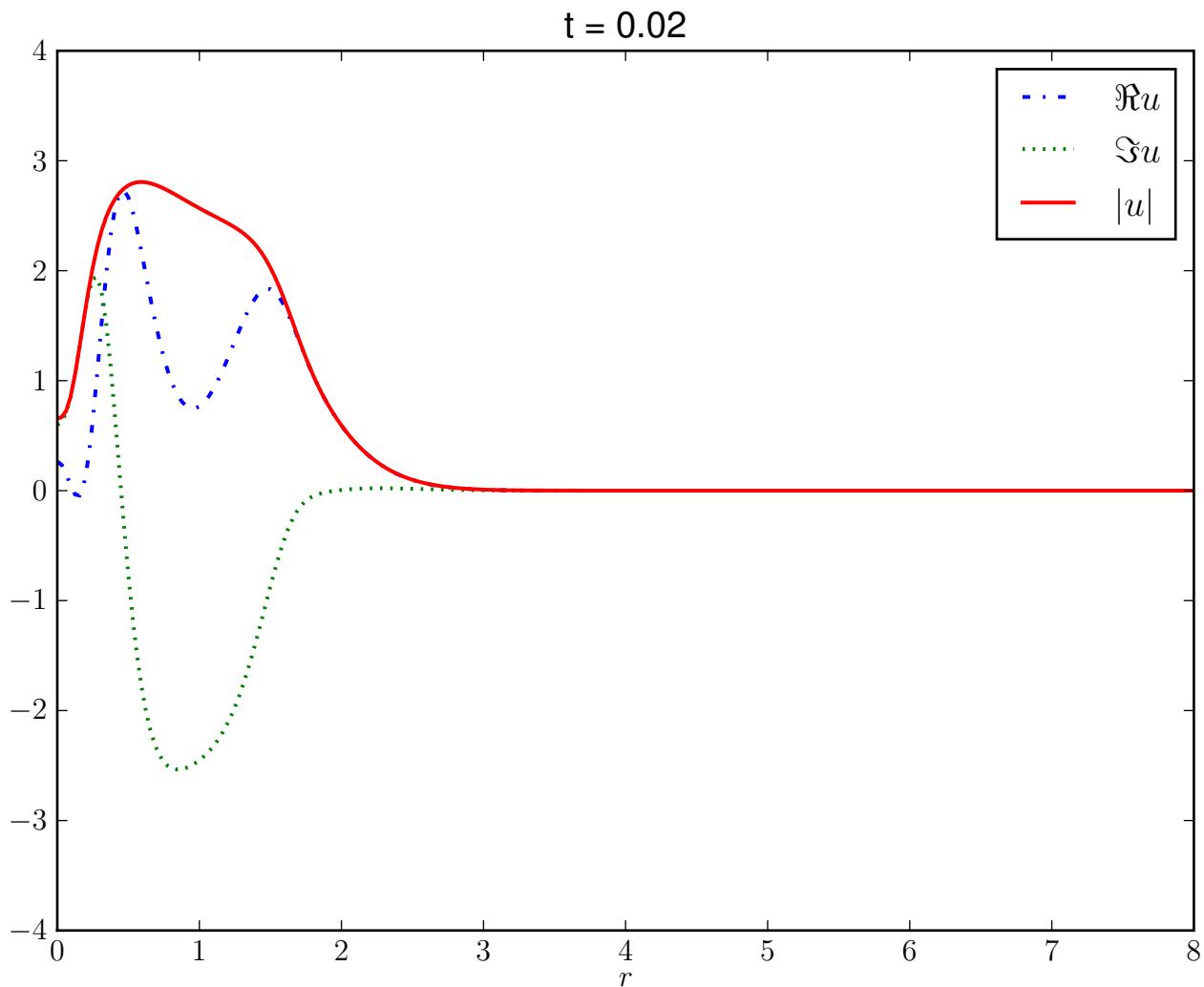
$t = 0.1$



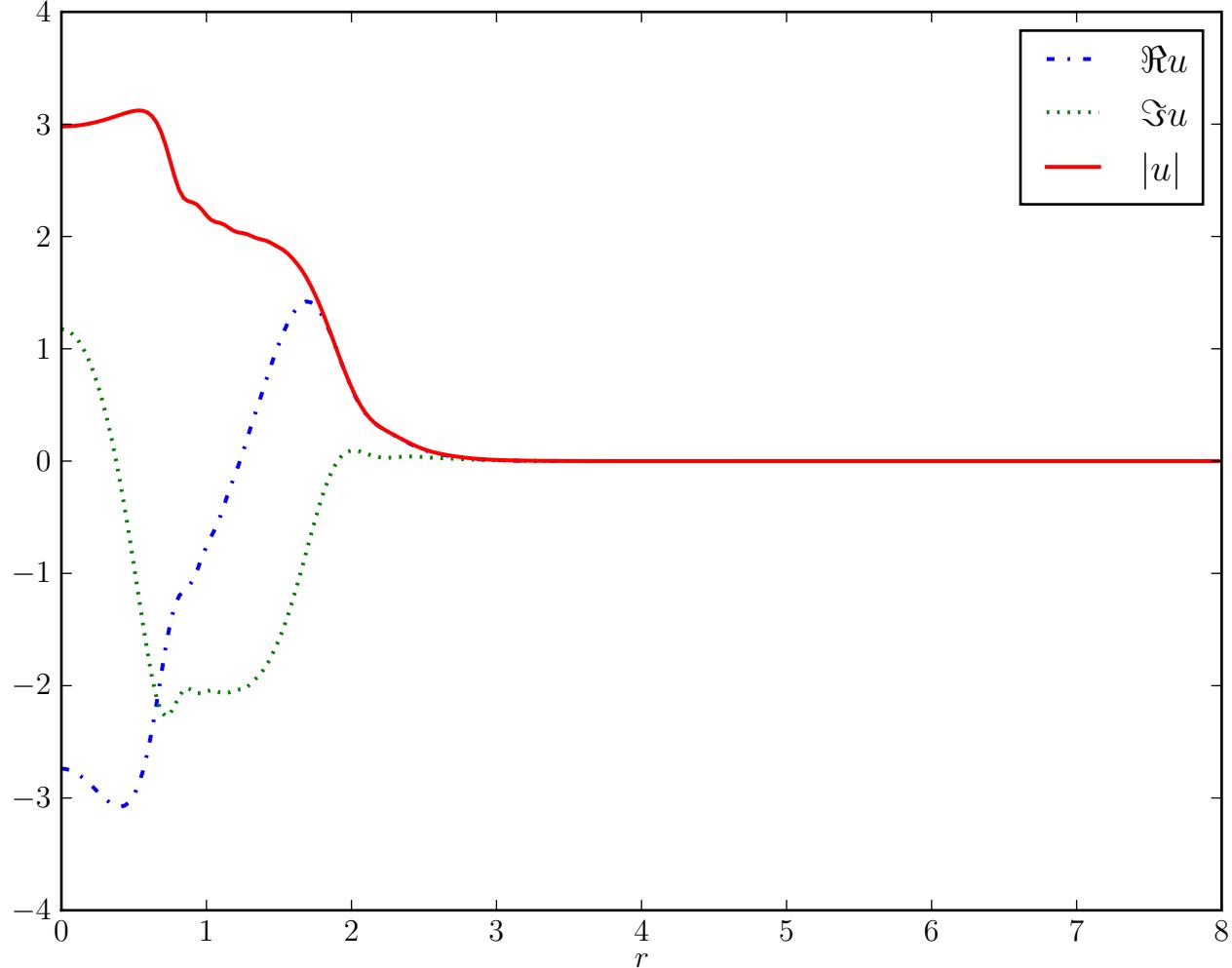
Spherical Ring Initial Data

$t = 0$

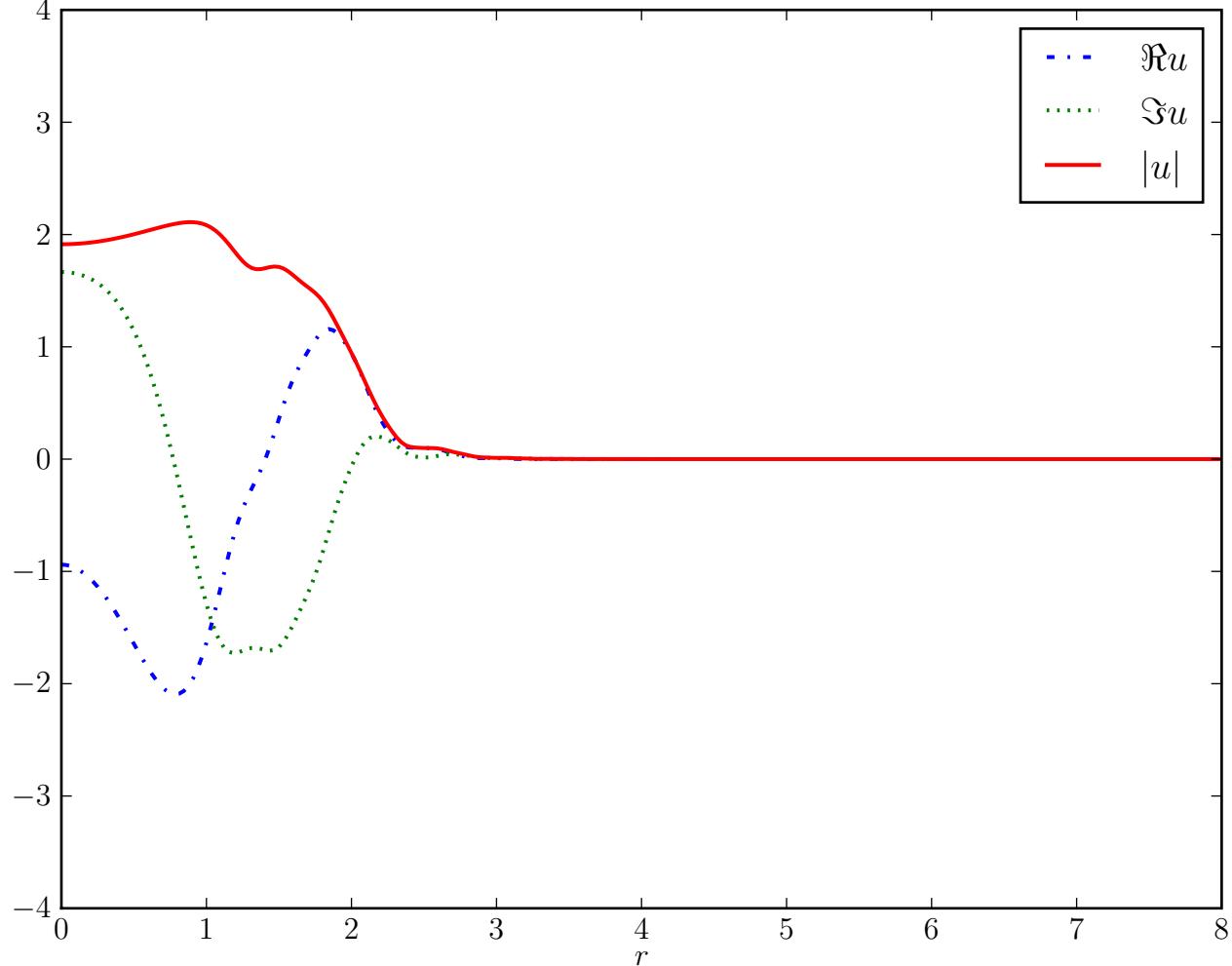




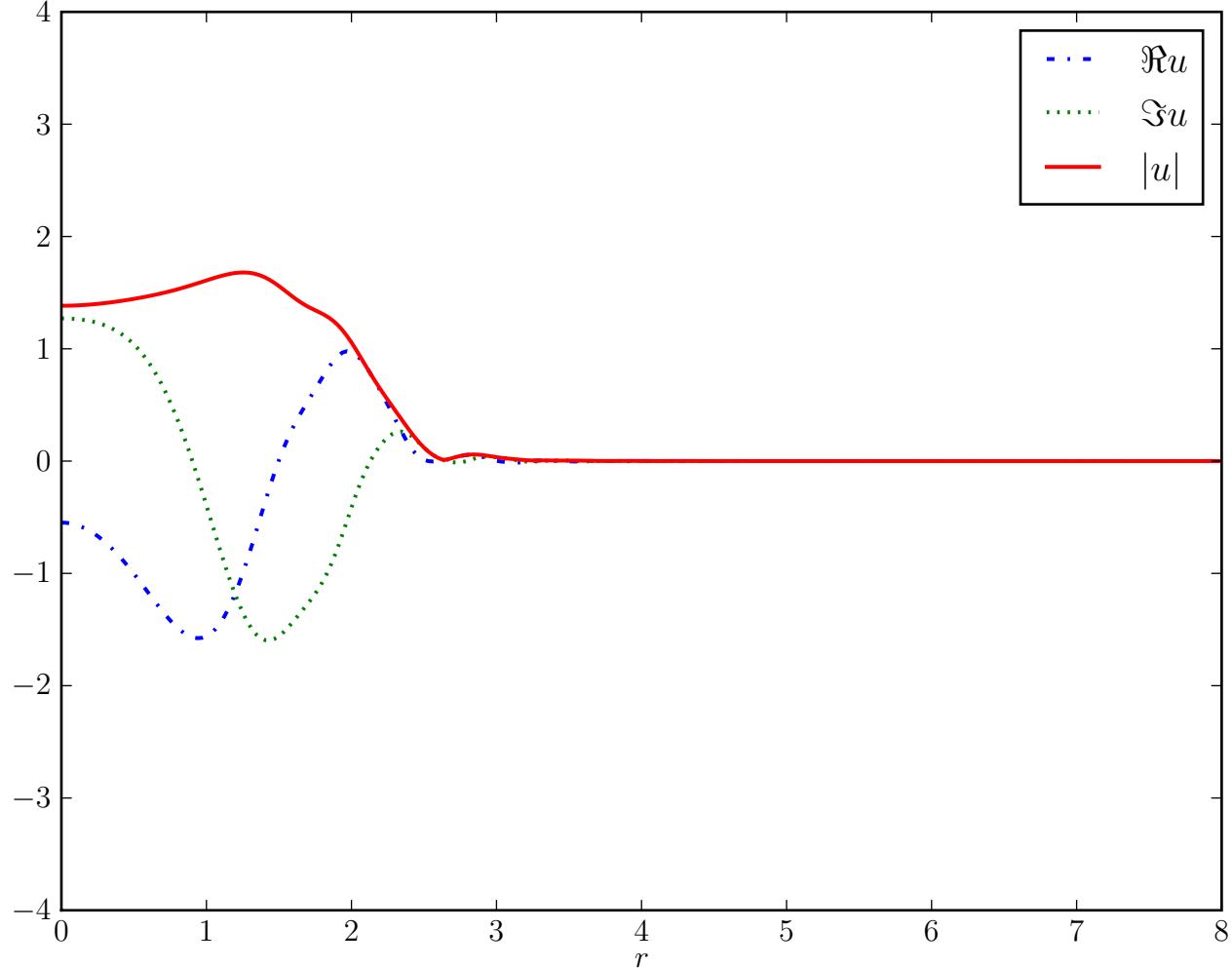
$t = 0.04$

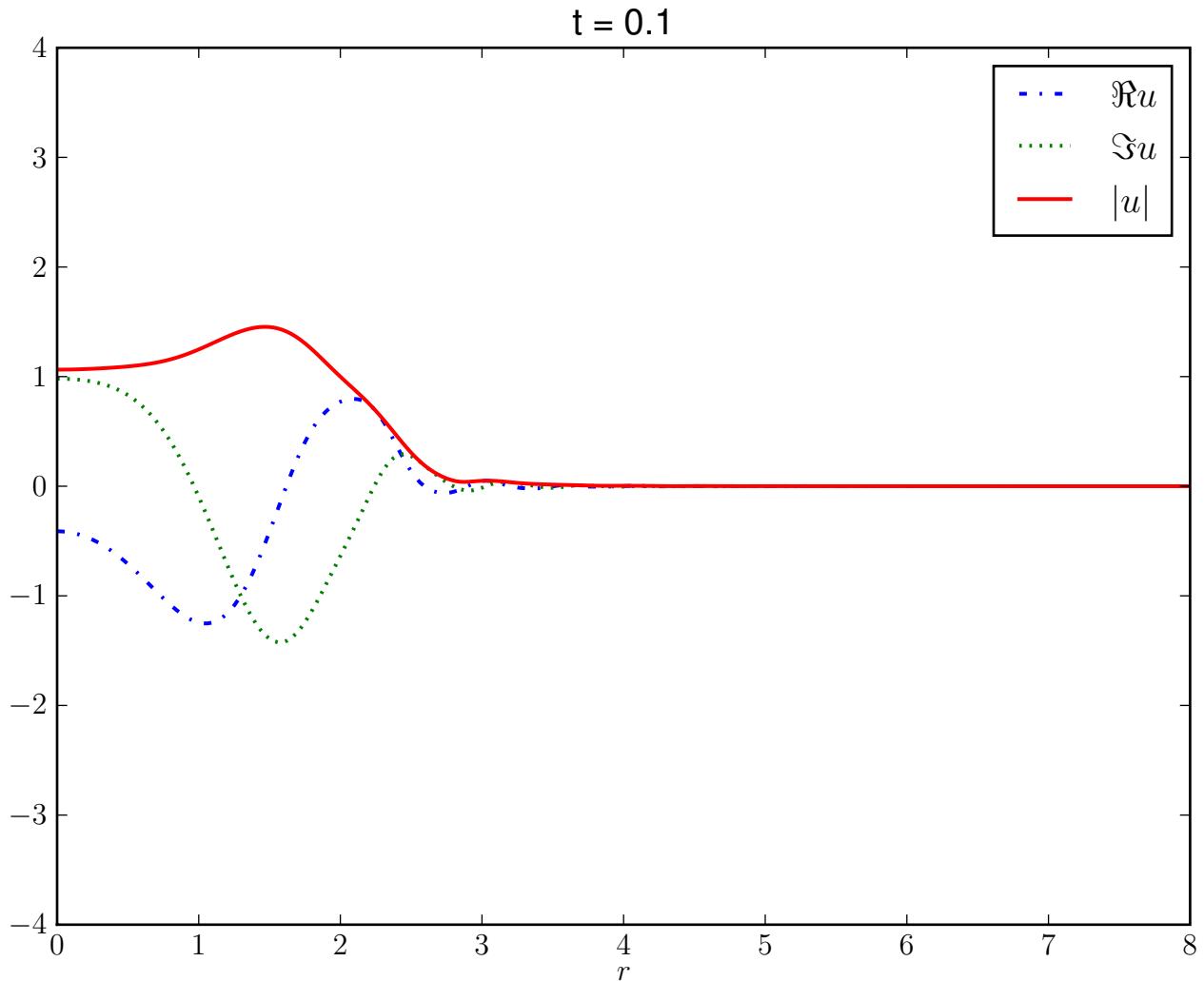


$t = 0.06$

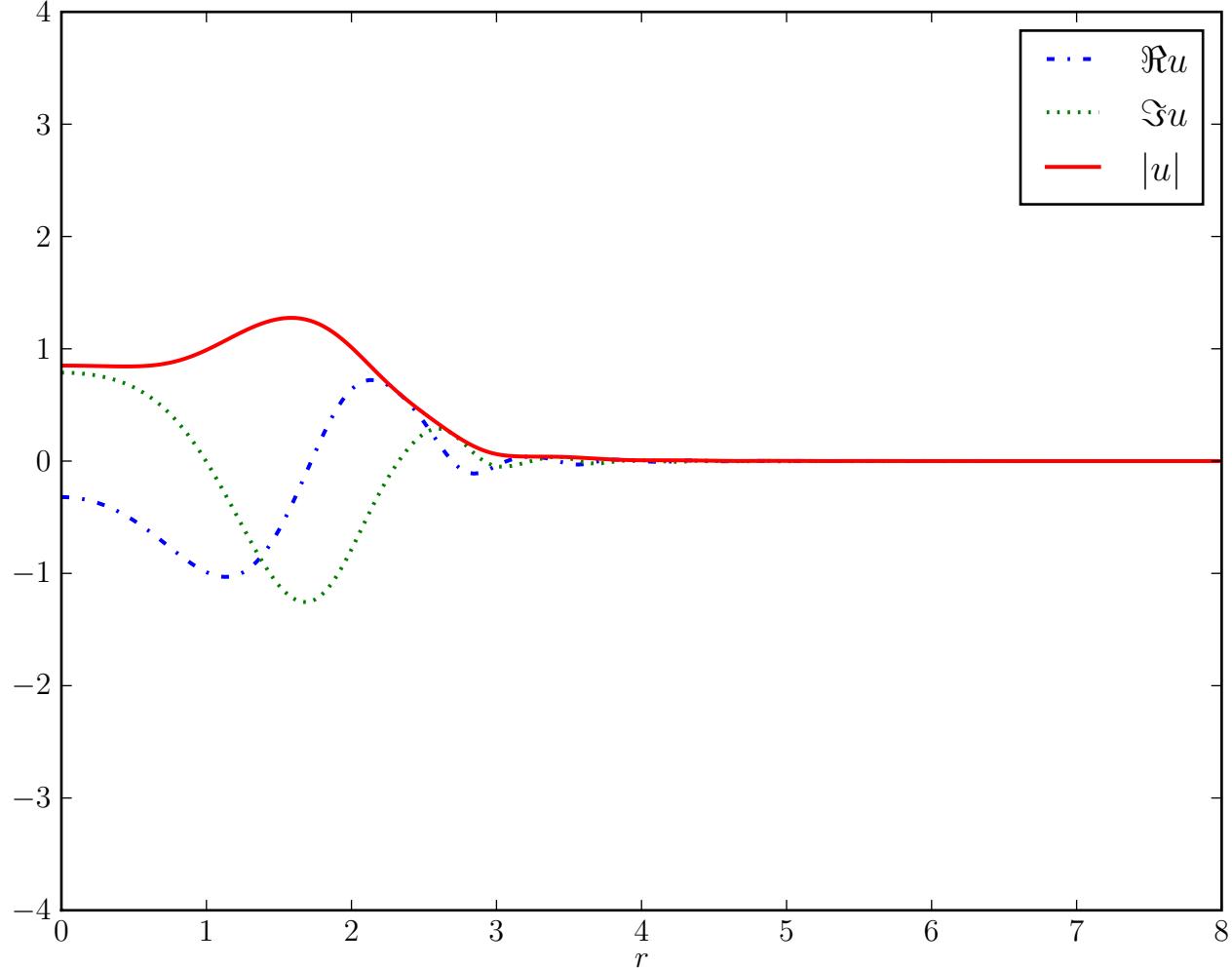


$t = 0.08$

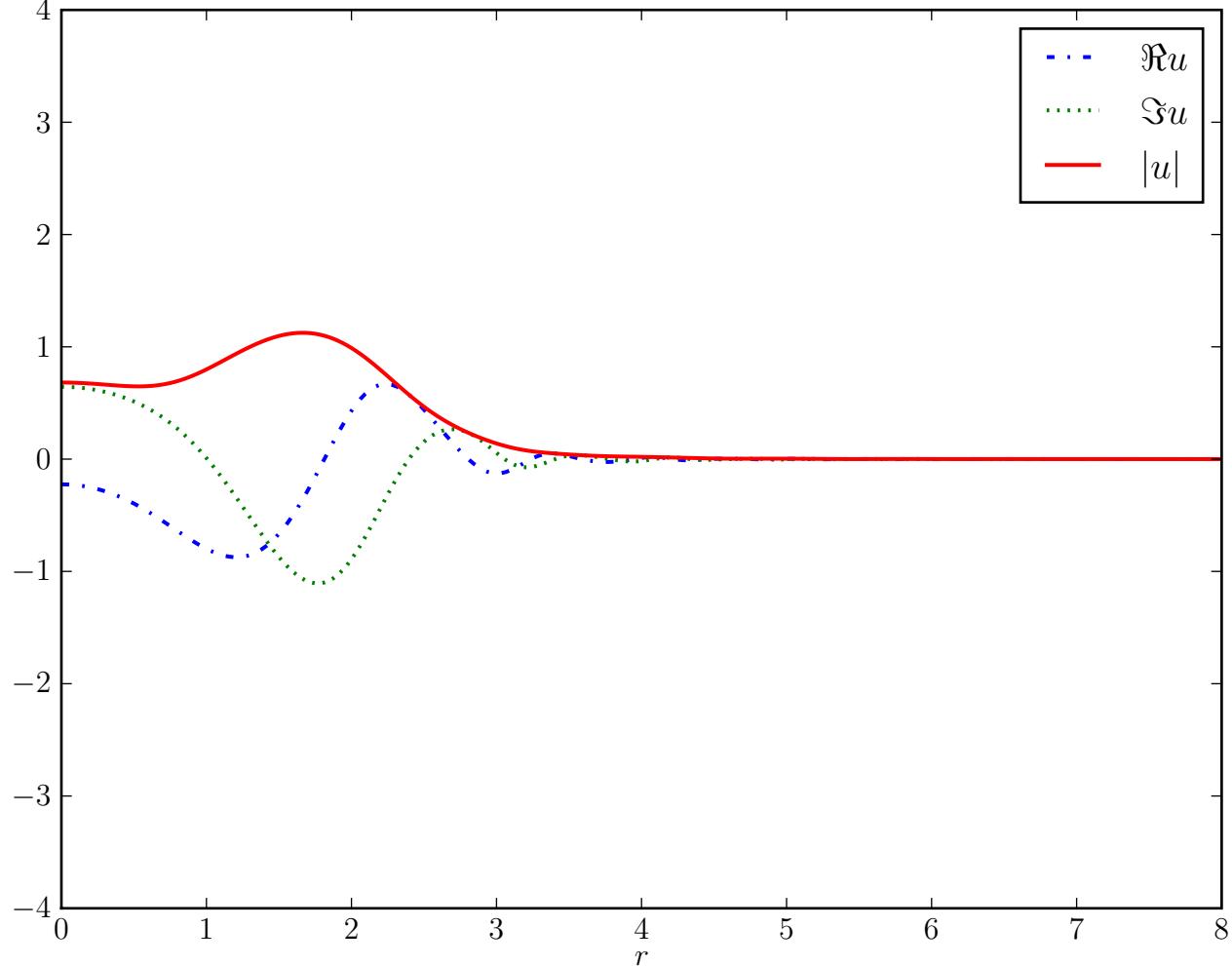




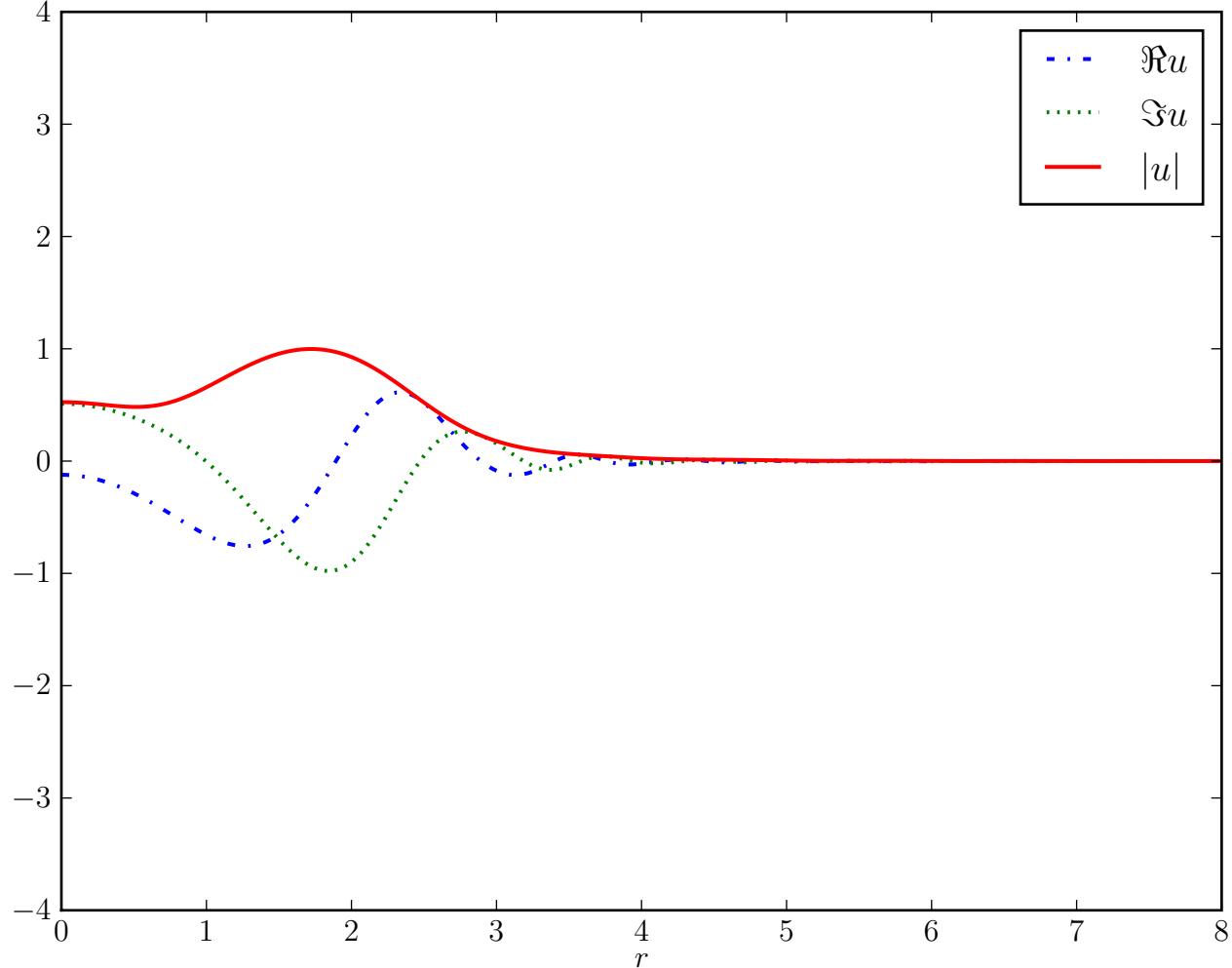
$t = 0.12$



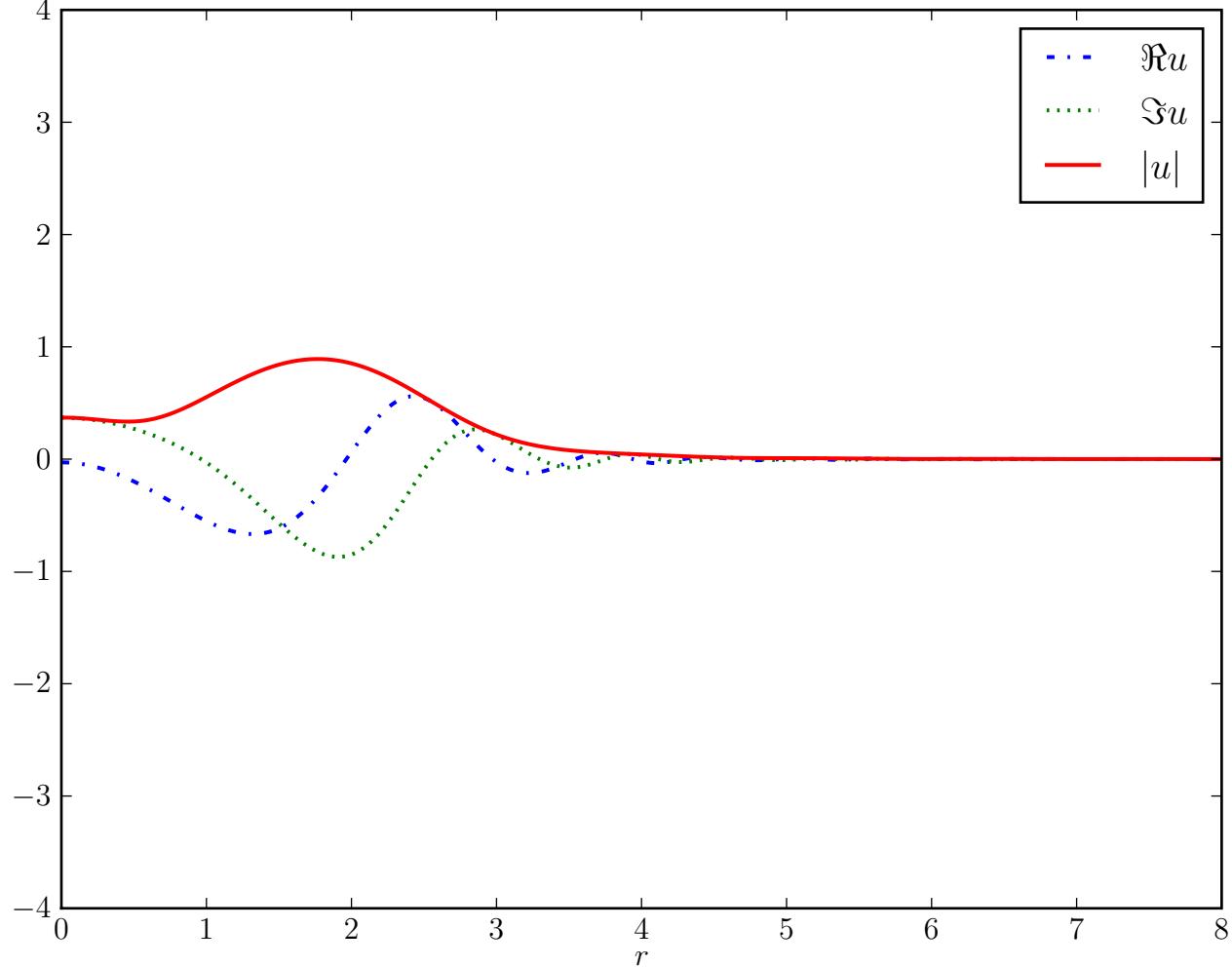
$t = 0.14$

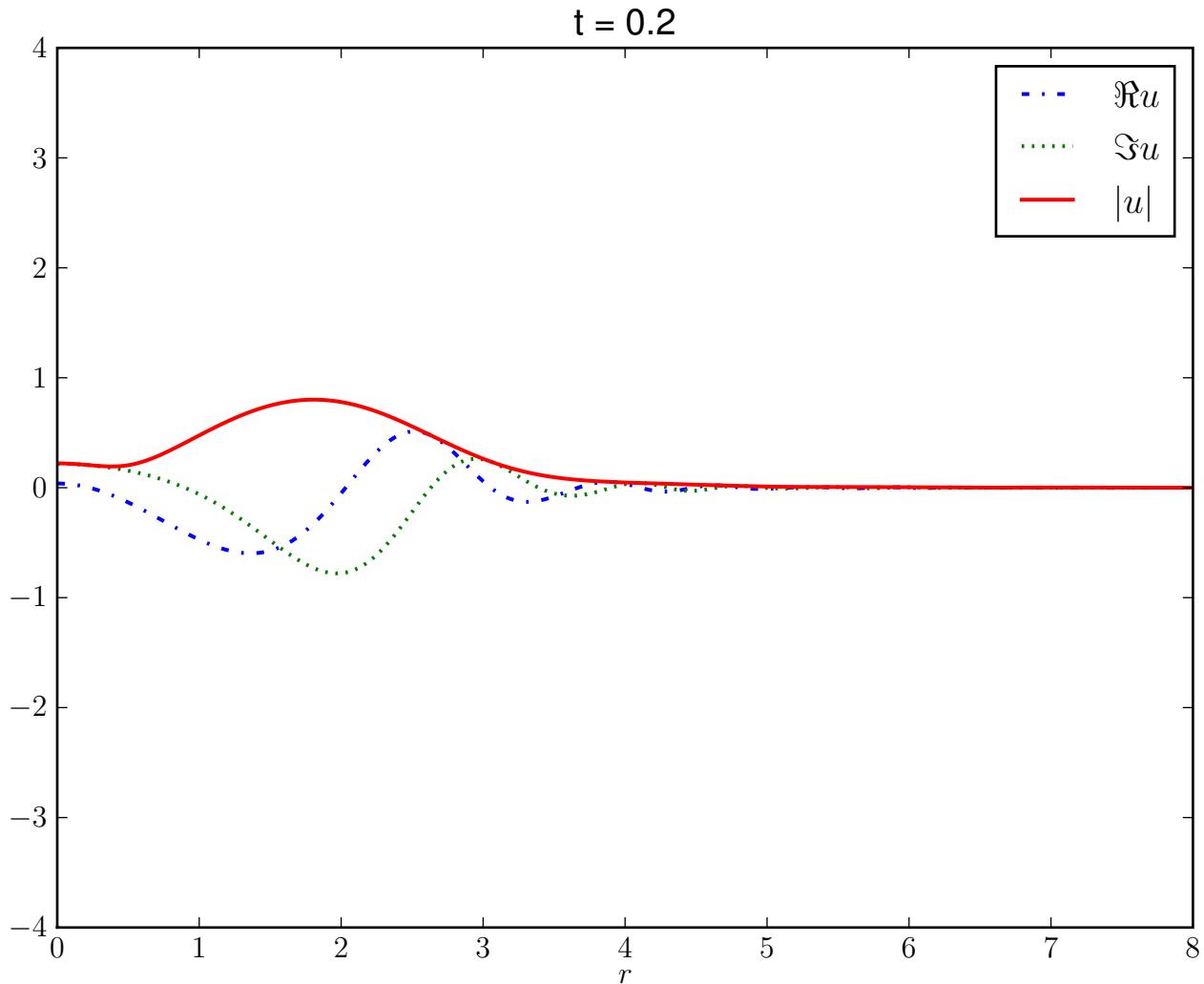


$t = 0.16$

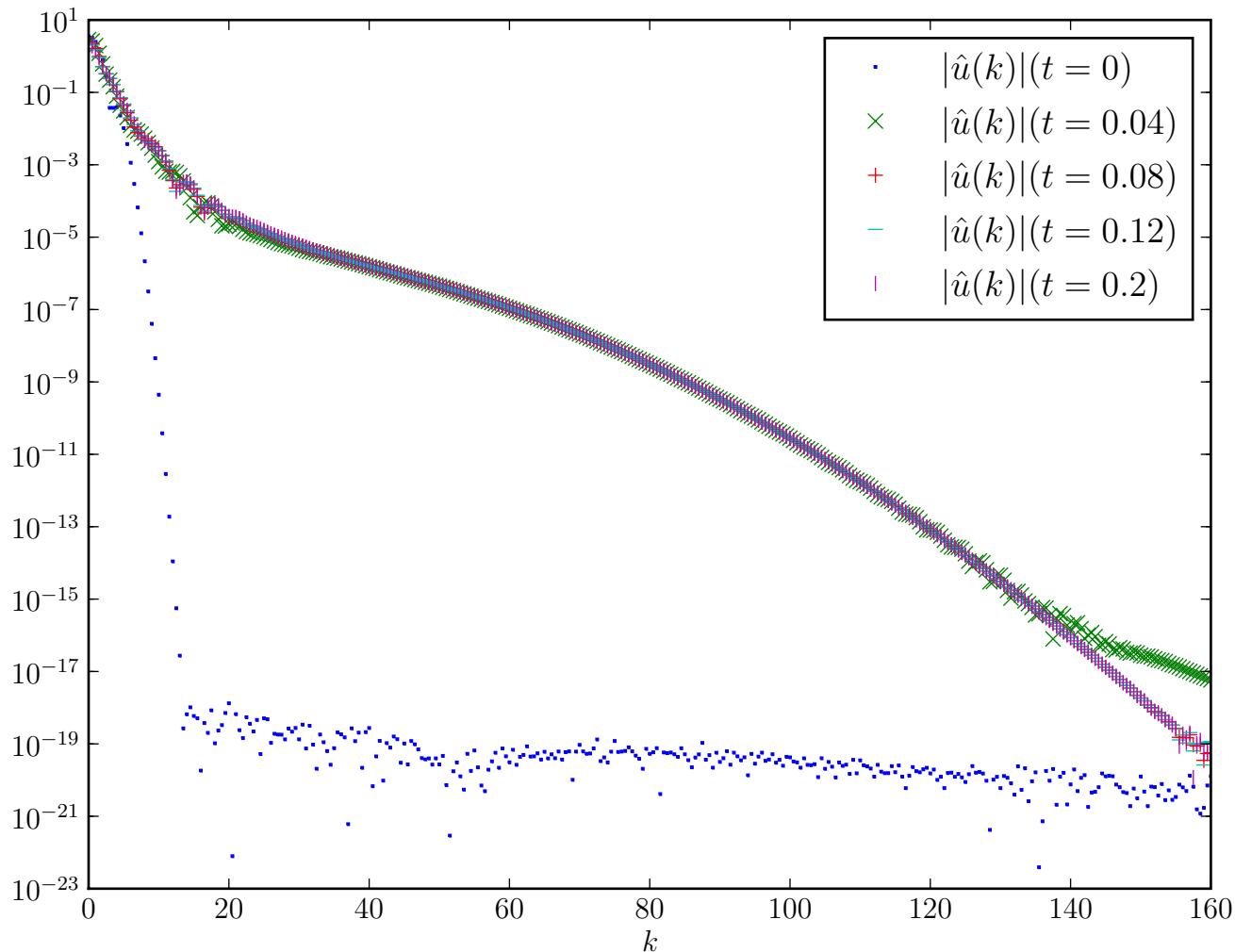


$t = 0.18$

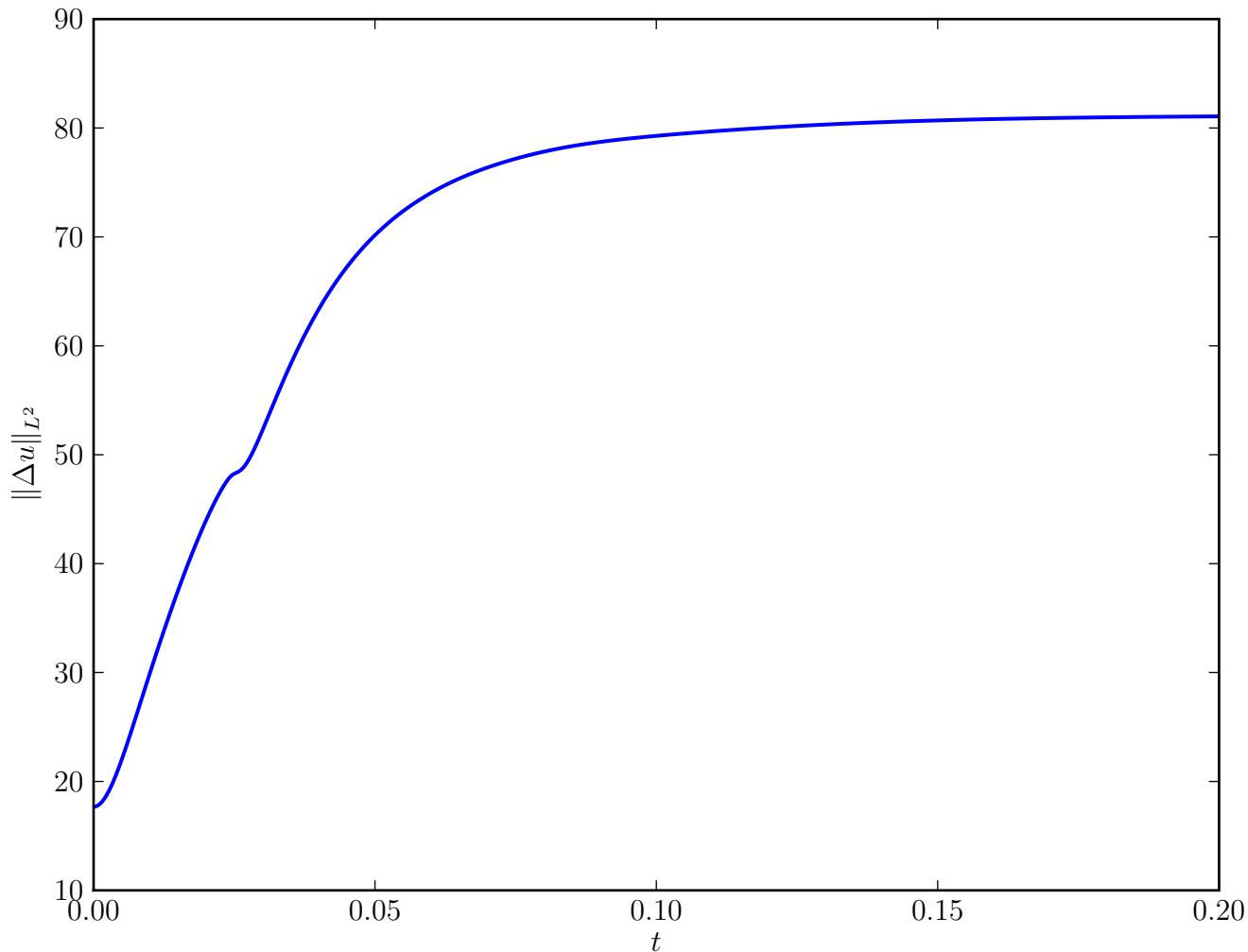




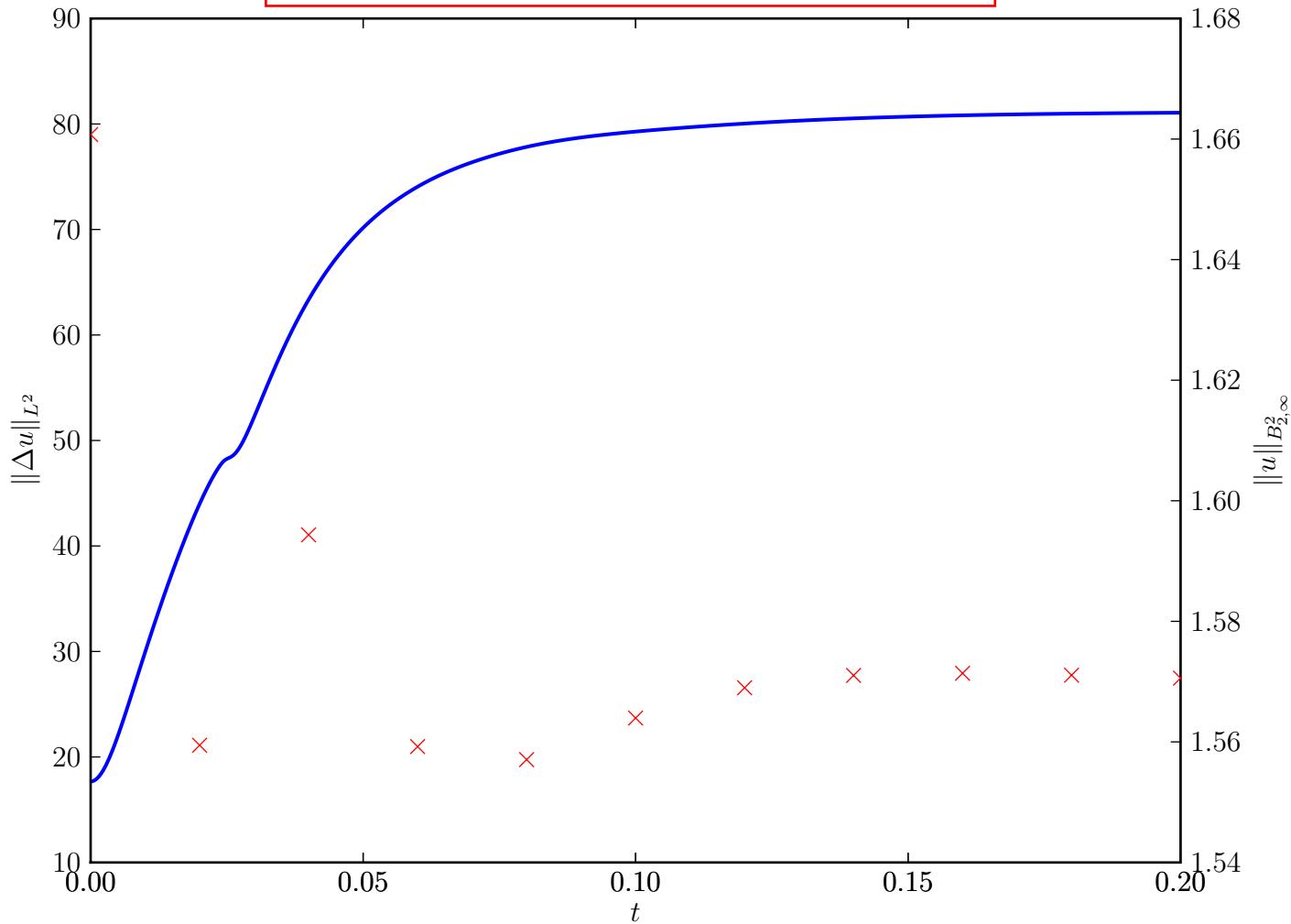
Spherical Ring Fourier transform snapshots along nonlinear flow



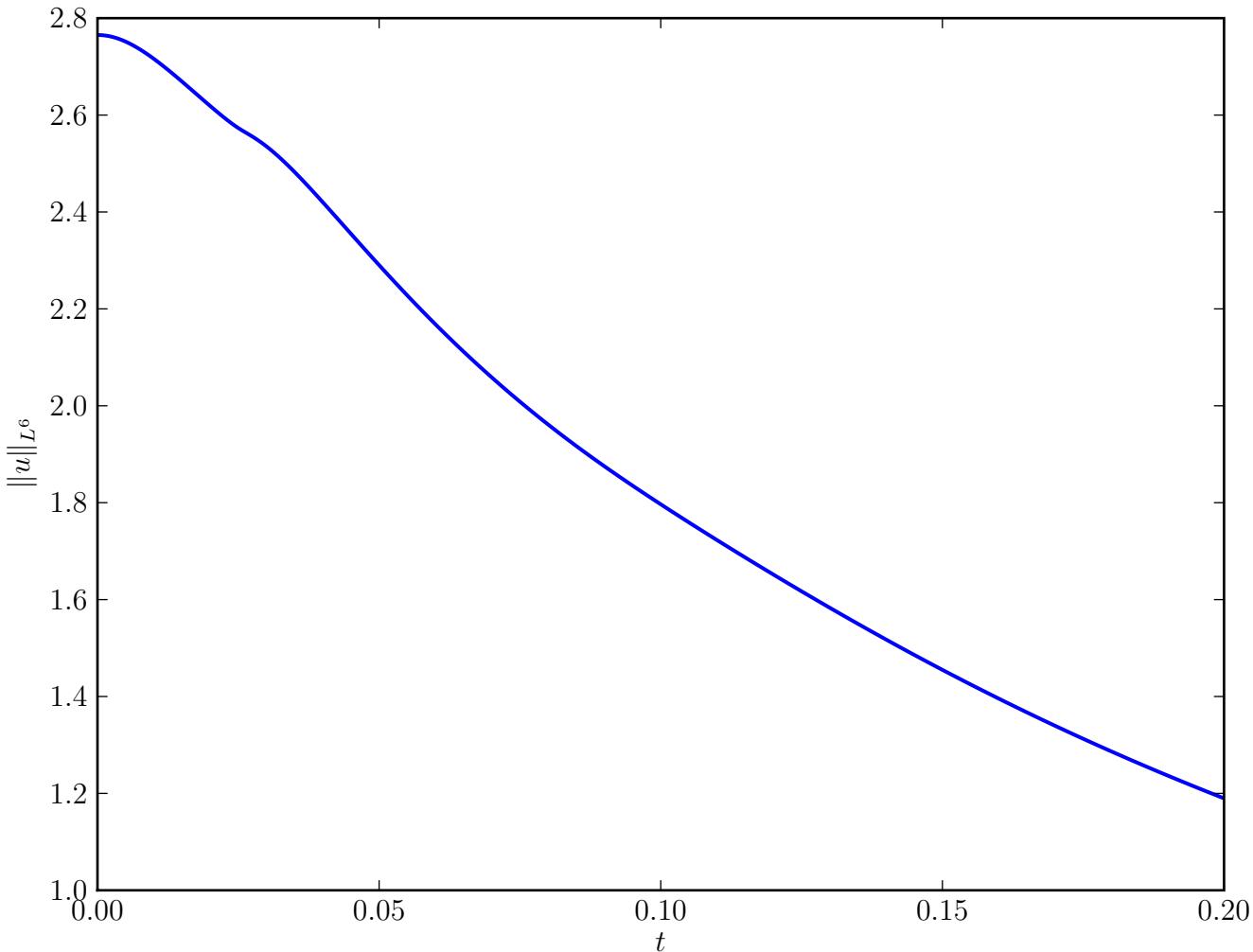
H² norm of Spherical Ring along nonlinear flow



Sobolev vs. Besov: Spherical Ring along nonlinear flow

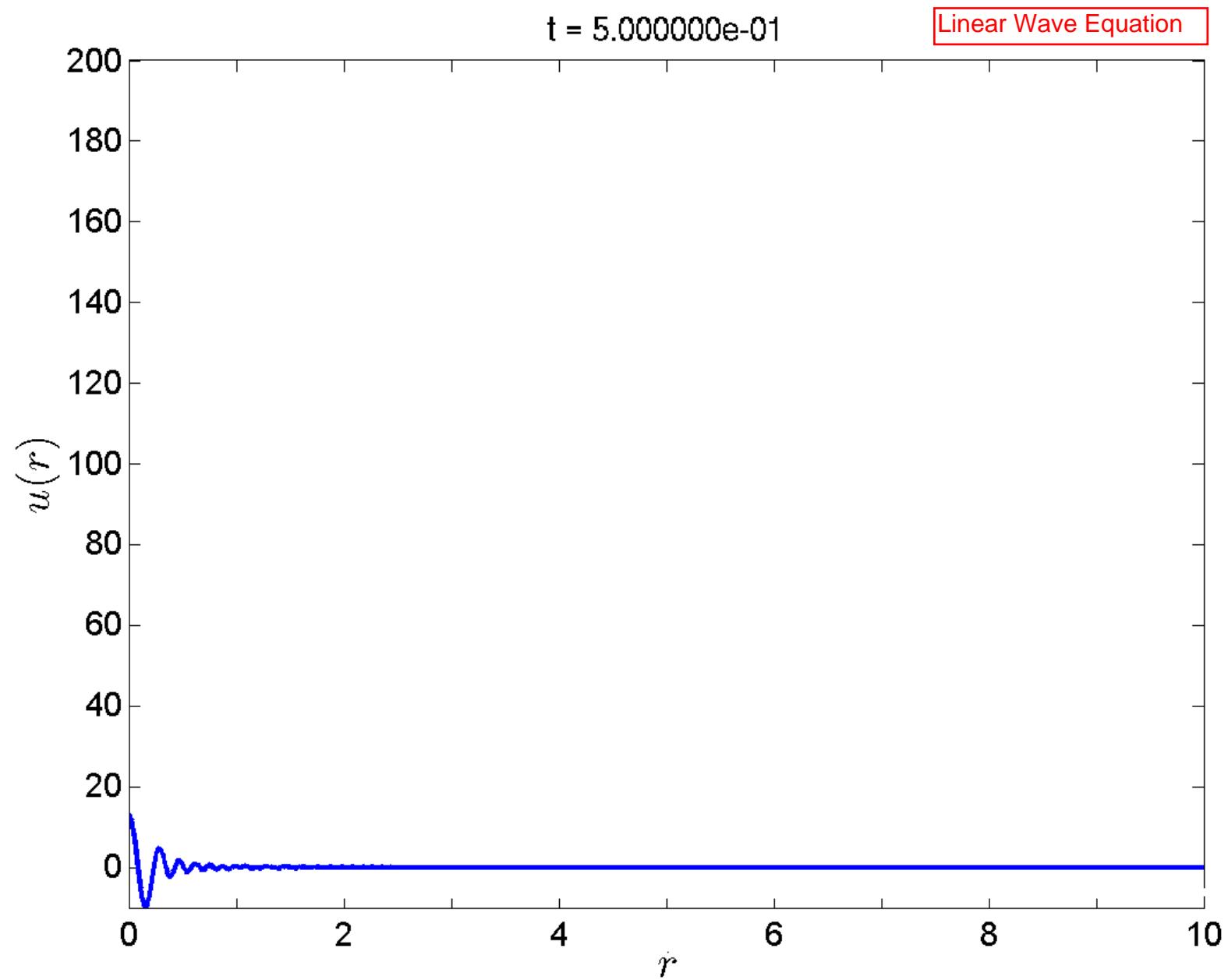


Potential Energy Norm Decay

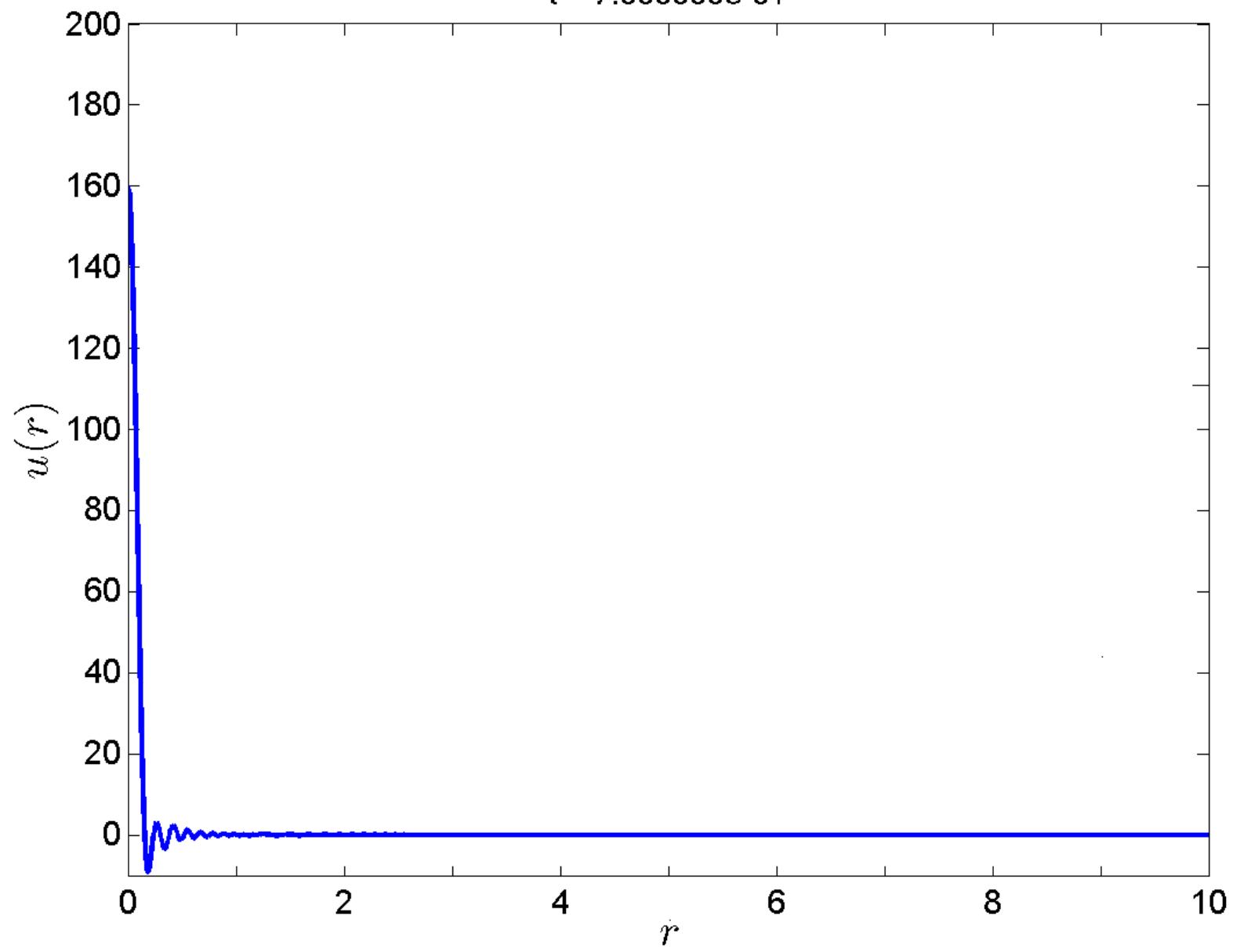


C-Pang-Simpson-Sulem simulations

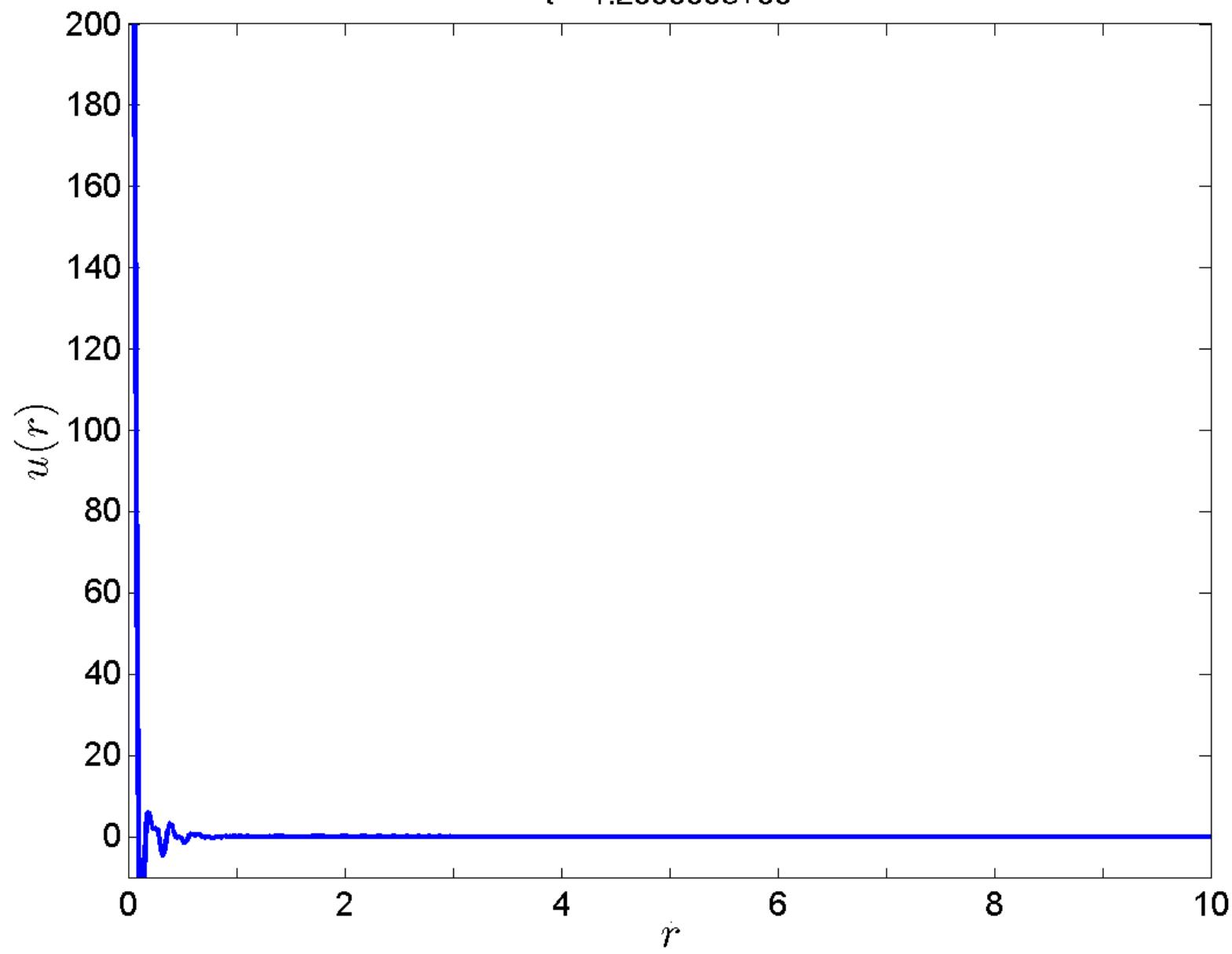
NLW_5 (R^5)



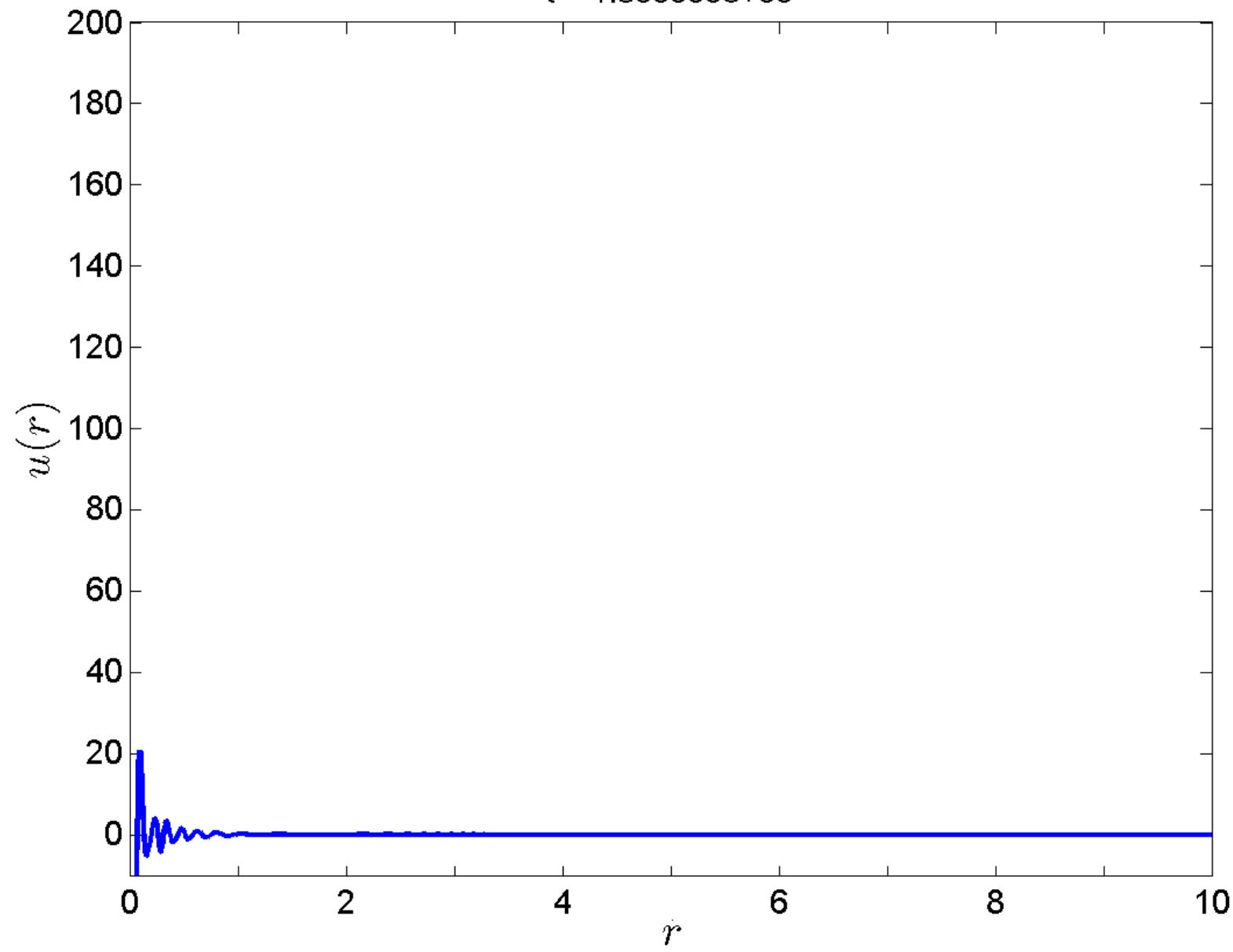
$t = 7.000000e-01$

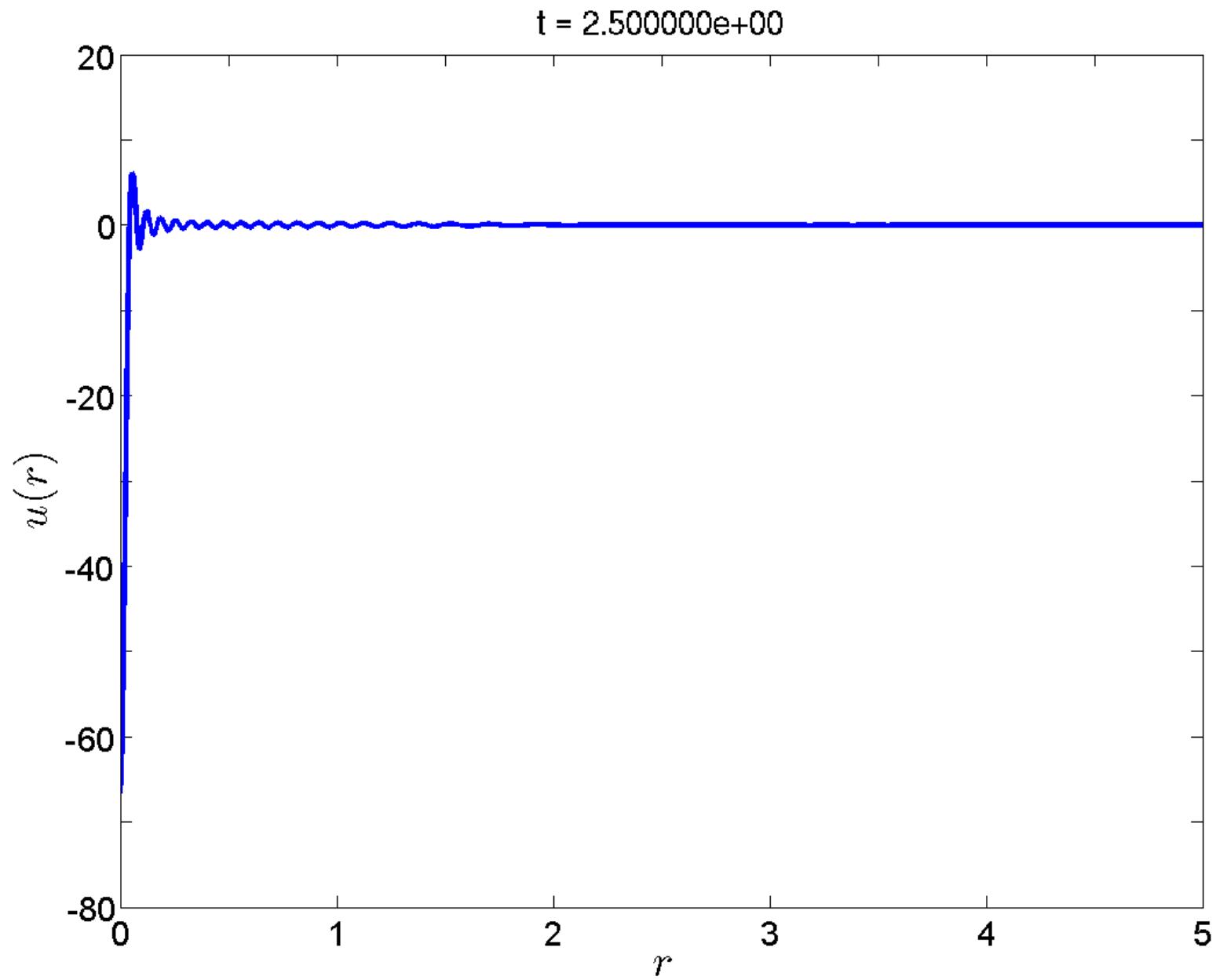


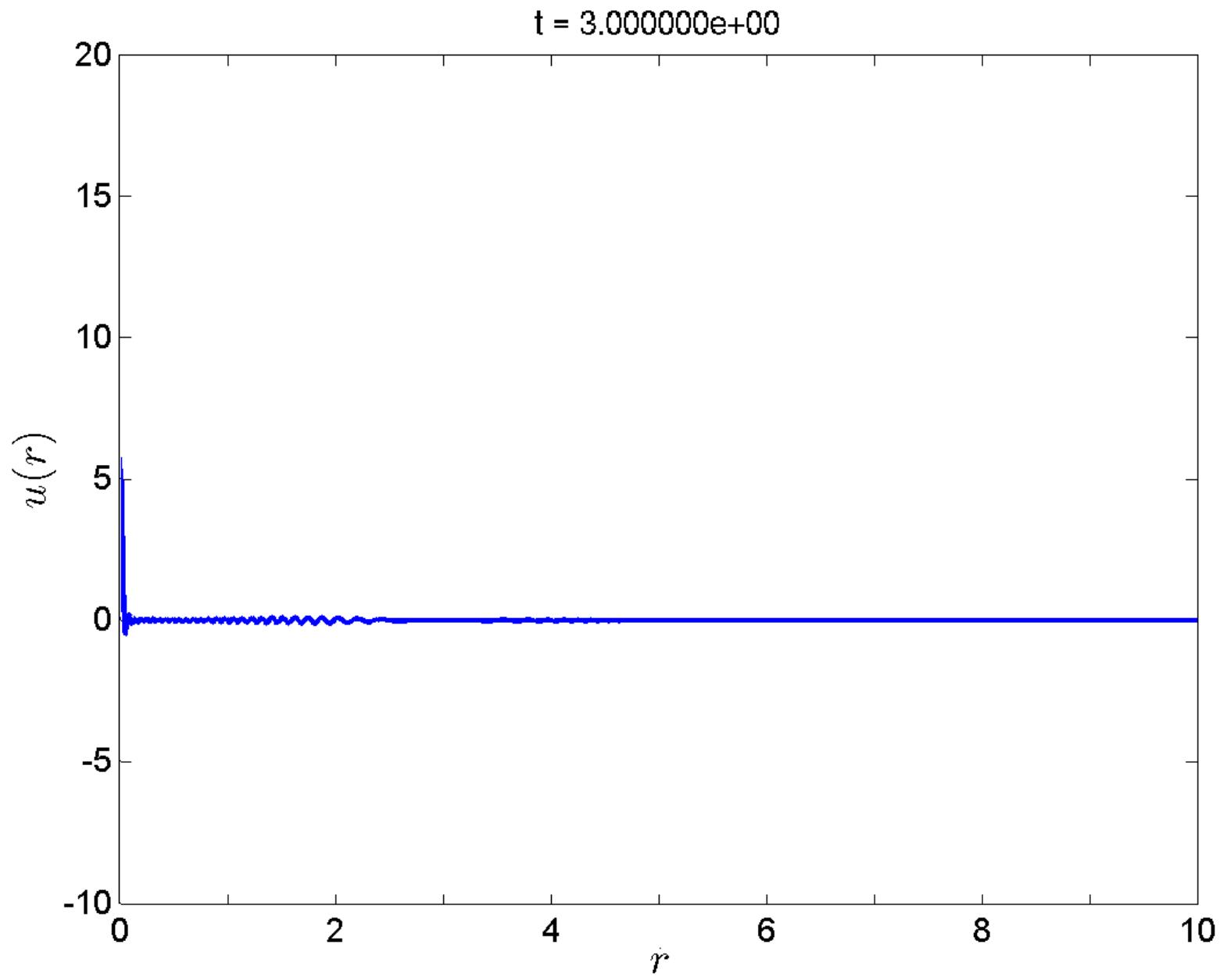
$t = 1.200000e+00$

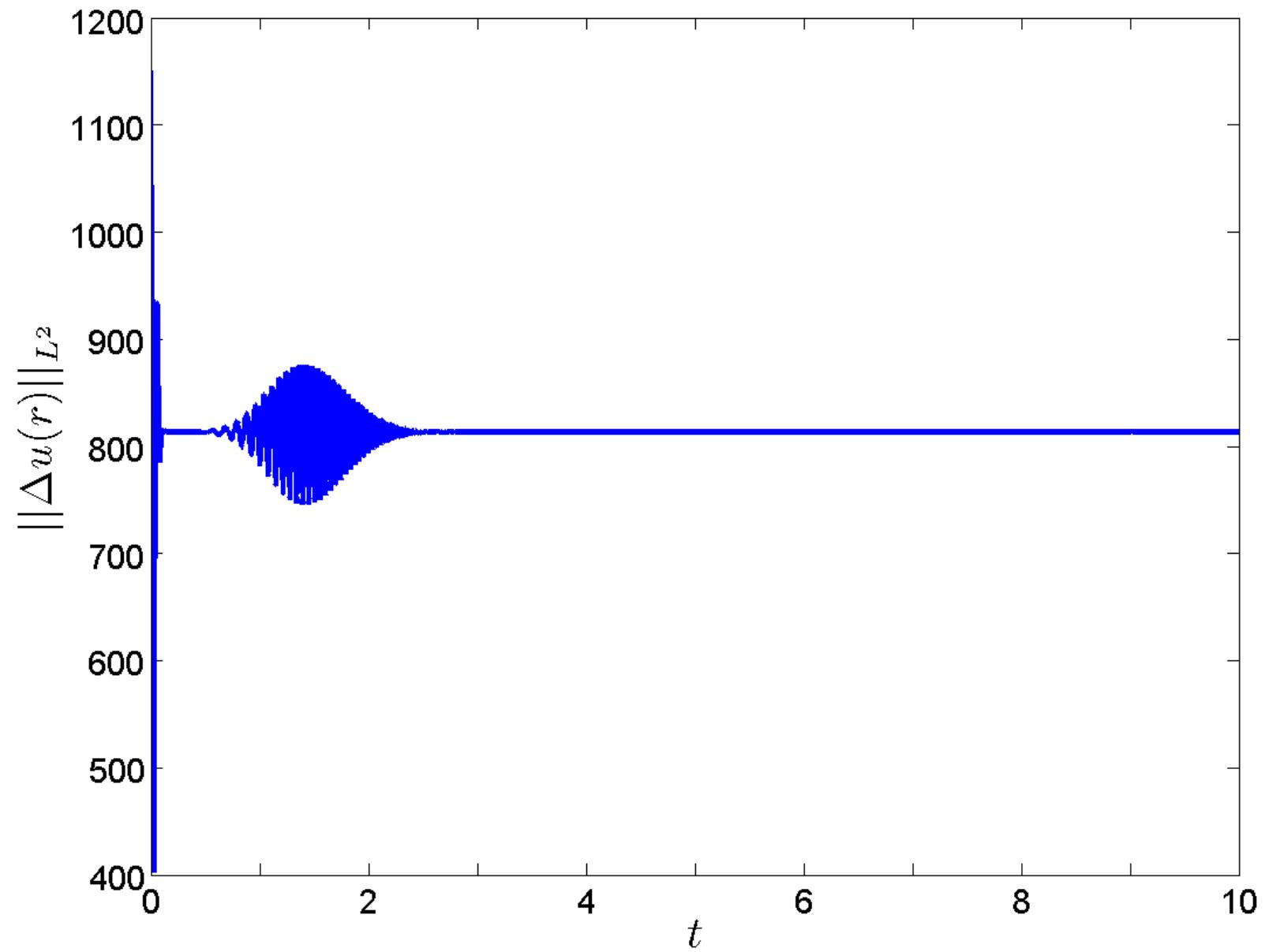


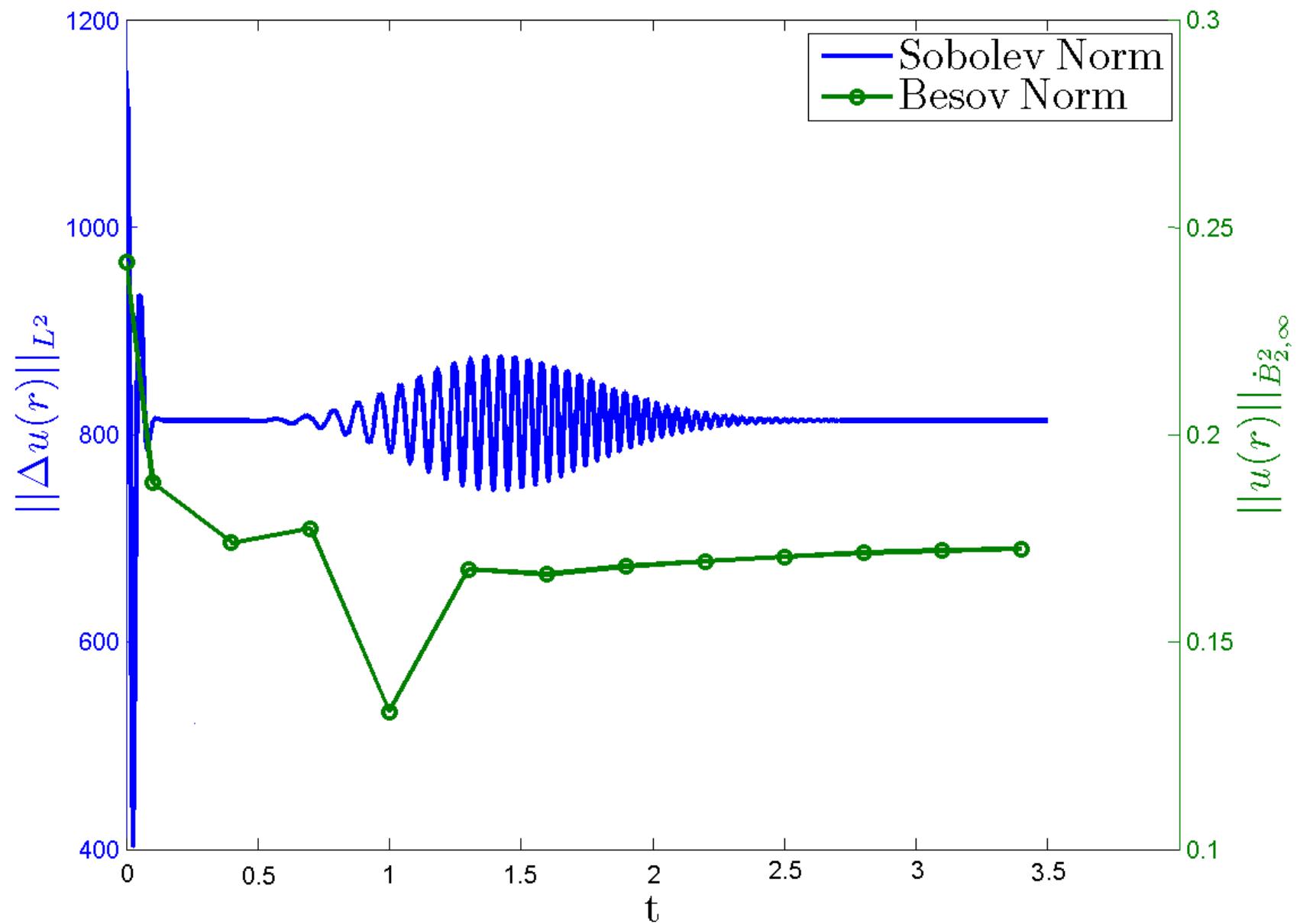
$t = 1.600000e+00$

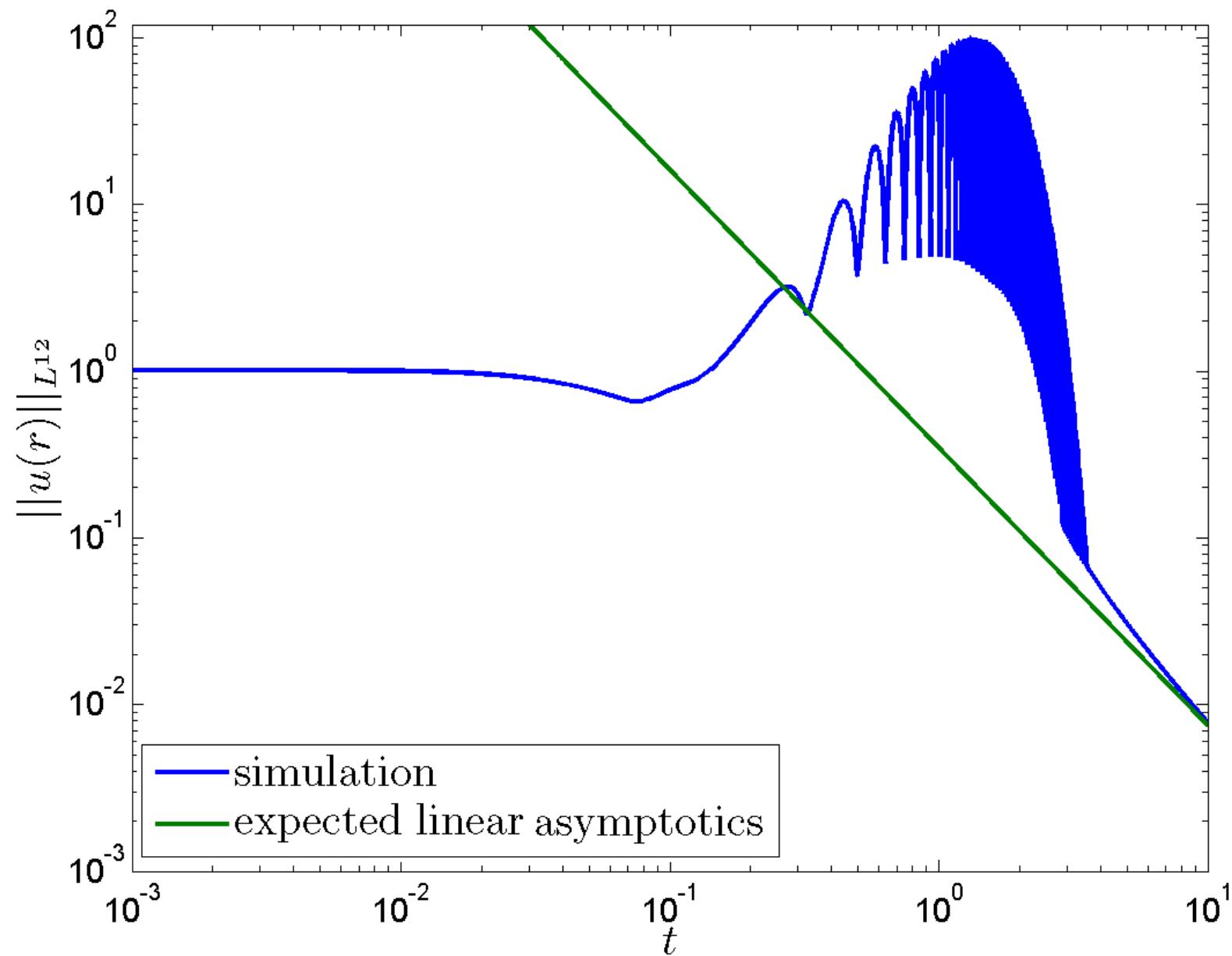


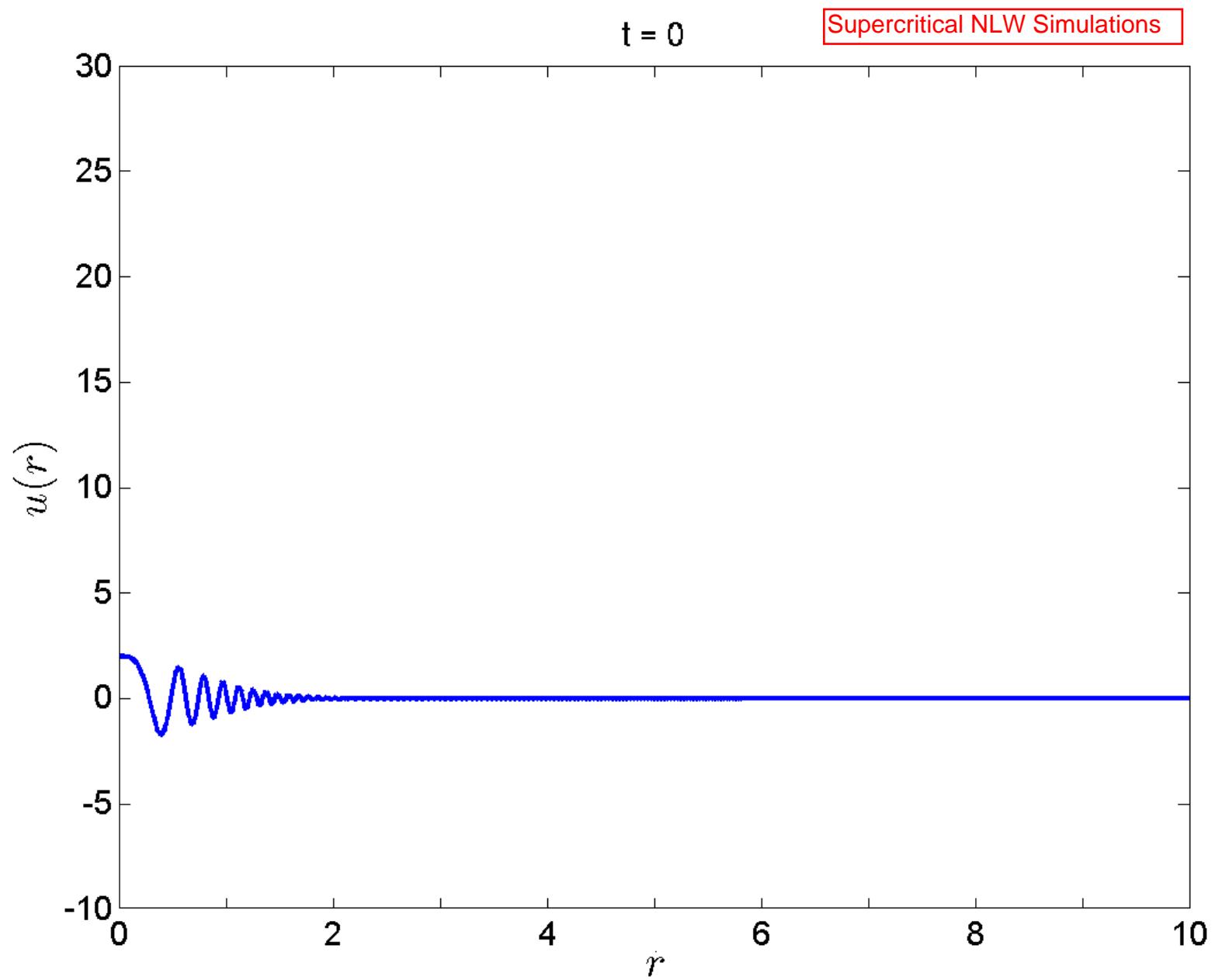


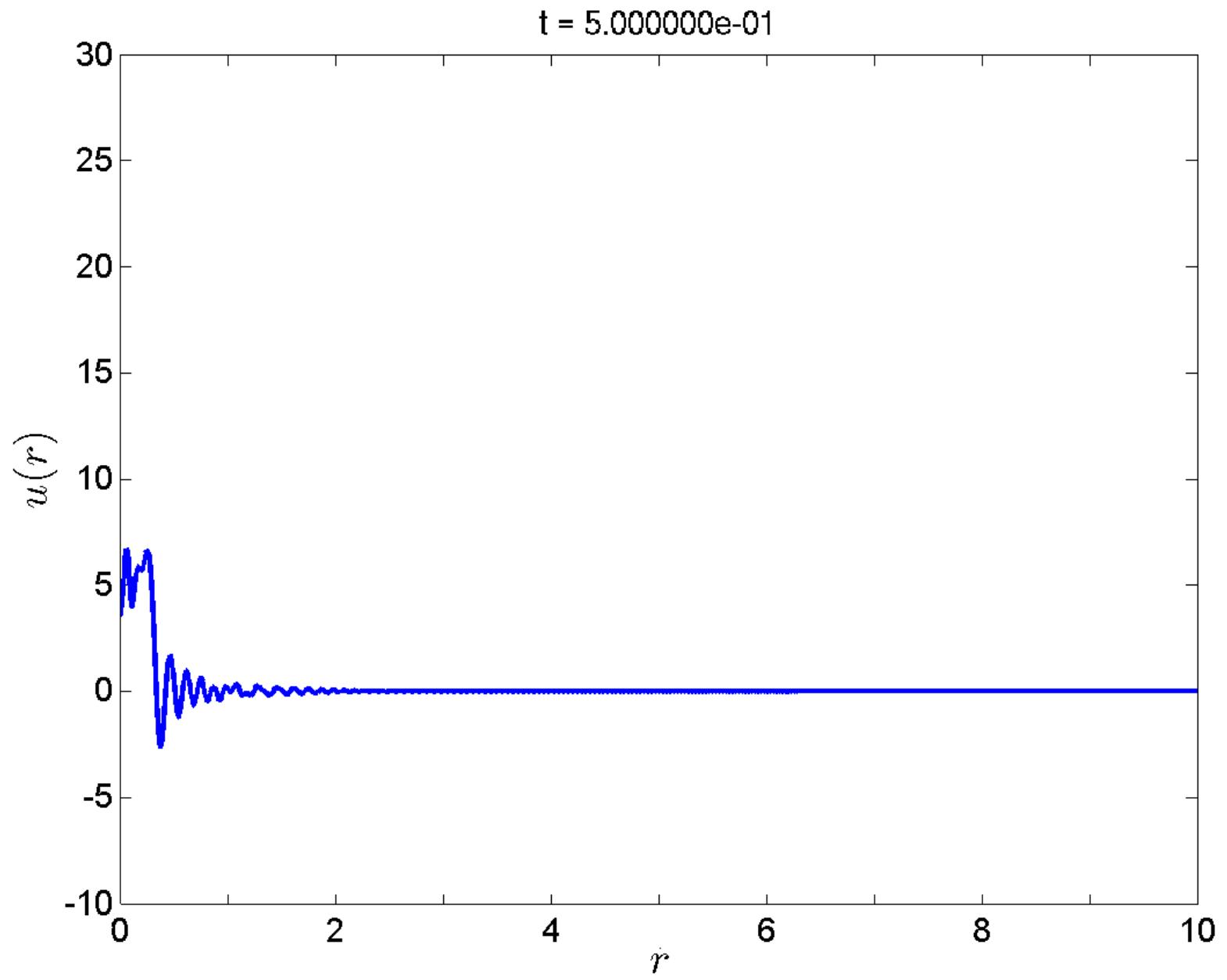


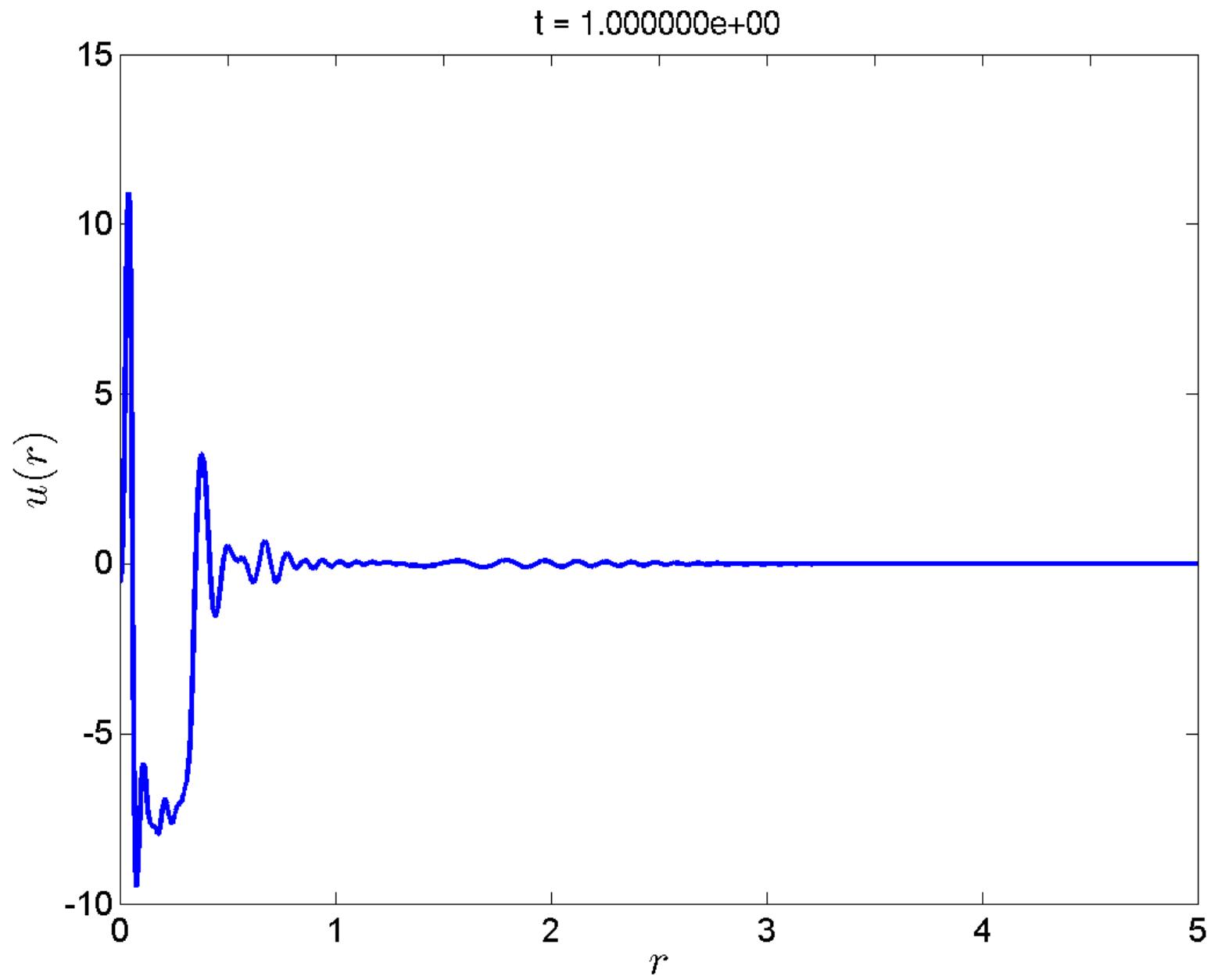


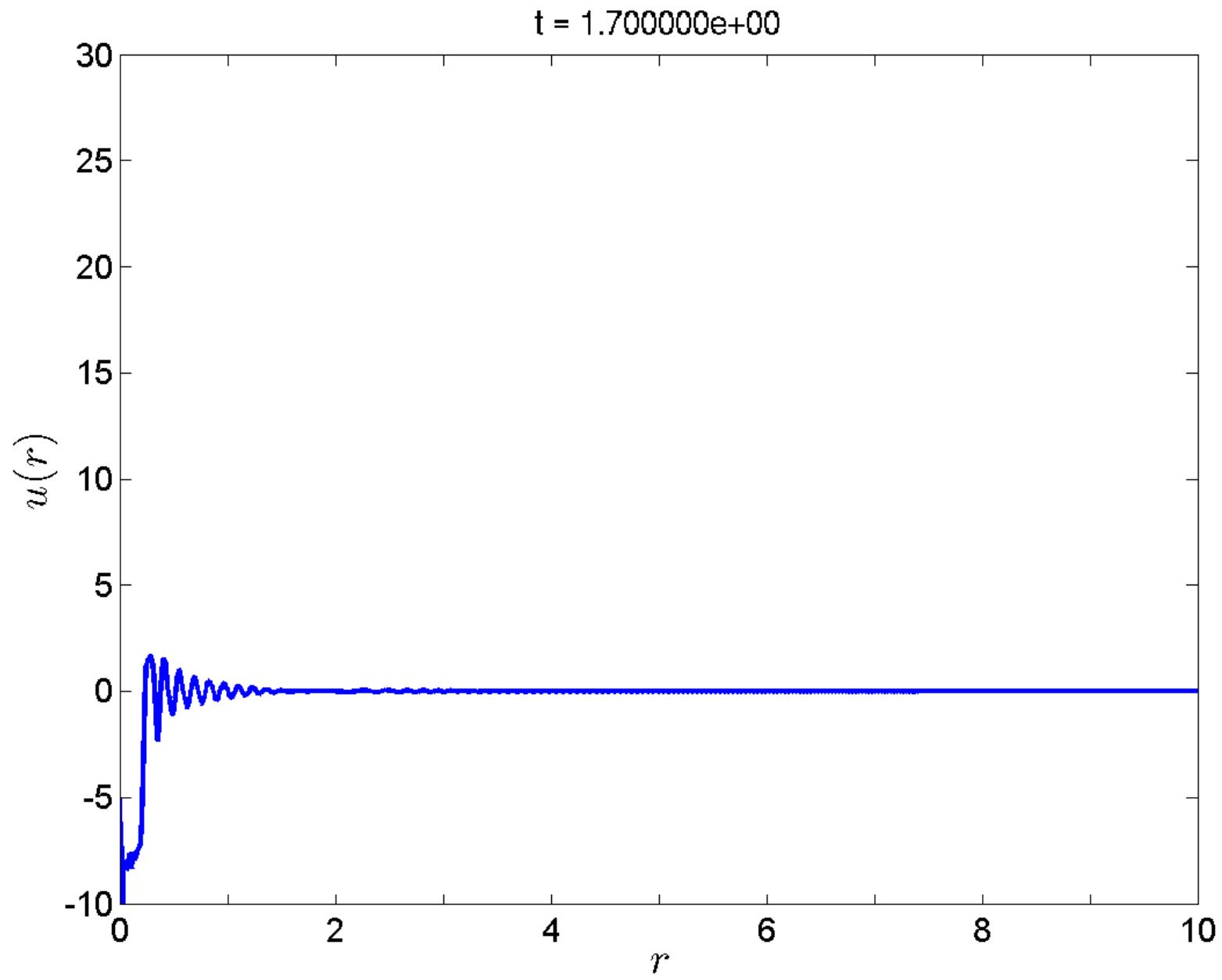


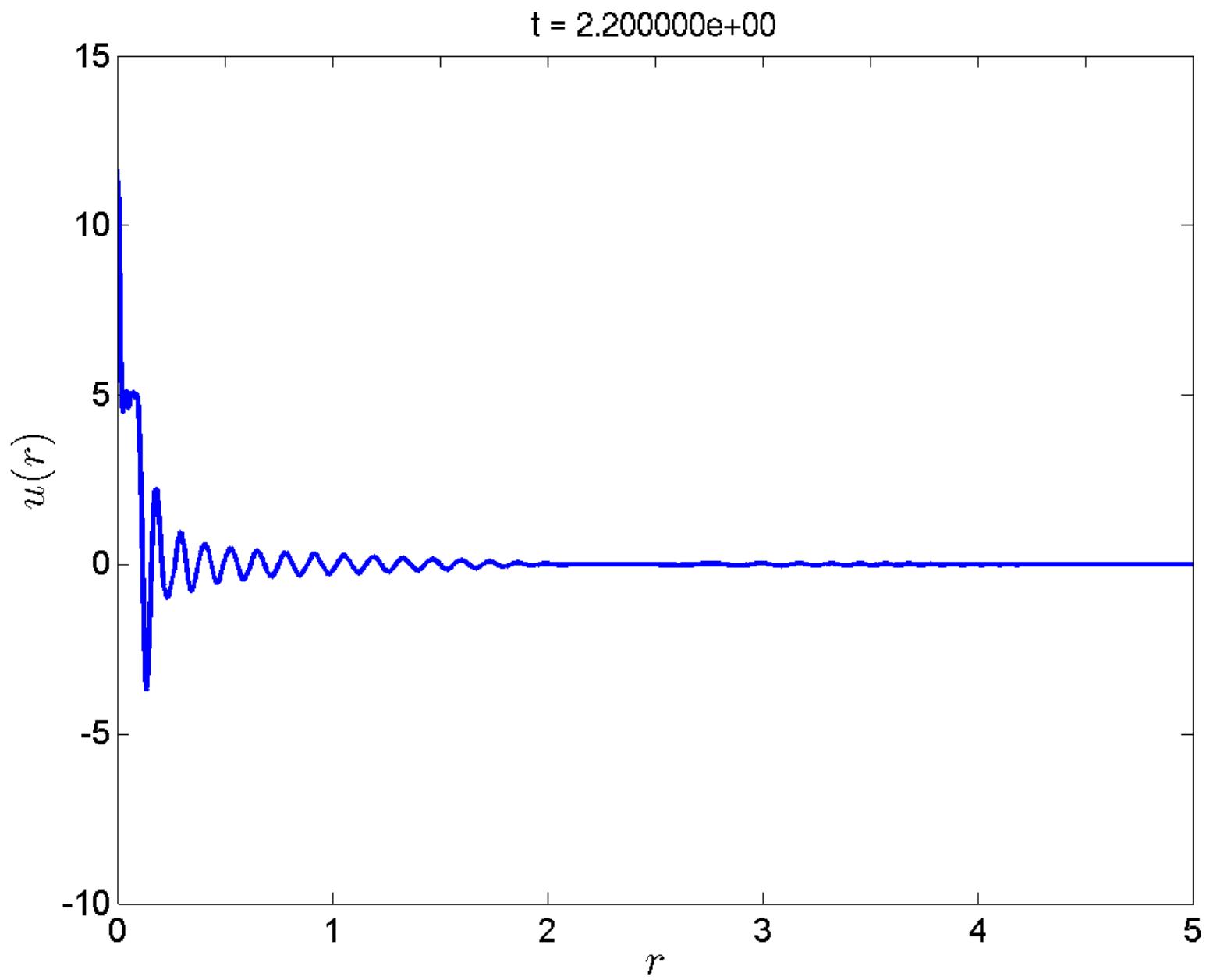


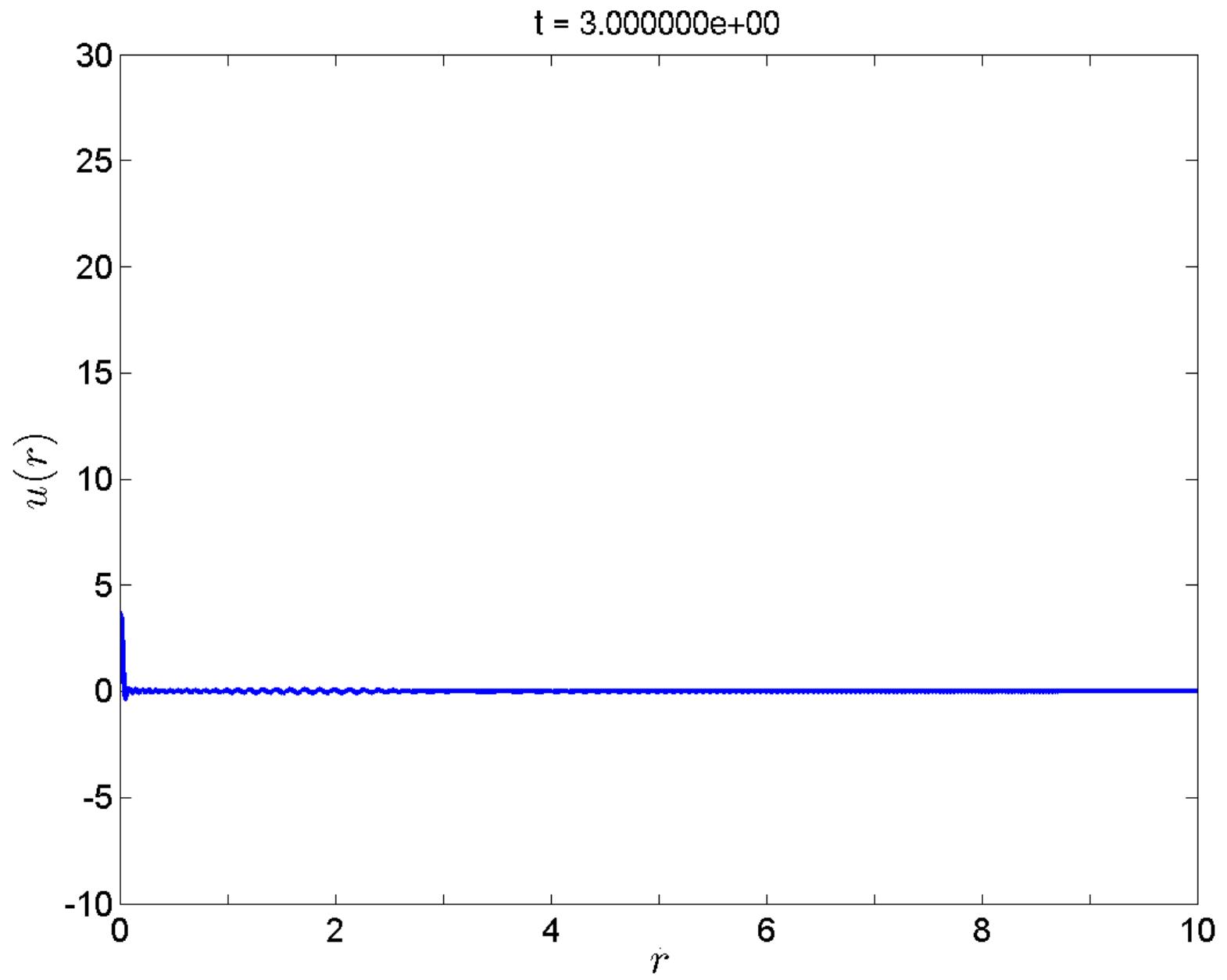


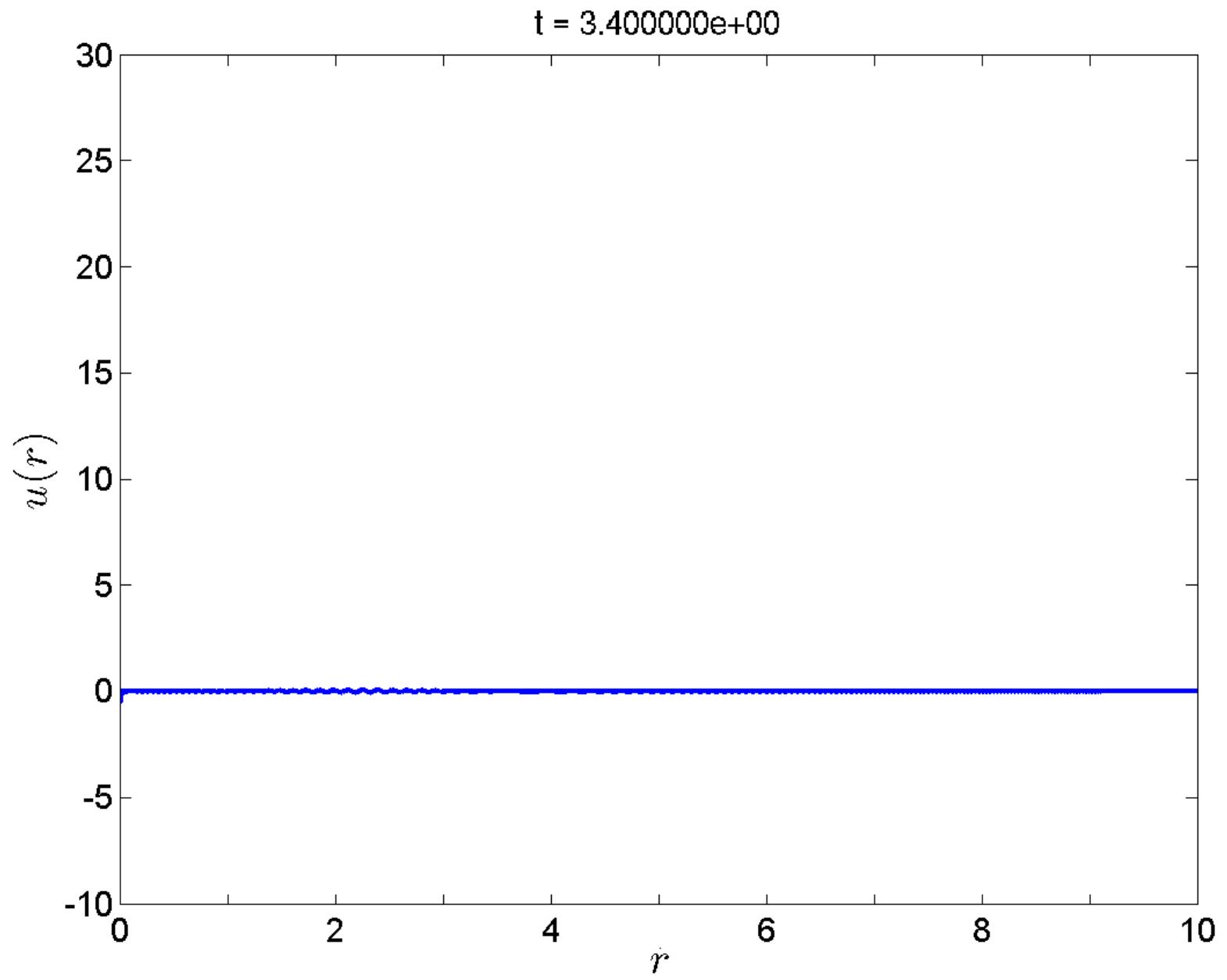


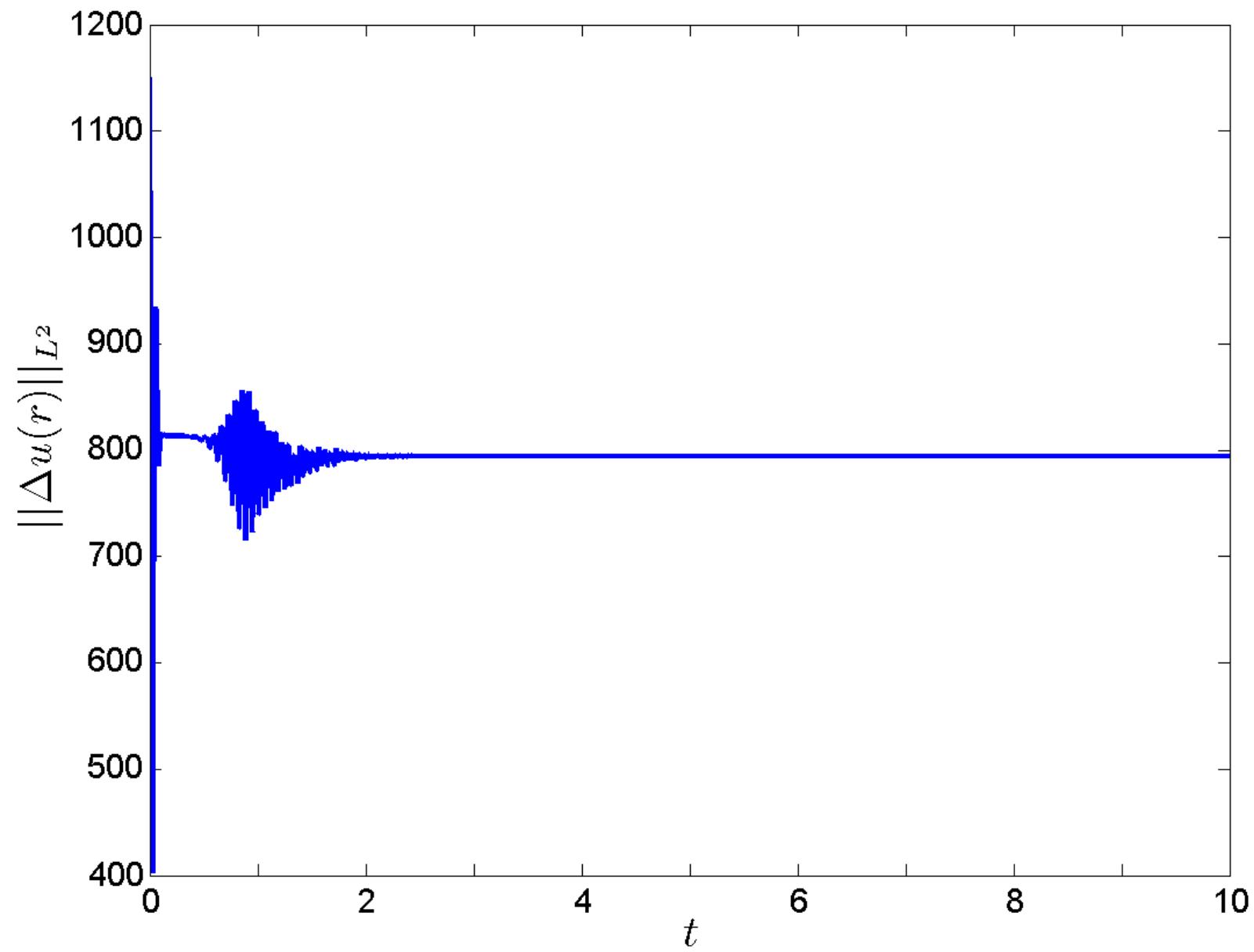


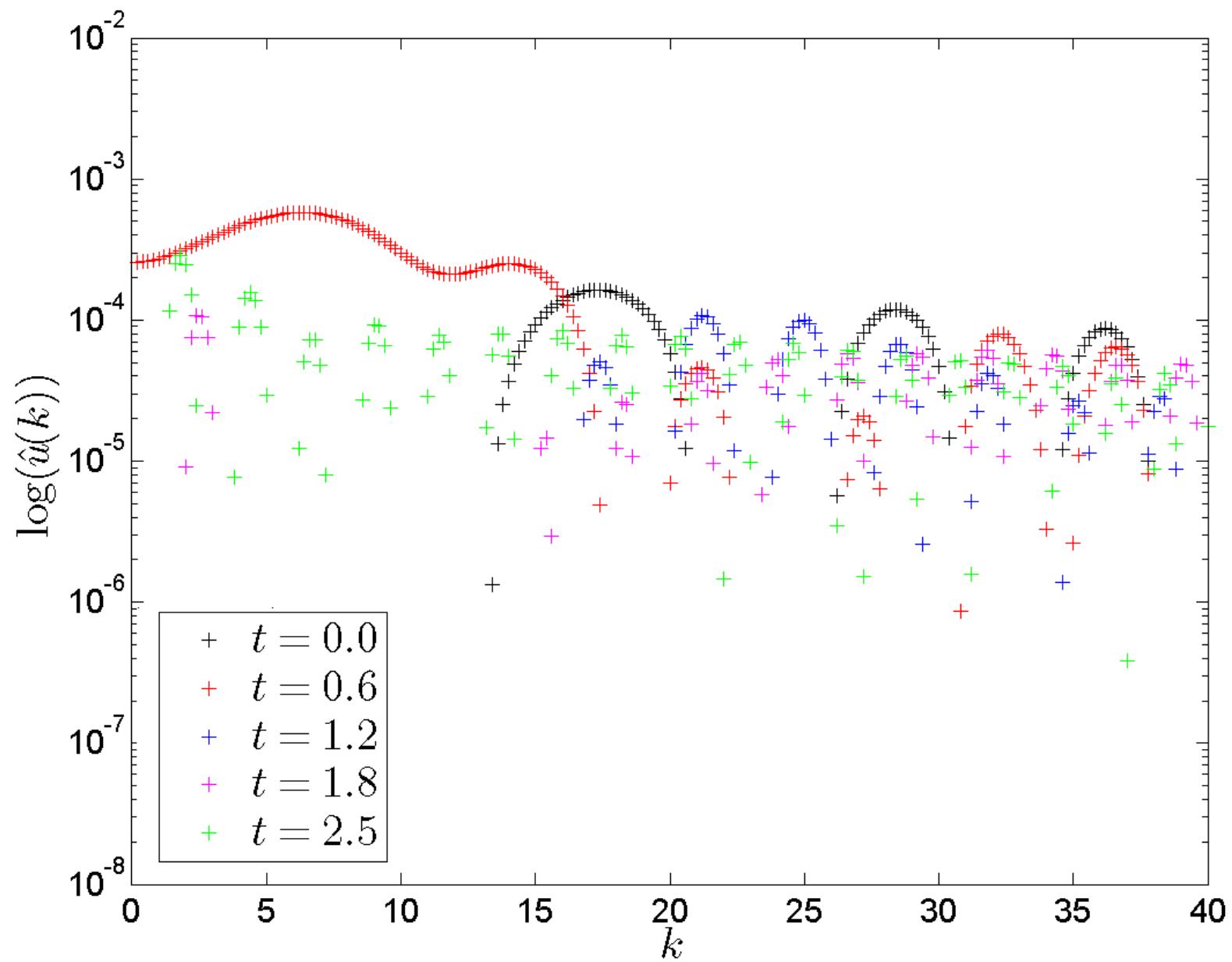


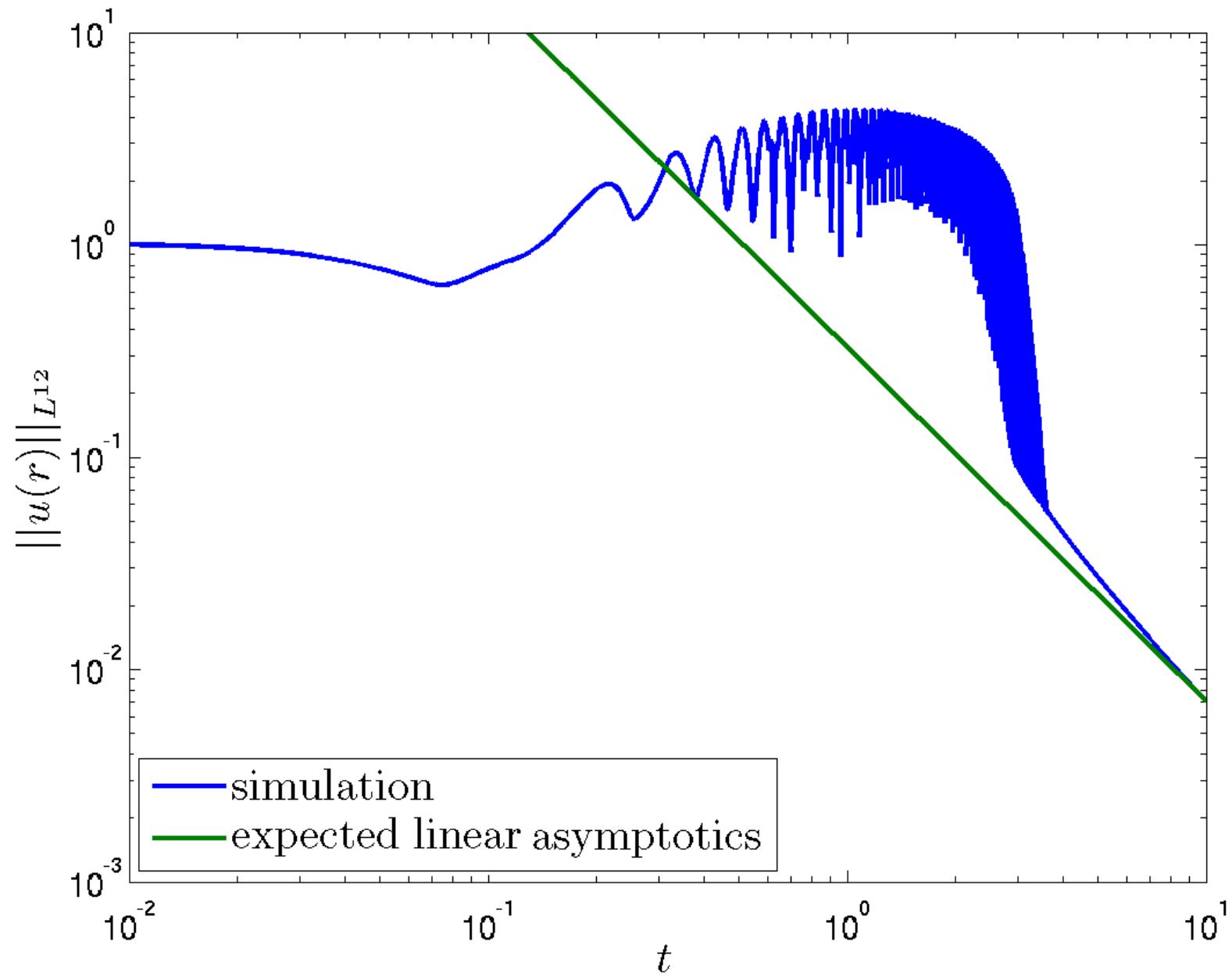












4. Ideas Towards a Supercritical Theory?

Ideas Towards a Supercritical Theory?

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GAFA Geometric And Functional Analysis

PROBLEMS IN HAMILTONIAN PDE'S

J. BOURGAIN

1 Introduction

The purpose of this exposé is to describe a line of research and problems, which I believe, will not be by any means completed in the near future. As such, we certainly hope to encourage further investigations. The list of topics in this field is fairly extensive and only a few will be commented on here. Their choice was mainly dictated by personal research involvement. It should also be mentioned that the different groups of researchers may have very different styles and aims. As a science, claims and results range from pure experimentation to rigorous mathematical proofs. **(Although my primary interest is this last aspect, I have no doubt that numerics or heuristic argumentation may be equally interesting and important.)** The history of the Korteweg-de-Vries equation for instance is a striking example of how a problem may evolve through these different interacting stages to eventually create a beautiful theory. As a mathematician I feel however that



Ideas Towards a Supercritical Theory?

PROBLEM. Is there global scattering in the energy space for $p = 2 + \frac{4}{d}$?

(See also [C1,2] for other results on scattering).

(iv) We like to sketch the theoretical possibility for computer assisted proofs of global existence and scattering, for a given data ϕ . Consider for instance the 3D supercritical problem

$$\begin{cases} iu_t + \Delta u - u|u|^6 = 0 \\ u(0) = \phi \end{cases} \quad (3.22)$$

where ϕ is a given smooth function. We do expect a global smooth solution + scattering. For this to hold, it is sufficient to show that for some time, $0 < T < \infty$,

- (3.22) has a smooth solution on $[0, T]$. Equivalently, $T^* > T$, where T^* refers to Theorem 3.7
- The norm $\|e^{i(t-T)\Delta} u(T)\|_{L_{t\geq T}^{15} L_x^{15}} < \delta$

where $\delta > 0$ is some numerical constant (we do not explain the role of the L^{15} -norm here). About step (a). If we fix a time T , one may establish the result numerically. To do this, one first gathers sufficiently many discrete data and interpolates them with a (smooth) function $v = v(x, t)$, $t < T$. Assuming (3.22) has indeed a smooth solution, the function v will

Ideas Towards a Supercritical Theory?

Numerical Observations:

- The critical Sobolev norm $\|u(t)\|_H^{s_c}$ stays bounded but is not monotone.
- The H^{s_c} density spreads out in Fourier space.
- The Besov $B_{2,\infty}^2$ norm shrinks to small values as time advances.



Ideas Towards a Supercritical Theory?

From [Bou98], the *small data* theory is relaxed:

$$\|(e^{it\Delta}\phi)(x)\|_{L^4(dxdt)} \leq C\|\phi\|_2^{1-\theta} \left\{ \sup_{A,\tau \in \mathcal{C}_A} \left\| \frac{1}{A} \widehat{\phi}|_\tau \right\|_1 \right\}^\theta.$$

(57)

Inequality (57) permits in particular to generate initial data $\phi \in L^2(\mathbb{R}^2)$ with arbitrarily large $\|\phi\|_2$ -norm such that $\|e^{it\Delta}\phi\|_4$ is small; hence the 2D-Cauchy problem

$$\begin{cases} iu_t + \Delta u + \lambda|u|^2u = 0 \\ u(0) = \phi \end{cases} \quad (58)$$

is globally well posed with a solution u in $C([- \infty, \infty]; L^2(\mathbb{R}^2)) \cap L^4([- \infty, \infty], L^4(\mathbb{R}^2))$. Observe that in the small data theory for (58), the smallness of the L^2 -norm $\|\phi\|_2$ is only used to ensure smallness of $\|e^{it\Delta}\phi\|_4$.

Ideas Towards a Supercritical Theory?

The motif from [Pla02] sets the stage:

1 Existence and regularity theorems

Theorem 1. Let $n = 2$, and $u_0 \in \dot{B}_2^{0,\infty}$ be such that

$$(4) \quad \|u_0\|_{\dot{B}_2^{0,\infty}} < \varepsilon_0.$$

Then there exists a global solution of (1) such that

$$(5) \quad u(x, t) \in L_t^\infty(\dot{B}_2^{0,\infty}),$$

$$(6) \quad u(x, t) \xrightarrow[t \rightarrow 0]{} u_0(x) \quad \text{weakly.}$$

Moreover, this solution is unique under the additional assumption

$$u \in \dot{B}_{4+}^{0-, \infty}(\mathcal{L}_t^{4-}) \cap \dot{B}_{4-}^{0+, \infty}(\mathcal{L}_t^{4+})$$

and

$$(7) \quad \sup_j (2^{0^+ j} \|\Delta_j u\|_{L_x^{4-}(L_t^{4+})} + 2^{0^- j} \|\Delta_j u\|_{L_x^{4+}(L_t^{4-})}) < \varepsilon_1.$$

small Besov



GWP

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