



# Computation without Representation

$\omega\epsilon\beta := \text{http://drorbn.net/o19/}$

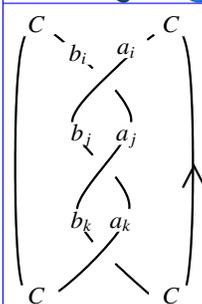
**Abstract.** A major part of “quantum topology” is the definition and computation of various knot invariants by carrying out computations in quantum groups. Traditionally these computations are carried out “in a representation”, but this is very slow: one has to use tensor powers of these representations, and the dimensions of powers grow exponentially fast.

In my talk, I will describe a direct method for carrying out such computations without having to choose a representation and explain why in many ways the results are better and faster. The two key points we use are a technique for composing infinite-order “perturbed Gaussian” differential operators, and the little-known fact that every semi-simple Lie algebra can be approximated by solvable Lie algebras, where computations are easier.

**KiW 43 Abstract** ( $\omega\epsilon\beta$ /kiw). Whether or not you like the formulas on this page, they describe the strongest truly computable knot invariant we know. (experimental analysis @ $\omega\epsilon\beta$ /kiw)

## Knotted Candies

$\omega\epsilon\beta$ /kc

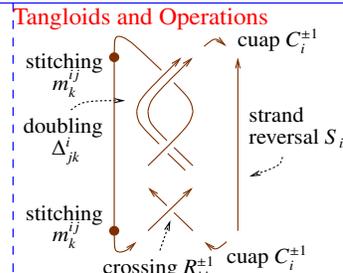


**The Yang-Baxter Technique.** Given an algebra  $U$  (typically  $\hat{U}(\mathfrak{g})$  or  $\hat{U}_q(\mathfrak{g})$ ) and elements

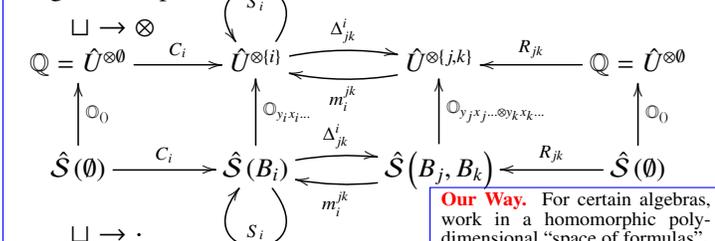
$$R = \sum a_i \otimes b_i \in U \otimes U \quad \text{and} \quad C \in U,$$
$$\text{form} \quad Z = \sum_{i,j,k} C a_i b_j a_k C^2 b_i a_j b_k C.$$

**Problem.** Extract information from  $Z$ .  
**The Dogma.** Use representation theory. In principle finite, but slow.

- A Knot Theory Portfolio.**
- Has operations  $\sqcup, m_{ij}^{ij}, \Delta_{jk}^i, S_i$ .
- All tangleoids are generated by  $R^{\pm 1}$  and  $C^{\pm 1}$  (so “easy” to produce invariants).
- Makes some knot properties (“genus”, “ribbon”) become “definable”.

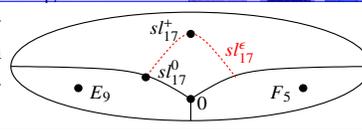


**A “Quantum Group” Portfolio** consists of a vector space  $U$  along with maps (and some axioms...)

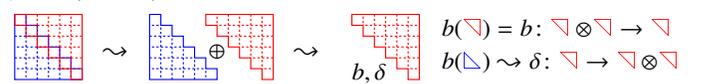


**Our Way.** For certain algebras, work in a homomorphic poly-dimensional “space of formulas”.

**The (fake) moduli** of Lie algebras on  $V$ , a quadratic variety in  $(V^*)^{\otimes 2} \otimes V$  is on the right. We care about  $sl_{17}^k := sl_{17}^\epsilon / (\epsilon^{k+1} = 0)$ .



**Solvable Approximation.** In  $gl_n$ , half is enough! Indeed  $gl_n \oplus \mathfrak{a}_n = \mathcal{D}(\nabla, b, \delta)$ :



Now define  $gl_n^\epsilon := \mathcal{D}(\nabla, b, \epsilon\delta)$ . Schematically, this is  $[\nabla, \nabla] = \nabla$ ,  $[\Delta, \Delta] = \epsilon\Delta$ , and  $[\nabla, \Delta] = \Delta + \epsilon\nabla$ . The same process works for all semi-simple Lie algebras, and at  $\epsilon^{k+1} = 0$  always yields a solvable Lie algebra.

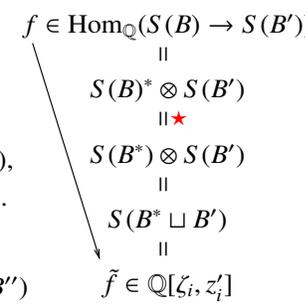
**CU and QU.** Starting from  $sl_2$ , get  $CU_\epsilon = \langle y, a, x, t \rangle / ([t, -] = 0, [a, y] = -y, [a, x] = x, [x, y] = 2\epsilon a - t)$ . Quantize using standard tools (I’m sorry) and get  $QU_\epsilon = \langle y, a, x, t \rangle / ([t, -] = 0, [a, y] = -y, [a, x] = x, xy - e^{\hbar\epsilon}yx = (1 - T e^{-2\hbar\epsilon a})/\hbar)$ .

**PBW Bases.** The  $U$ ’s we care about always have “Poincaré-Birkhoff-Witt” bases; there is some finite set  $B = \{y, x, \dots\}$  of “generators” and isomorphisms  $\mathbb{O}_{y,x,\dots}: \hat{S}(B) \rightarrow U$  defined by “ordering monomials” to some fixed  $y, x, \dots$  order. The quantum group portfolio now becomes a “symmetric algebra” portfolio, or a “power series” portfolio.

## Operations are Objects.

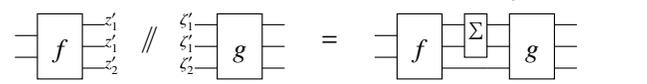
★  $B^* := \{z_i^m = \zeta_i^n : z_i \in B\}$ ,  $\langle z_i^m, \zeta_i^n \rangle = \delta_{mn} n!$ ,  $\langle \prod z_i^{m_i}, \prod \zeta_i^{n_i} \rangle = \prod \delta_{m_i n_i} n_i!$

in general, for  $f \in S(z_i)$  and  $g \in S(\zeta_i)$ ,  $\langle f, g \rangle = f(\partial_{z_i})g|_{z_i=0} = g(\partial_{z_i})f|_{z_i=0}$ .



## The Composition Law.

If  $S(B) \xrightarrow[\tilde{f} \in \mathbb{Q}[\zeta_i, z'_j]]{f} S(B') \xrightarrow[\tilde{g} \in \mathbb{Q}[\zeta'_j, z''_k]]{g} S(B'')$  then  $(\tilde{f}\tilde{g}) = (\tilde{g} \circ \tilde{f}) = \left( \tilde{g}|_{z'_j \rightarrow \partial_{z'_j}} \tilde{f} \right)_{z'_j=0} = \left( \tilde{f}|_{z'_j \rightarrow \partial_{z'_j}} \tilde{g} \right)_{z'_j=0}$



1. The 1-variable identity map  $I: S(z) \rightarrow S(z)$  is given by  $\tilde{I}_1 = \mathbb{P}z\zeta$  and the  $n$ -variable one by  $\tilde{I}_n = \mathbb{P}z_1\zeta_1 + \dots + z_n\zeta_n$ :

$$\tilde{I}_1 = \square + \square + \frac{1}{2} \square + \frac{1}{6} \square + \dots$$

2. The “archetypal multiplication map  $m_k^{ij}: S(z_i, z_j) \rightarrow S(z_k)$ ” has  $\tilde{m} = \mathbb{P}z_k(\zeta_i + \zeta_j)$ .
3. The “archetypal coproduct  $\Delta_{jk}^i: S(z_i) \rightarrow S(z_j, z_k)$ ”, given by  $z_i \rightarrow z_j + z_k$  or  $\Delta z = z \otimes 1 + 1 \otimes z$ , has  $\tilde{\Delta} = \mathbb{P}(z_j + z_k)\zeta_i$ .
4.  $R$ -matrices tend to have terms of the form  $e^{\hbar y_1 x_2} \in \mathcal{U}_q \otimes \mathcal{U}_q$ . The “baby  $R$ -matrix” is  $\tilde{R} = e^{\hbar y x} \in S(y, x)$ .
5. The “Weyl form of the canonical commutation relations” states that if  $[y, x] = tI$  then  $e^{\xi x} e^{\eta y} = e^{\eta y} e^{\xi x} e^{-\eta\xi t}$ . So with

$$sw_{xy} \left( S(y, x) \xrightarrow[\mathbb{O}_{yx}]{\mathbb{O}_{xy}} \mathcal{U}(y, x) \right) \text{ we have } \tilde{SW}_{xy} = \mathbb{P}^{\eta y + \xi x - \eta\xi t}$$