

Example. $G =$ the Heisenberg group.

$V = \mathbb{R}^n = V^*$ via $\langle \cdot, \cdot \rangle =$ standard inner product

$$G = \left\{ \begin{bmatrix} 1 & a & c \\ 0 & I & b \\ 0 & 0 & 1 \end{bmatrix} \right\} \quad \begin{array}{l} a = (a_1, \dots, a_n) \in V^* \\ b = \begin{pmatrix} b_1 \\ \vdots \\ b_n \end{pmatrix} \in V \\ c \in \mathbb{R} \end{array}$$

$$\begin{bmatrix} 1 & a & c \\ 0 & I & b \\ 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & a' & c' \\ 0 & I & b' \\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & a+a' & c+c'+\langle a, b' \rangle \\ 0 & I & b+b' \\ 0 & 0 & 1 \end{bmatrix}$$

$$1 \longrightarrow \mathbb{R} \longrightarrow G \longrightarrow (\mathbb{R}^{2n}, +) \longrightarrow 1$$

central extension

Its Lie algebra: $\mathfrak{g} = \left\{ \begin{bmatrix} 0 & \alpha & \gamma \\ 0 & 0 & \beta \\ 0 & 0 & 0 \end{bmatrix} \right\} \quad \begin{array}{l} \alpha = (\alpha_1, \dots, \alpha_n) \\ \beta = \begin{pmatrix} \beta_1 \\ \vdots \\ \beta_n \end{pmatrix} \end{array} \quad \gamma \in \mathbb{R}$

$\{\text{mathfrak{g}}\}$

$$\exp\left(\begin{bmatrix} 0 & \alpha & \gamma \\ 0 & 0 & \beta \\ 0 & 0 & 0 \end{bmatrix}\right) = \begin{bmatrix} 1 & \alpha & \gamma + \frac{1}{2}\langle \alpha, \beta \rangle \\ 0 & I & \beta \\ 0 & 0 & 1 \end{bmatrix}$$

$$\text{Ad}\left(\begin{bmatrix} 1 & a & c \\ 0 & I & b \\ 0 & 0 & 1 \end{bmatrix}\right) \begin{bmatrix} 0 & \alpha & \gamma \\ 0 & 0 & \beta \\ 0 & 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & \alpha & \gamma + \langle a, \beta \rangle - \langle a, b \rangle \\ 0 & 0 & \beta \\ 0 & 0 & 0 \end{bmatrix}$$

Adjoint orbits :

$$\begin{cases} d\alpha \times d\beta \times d\mathbb{R} & \text{if } \alpha, \beta \text{ not both } 0 \\ d\alpha \times d\alpha \times d\mathbb{R} & \text{if } \alpha = \beta = 0 \end{cases}$$

Their dimension: 1 or 0.

⚡ not symplectic!

$$\mathfrak{g}^* \cong \mathbb{R}^{2n+1} \ni (\alpha^*, \beta^*, \gamma^*) : \begin{pmatrix} 0 & \alpha & \gamma \\ 0 & 0 & \beta \\ 0 & 0 & 0 \end{pmatrix} \mapsto \langle \alpha^*, \alpha \rangle + \langle \beta^*, \beta \rangle + \langle \gamma^*, \gamma \rangle$$

$$\alpha^*, \beta^* \in \mathbb{R}^n, \gamma^* \in \mathbb{R}$$

The coadjoint action:

$$\left\langle \underbrace{\text{Ad}^* \begin{pmatrix} 1 & a & c \\ 0 & I & b \\ 0 & 0 & 1 \end{pmatrix}}_{\text{in } \mathfrak{g}^*} (\alpha^*, \beta^*, \gamma^*), \underbrace{\begin{pmatrix} 0 & \alpha & \gamma \\ 0 & 0 & \beta \\ 0 & 0 & 0 \end{pmatrix}}_{\text{in } \mathfrak{g}} \right\rangle$$

$$= \left\langle \underbrace{(\alpha^*, \beta^*, \gamma^*)}_{\text{in } \mathfrak{g}^*}, \underbrace{\text{Ad} \left(\begin{bmatrix} 1 & -a & -c + \langle a, b \rangle \\ 0 & I & -b \\ 0 & 0 & 1 \end{bmatrix} \right) \begin{pmatrix} 0 & \alpha & \gamma \\ 0 & 0 & \beta \\ 0 & 0 & 0 \end{pmatrix}}_{\text{in } \mathfrak{g}} \right\rangle$$

$$= \left\langle (\alpha^*, \beta^*, \gamma^*), \begin{bmatrix} 0 & \alpha & \gamma - \langle a, \beta \rangle + \langle a, b \rangle \\ 0 & 0 & \beta \\ 0 & 0 & 0 \end{bmatrix} \right\rangle$$

$$= \underbrace{\langle \alpha^*, \alpha \rangle}_{\mathbb{R}} + \underbrace{\langle \beta^*, \beta \rangle}_{\mathbb{R}^n} + \gamma^* (\gamma - \langle a, \beta \rangle + \langle a, b \rangle)$$

$$\begin{matrix} \uparrow \uparrow \\ \mathbb{R} \quad \mathbb{R}^n \end{matrix} = \langle \alpha^* + \gamma^* b, \alpha \rangle + \langle \beta^* - \gamma^* a, \beta \rangle + \gamma^* \gamma$$

$$\text{Thus, } \underbrace{\text{Ad}^* \begin{pmatrix} 1 & a & c \\ 0 & I & b \\ 0 & 0 & 1 \end{pmatrix}}_{\text{in } \mathfrak{g}} \underbrace{(\alpha^*, \beta^*, \gamma^*)}_{\text{in } \mathfrak{g}^*} = (\underbrace{\alpha^* + \gamma^* b}_{\mathbb{R}^n}, \underbrace{\beta^* - \gamma^* a}_{\mathbb{R}^n}, \underbrace{\gamma^*}_{\mathbb{R}})$$

Coadjoint orbits: $\mathbb{R}^{2n} \times \{\gamma^*\}$ if $\gamma^* \neq 0$
 singletons $\{(\alpha^*, \beta^*, 0)\}$ if $\gamma^* = 0$

← dim = 2n

← dim = 0

when $\gamma \neq 0$: $\omega_{\text{FKS}} = \pm \gamma^* \cdot \omega_{\text{std}}$
 $\mathcal{L}$ on $\mathbb{R}^{2n}$

$H := (\mathbb{R}^{2n}, +) \hookrightarrow \mathbb{R}^{2n}$ by translation

does not have an (equivariant) momentum map
 i.e. is not a Hamiltonian action.

proof The group is abelian
 so Ad^* is trivial

so $\mu: \mathbb{R}^{2n} \rightarrow \mathfrak{h}^*$ being equivariant
 means being invariant.

The action is transitive
 so μ invariant $\Rightarrow \mu$ constant
 $\Rightarrow$ the action is trivial.

Recall:

$$1 \rightarrow \mathbb{R} \rightarrow G \xrightarrow{\pi} H \rightarrow 1.$$

$\underbrace{\hspace{1.5cm}}_{\text{the Heisenberg group}} \quad \underbrace{\hspace{1.5cm}}_{(\mathbb{R}^{2n}, +)}$

Let $G \xrightarrow{\pi} H \hookrightarrow \mathbb{R}^{2n}$

nonfaithful action by translations.

does have an equivariant momentum map.

up to $\pm$:

$$\mu: \mathbb{R}^{2n} \rightarrow \mathfrak{g}^* = \left\{ \overset{\mathbb{R}^n}{\alpha^*}, \overset{\mathbb{R}^n}{\beta^*}, \overset{\mathbb{R}}{\gamma^*} \right\}$$

$$(x, y) \mapsto (-y, x, -1)$$

$\mu^1 \quad \mu^2 \quad \mu^3$

is this a momentum map?

assume $n=1$. ^{elements} basis for $\mathfrak{g}$ act by $\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, 0$.

$$\left(\begin{array}{ccc} \mathfrak{G} & \longrightarrow & \mathfrak{H} \\ \begin{pmatrix} 1 & a & c \\ 0 & I & b \\ 0 & 0 & 1 \end{pmatrix} & \longmapsto & \begin{array}{l} \text{translations} \\ x \mapsto x+a \\ y \mapsto y+b \end{array} \end{array} \right)$$

$$2\left(\frac{\partial}{\partial x}\right) dx \wedge dy = dy = -d\mu^1$$

$$2\left(\frac{\partial}{\partial y}\right) dx \wedge dy = -dx = -d\mu^2$$

$$2(0) dx \wedge dy = 0 = -d\mu^3$$

$\therefore$ Hamilton's equation holds.

Equivariance:

$$(x, y) \xrightarrow{\mu} \begin{matrix} \alpha^* & \beta^* & \gamma^* \\ (-y, x, -1) \end{matrix}$$

$$\begin{pmatrix} 1 & a & c \\ 0 & I & b \\ 0 & 0 & 1 \end{pmatrix} :$$

translating $\downarrow$

$$(x+a, y+b)$$

$\downarrow$ coadjoint

$$(-y-b, x+a, -1)$$



Compact Lie groups

Invariant measures on compact Lie groups:

Let G be a compact Lie group.

Then there exists a (unique) smooth probability measure on G that is invariant under left & right translations and under inversion.

Left translations are the maps $L_a: g \mapsto ag$ for $a \in G$
 Right translations are the maps $R_a: g \mapsto ga$ — " —
 inversion is the map $g \mapsto g^{-1}$.

idea. trivialize $TG \cong G \times \mathfrak{g}$ using
 $a \cdot \xi \leftrightarrow (a, \xi)$

def M is parallelizable if
 $TM \cong M \times \mathbb{R}^n$

Let $n = \dim G$.
 Take any n -covector at 1:

$$0 \neq v \in \Lambda^n \mathfrak{g}^*$$

Sweep it by left translation to get a left invariant volume form.

$$vol \in \Omega^n(G), \quad vol|_a = (L_a^{-1})^* v$$

$$L_a: G \rightarrow G$$

$$1 \mapsto a$$

$$(L_a)_*: \mathfrak{g} \mapsto T_a \mathfrak{g}$$

$$(L_a)^*: \mathfrak{g}^* \mapsto T_a^* \mathfrak{g}$$

$$\Lambda^n \mathfrak{g}^* \mapsto \Lambda^n T_a^* \mathfrak{g}$$

divide by $\int_G vol$
 to get
 a new volume
 form
 with $\int_G vol = 1$

the orientation
 determined by vol

The measure of a "nice" set A is $\int_A \text{vol}$
 wrt the orientation determined by vol .

Claim. This is the unique left invt volume form
 with $\int_G \text{vol} = 1$. up to $\pm$

But vol & $-\text{vol}$ induce the same measure
 b/c we integrate each wrt its orientation.

call this measure μ . Then for nice $f: G \rightarrow \mathbb{R}$

$$\int f d\mu = \int f \cdot \text{vol}$$

$\nwarrow$ wrt the orientation determined by vol .

claim $d\mu$ is invt under right translations
 & inversion.

Because left & right translations commute

$\forall a \quad R_a^* \text{vol}$ is also left-invt

This $\Rightarrow \quad R_a^* \text{vol} = \pm \text{vol}$

define the same measure

so $d\mu$ is invt under right translations too.

Since vol is right-invariant, $(\text{inversion})^* \text{vol}$ is
 left-invariant, so $(\text{inversion})^* \text{vol} = \pm \text{vol}$ & again
 we get the same measure.

Examples:

$$G = S^1 = \{e^{i\theta}\}$$

$$\text{vol} = \frac{d\theta}{2\pi}$$

$$(\text{inversion})^* \text{vol} = -\text{vol}.$$

$$G = \{\pm 1\} \ltimes S^1$$

$$(\epsilon_1, z_1) \cdot (\epsilon_2, z_2) = \begin{cases} (\epsilon_1 \epsilon_2, z_1 z_2) & \text{if } \epsilon_2 = 1 \\ (\epsilon_1 \epsilon_2, z_1^{-1} z_2) & \text{if } \epsilon_2 = -1 \end{cases}$$

$\epsilon_1, \epsilon_2 \in \{\pm 1\}$
 $|z_1| = |z_2| = 1$

conjugation by $(\epsilon, 1)$ with $\epsilon = -1$

takes (ϵ, z) to (ϵ, z')

so it flips orientation.

so vol cannot be invariant under both left & right translation.

The measure is given by vol
1-density

on an n dim'l mfd M :

top degree diff'l forms are sections of $\wedge^n T^*M$.

$\hat{M}$

$\downarrow$
 M

the orientation double covering.

$\hat{M}_1$

M_1

$\longrightarrow$

$\Uparrow$ lifts

diffeomorphism

$\hat{M}_2$

M_2

$$|\Lambda^n T^*M| := \underbrace{\Lambda^n T^*M}_{\text{rank 1 real line bundle}} \otimes_{\mathbb{Z}_2} \underbrace{\hat{M}}_{\text{rank 1 real line bundle}}$$

fibrewise!
over each point of M :
 $\mathbb{R} \otimes_{\mathbb{Z}_2} \mathbb{Z}_2 = \mathbb{R}$

a 1-density of M is a section of $|\Lambda^n T^*M|$
 $n = \dim M$.

over a small open set: it's represented
 by a pair

$$(\alpha, \theta) \sim (-\alpha, -\theta)$$

$\nearrow$ n -form $\uparrow$ orientation $\uparrow$ the opposite orientation

For any compactly supported density φ , $\int_M \varphi$ is well defined.

even if M was not orientable.
 (in our case $M = G$ is parallelizable)
 hence orientable.

$\nexists$ volume form vol
 determines a density $|\text{vol}|$.

Averaging

Lemma. G $\xrightarrow{\text{linearly}}$ V
 compact Lie group finite dim'd real vector space.

Then $\exists G$ -invariant inner product $\langle, \rangle$.

Proof. Take any inner product $\langle, \rangle'$.

Note: $\forall a \in G \quad (u, v) \mapsto \langle a \cdot u, a \cdot v \rangle'$
 is also an inner product

$$\langle u, v \rangle := \int_{a \in G} \langle a \cdot u, a \cdot v \rangle' \quad da$$

an invariant probability measure

is then a G -invariant inner product on V .

Special case. G a compact Lie group

$$\text{Ad}: G \rightarrow \mathfrak{g}$$

Then $\exists \text{Ad}$ -invariant inner product $\langle, \rangle$ on $\mathfrak{g}$.

Which we use to identify $\mathfrak{g}$ with $\mathfrak{g}^*$:

$$\mathfrak{g} \xrightarrow{\sim} \mathfrak{g}^*$$

$$\xi \mapsto \langle \xi, \cdot \rangle$$

an Ad -invariant inner product on $\mathfrak{g}$

Claim this linear isomorphism intertwines $\text{Ad}: G \rightarrow \mathfrak{g}$ with $\text{Ad}^*: G \rightarrow \mathfrak{g}^*$.

indeed: $\forall a \in G$

$$\text{Ad}(a)(\xi) \longmapsto \langle \text{Ad}(a)(\xi), \cdot \rangle$$

$$\begin{aligned} &= \langle \xi, \text{Ad}(a^{-1}) \cdot \rangle \\ &= \text{Ad}^*(a) \langle \xi, \cdot \rangle \end{aligned}$$

because $\langle \cdot, \cdot \rangle$ is Ad-invariant

∴ choice of Ad-invariant inner product on $\mathfrak{g}$ identifies adjoint orbits with coadjoint orbits.
 so adjoint orbits are even dimensional.

Example. $G = \text{SU}(n) = \{A \in \mathbb{C}^{n \times n} \mid \overline{A}^T A = I \text{ and } \det A = 1\}$

$$= \left\{ A \in \mathbb{C}^{n \times n} \mid \begin{aligned} &H(Au, Av) = H(u, v) \\ &\forall u, v \in \mathbb{C}^n \text{ and } \det A = 1 \end{aligned} \right\}$$

where $H(z, w) = \sum_{j=1}^n \overline{z}_j w_j$.

$\mathfrak{g} = \mathfrak{su}(n) = \left\{ \begin{array}{l} \text{traceless} \\ \text{anti-Hermitian matrices} \end{array} \right\}$

$$= \{A \in \mathbb{C}^{n \times n} \mid \overline{A}^T + A = 0 \text{ and } \text{trace } A = 0\}.$$

Can take $\langle A, B \rangle = -\text{Re}(\text{Trace } AB)$

discussion

$$\text{Tr}(AB) = \text{Tr}(\bar{A}^T B) = - \sum_j (\bar{A}^T B)_{jj}$$

$$= - \sum_{j,l} (\bar{A}^T)_{jl} B_{lj}$$

$$= - \sum_{j,l} \bar{A}_{lj} B_{lj}$$